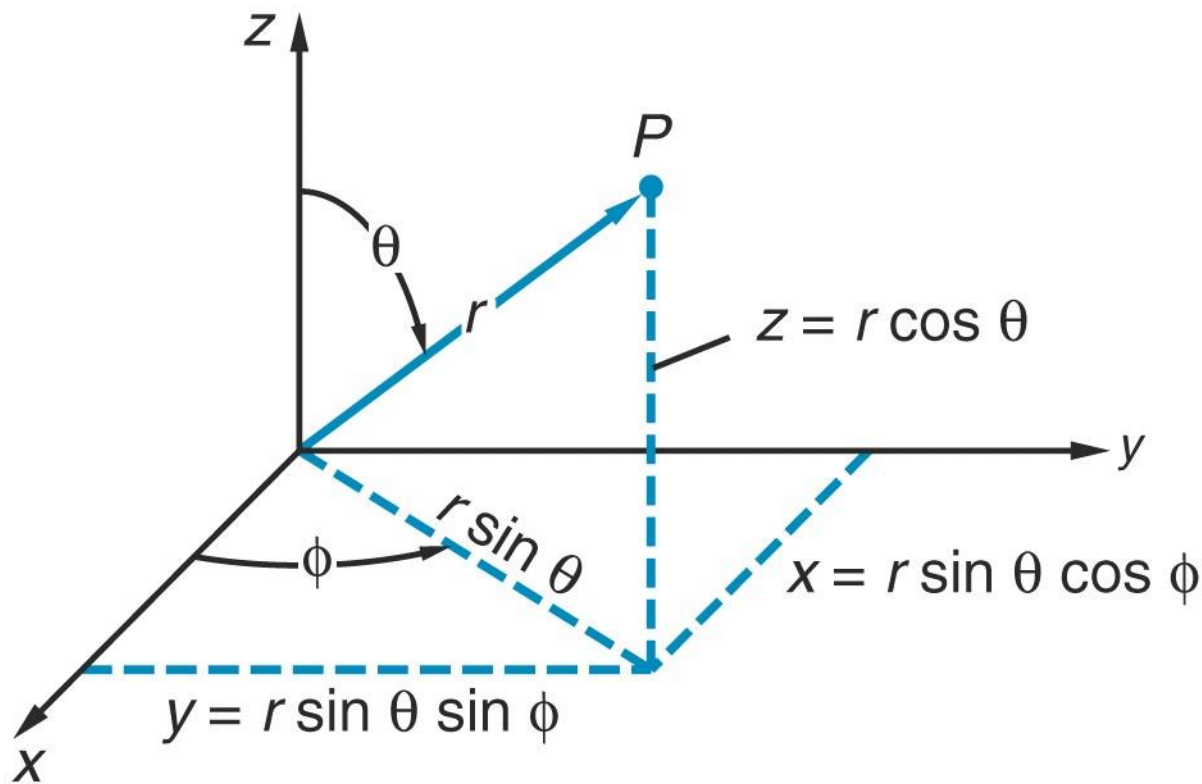


# Spherical Coordinates



Range of variables

Cartesian

$x, y, z: -\infty \rightarrow +\infty$

Spherical

$r: 0 \rightarrow +\infty$

$\theta: 0 \rightarrow \pi$

$\phi: 0 \rightarrow 2\pi$

$$\frac{\partial^2 F}{\partial x^2} + \frac{\partial^2 F}{\partial y^2} + \frac{\partial^2 F}{\partial z^2} = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial F}{\partial r} \right) + \frac{1}{r^2 \sin(\theta)} \frac{\partial}{\partial \theta} \left( \sin(\theta) \frac{\partial F}{\partial \theta} \right) + \frac{1}{r^2 \sin^2(\theta)} \frac{\partial^2 F}{\partial \phi^2}$$

# Angular Momentum, Classical

The definition of angular momentum is

$$\vec{L} = \vec{r} \times \vec{p} \quad L_x = y p_z - z p_y \quad L_y = z p_x - x p_z \quad L_z = x p_y - y p_x$$

Properties of cross products?

Direction perpendicular to both vectors.

Right hand rule.

Magnitude is the product of magnitudes times sin of angle between

Cyclic nature & distributive property

$$\hat{i} = \hat{j} \times \hat{k} \quad \hat{k} = \hat{i} \times \hat{j} \quad \hat{j} = \hat{k} \times \hat{i}$$

$$\hat{i} = -\hat{k} \times \hat{j} \quad \hat{k} = -\hat{j} \times \hat{i} \quad \hat{j} = -\hat{i} \times \hat{k}$$

# Question #1

Classical: An object at position  $\mathbf{r}$  has an angular momentum  $\mathbf{L} = \mathbf{r} \times \mathbf{p}$ . If  $\mathbf{L}$  is constant which statement is correct?

- (a) The force on the object must be 0.
- (b) The force in the direction of  $\mathbf{r}$  must be 0.
- (c) The force perpendicular to the direction  $\mathbf{r}$  must be 0.
- (d) None of the above are correct.

## Question #2

Classical: Which statement is true for an object at position  $\mathbf{r}(t)$  with  $\mathbf{L}$  constant?  $|\mathbf{r}| = r = (x^2 + y^2 + z^2)^{1/2}$

- (a) The magnitude of the force only depends on  $r$ .
- (b) The potential energy,  $U$ , only depends on  $r$ .
- (c) The potential energy perpendicular to  $r$  must be 0.
- (d) The force must be a constant.
- (e) The potential energy depends linearly on  $x$  or  $y$  or  $z$ .

# Angular Momentum, Important

The eigenstate of one component is not an eigenstate of the other two components (except  $j = n = m = 0$ ). **It is not possible to have a quantum system with defined  $L_x$ ,  $L_y$ , and  $L_z$  at the same time.**

By convention  $L_z$  is almost always the special, well defined component of  $\mathbf{L}$ .

$m$  is an integer

In spherical coordinates [ $x = r \sin(\theta) \cos(\phi)$ ,  $y = r \sin(\theta) \sin(\phi)$ ] the eigenstate has simple form.

$$(x + i y)/r = \sin(\theta) e^{i \phi} \quad \Rightarrow \quad [(x + i y)/r]^m = \sin^m(\theta) e^{i m \phi}$$

# Total Angular Momentum, Quantum

The squared angular momentum arises from summing and squaring the components. In terms of spherical coordinates

( $x = r \sin(\theta) \cos(\phi)$ ,  $y = r \sin(\theta) \sin(\phi)$ ,  $z = r \cos(\theta)$  ),  $L^2$ , has a relatively simple form

$$L_{\text{op}}^2 = -\hbar^2 \left[ \frac{1}{\sin(\theta)} \frac{\partial}{\partial \theta} \left( \sin(\theta) \frac{\partial}{\partial \theta} \right) + \frac{1}{\sin^2(\theta)} \frac{\partial^2}{\partial \phi^2} \right]$$

The eigenstates are the spherical harmonics which are also eigenstates of  $L_z$ . The eigenvalues are  $\ell(\ell+1)\hbar^2$

$$L_{\text{op}}^2 Y_{\ell m}(\theta, \phi) = \hbar^2 \ell(\ell+1) Y_{\ell m}(\theta, \phi)$$

$$L_z Y_{\ell m}(\theta, \phi) = \hbar m Y_{\ell m}(\theta, \phi)$$

# Spherical Harmonics

TABLE 7-1 Spherical harmonics

$l = 0$	$m = 0$	$Y_{00} = \sqrt{\frac{1}{4\pi}}$
$l = 1$	$m = 1$	$Y_{11} = -\sqrt{\frac{3}{8\pi}} \sin \theta e^{i\phi}$
	$m = 0$	$Y_{10} = \sqrt{\frac{3}{4\pi}} \cos \theta$
	$m = -1$	$Y_{1-1} = \sqrt{\frac{3}{8\pi}} \sin \theta e^{-i\phi}$
$l = 2$	$m = 2$	$Y_{22} = \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{2i\phi}$
	$m = 1$	$Y_{21} = -\sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{i\phi}$
	$m = 0$	$Y_{20} = \sqrt{\frac{5}{16\pi}} (3 \cos^2 \theta - 1)$
	$m = -1$	$Y_{2-1} = \sqrt{\frac{15}{8\pi}} \sin \theta \cos \theta e^{-i\phi}$
	$m = -2$	$Y_{2-2} = \sqrt{\frac{15}{32\pi}} \sin^2 \theta e^{-2i\phi}$

Note: A 3-D color representation of the spherical harmonics is on the Internet at <http://ww3.uniovi.es/~quimica.fisica/qeg/harmonics/charmonics.html>

# Angular Momentum, Values

$\ell$  can be 0, 1, 2, 3, ...

$m$  can go from  $-\ell$  to  $\ell$  and is an integer

The number of oscillations in the  $\phi$  coordinate is  $|m|$ .

The number of oscillations in the  $\theta$  coordinate is  $\ell - |m|$

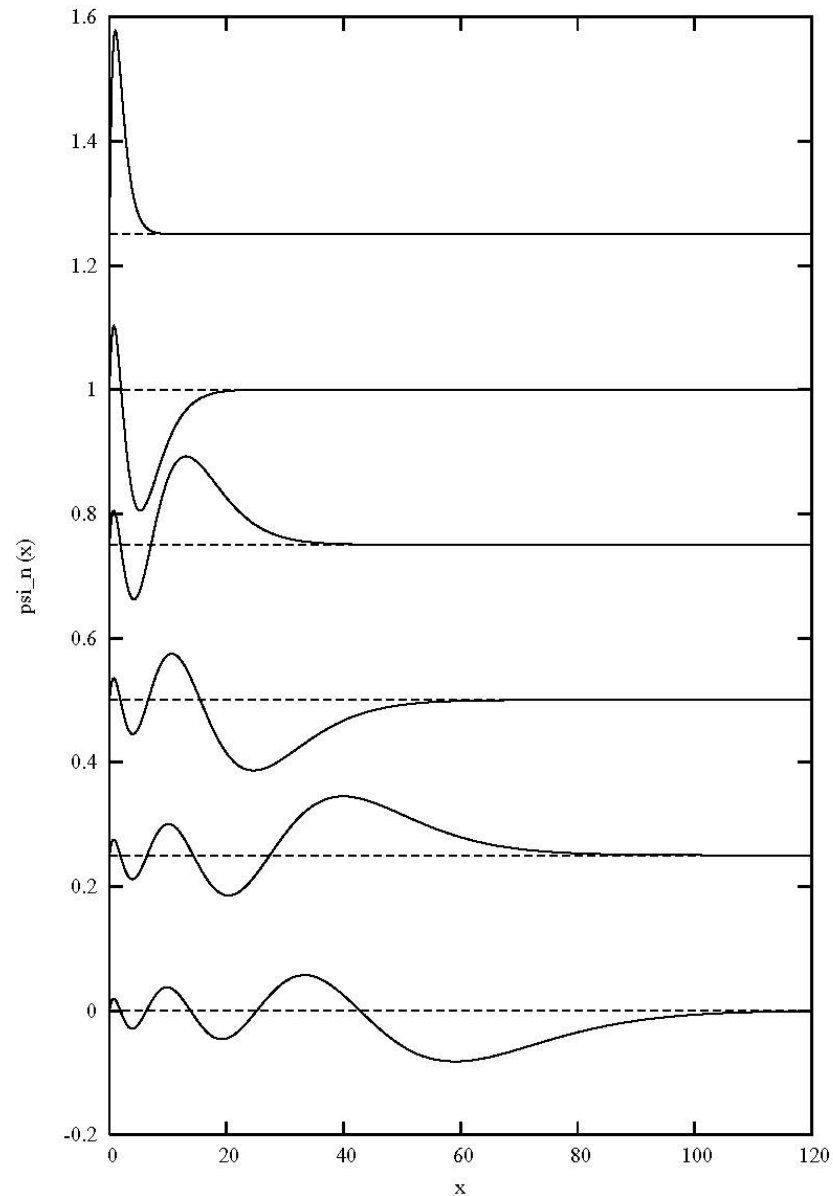


**TABLE 7.1 Some Hydrogen Atom Wave Functions**

$n$	$l$	$m_l$	$R(r)$	$\Theta(\theta)$	$\Phi(\phi)$
1	0	0	$\frac{2}{a_0^{3/2}} e^{-r/a_0}$	$\frac{1}{\sqrt{2}}$	$\frac{1}{\sqrt{2\pi}}$
2	0	0	$\frac{1}{(2a_0)^{3/2}} \left(2 - \frac{r}{a_0}\right) e^{-r/2a_0}$	$\frac{1}{\sqrt{2}}$	$\frac{1}{\sqrt{2\pi}}$
2	1	0	$\frac{1}{\sqrt{3}(2a_0)^{3/2}} \frac{r}{a_0} e^{-r/2a_0}$	$\sqrt{\frac{3}{2}} \cos \theta$	$\frac{1}{\sqrt{2\pi}}$
2	1	$\pm 1$	$\frac{1}{\sqrt{3}(2a_0)^{3/2}} \frac{r}{a_0} e^{-r/2a_0}$	$\mp \frac{\sqrt{3}}{2} \sin \theta$	$\frac{1}{\sqrt{2\pi}} e^{\pm i\phi}$
3	0	0	$\frac{2}{(3a_0)^{3/2}} \left(1 - \frac{2r}{3a_0} + \frac{2r^2}{27a_0^2}\right) e^{-r/3a_0}$	$\frac{1}{\sqrt{2}}$	$\frac{1}{\sqrt{2\pi}}$
3	1	0	$\frac{8}{9\sqrt{2}(3a_0)^{3/2}} \left(\frac{r}{a_0} - \frac{r^2}{6a_0^2}\right) e^{-r/3a_0}$	$\sqrt{\frac{3}{2}} \cos \theta$	$\frac{1}{\sqrt{2\pi}}$
3	1	$\pm 1$	$\frac{8}{9\sqrt{2}(3a_0)^{3/2}} \left(\frac{r}{a_0} - \frac{r^2}{6a_0^2}\right) e^{-r/3a_0}$	$\mp \frac{\sqrt{3}}{2} \sin \theta$	$\frac{1}{\sqrt{2\pi}} e^{\pm i\phi}$
3	2	0	$\frac{4}{27\sqrt{10}(3a_0)^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0}$	$\sqrt{\frac{5}{8}} (3 \cos^2 \theta - 1)$	$\frac{1}{\sqrt{2\pi}}$
3	2	$\pm 1$	$\frac{4}{27\sqrt{10}(3a_0)^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0}$	$\mp \sqrt{\frac{15}{4}} \sin \theta \cos \theta$	$\frac{1}{\sqrt{2\pi}} e^{\pm i\phi}$
3	2	$\pm 2$	$\frac{4}{27\sqrt{10}(3a_0)^{3/2}} \frac{r^2}{a_0^2} e^{-r/3a_0}$	$\frac{\sqrt{15}}{4} \sin^2 \theta$	$\frac{1}{\sqrt{2\pi}} e^{\pm 2i\phi}$

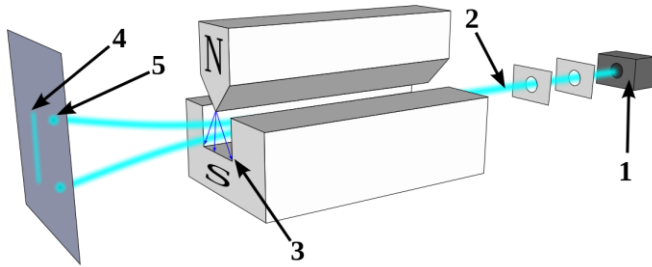
# Coulomb Potential $(-1/r)$ , $\ell = 0$

$r R(r)$

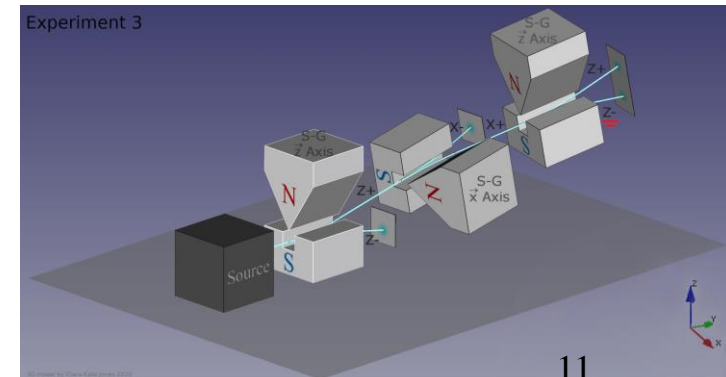
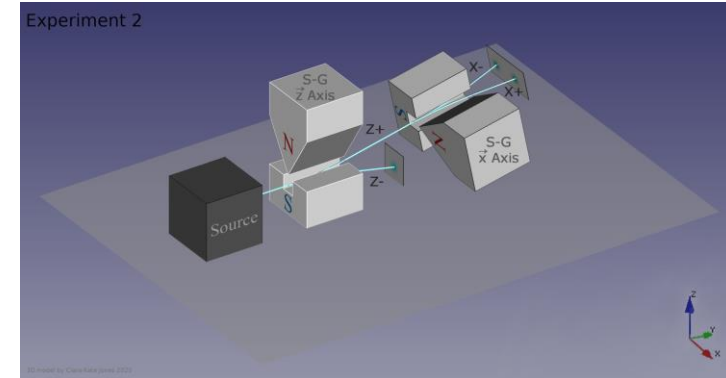
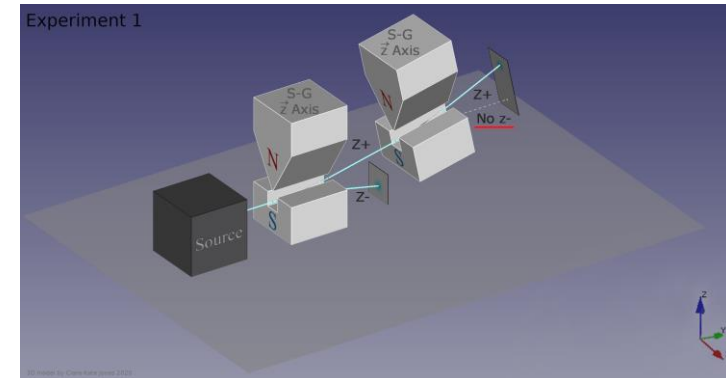
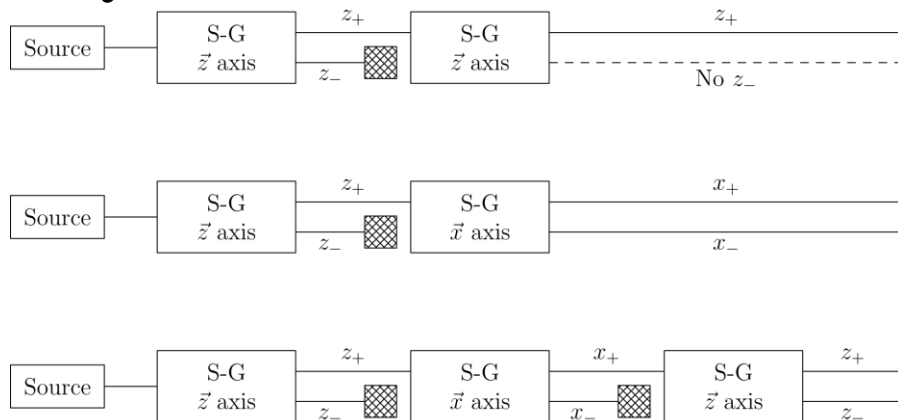


# Stern-Gerlach Experiment

Basic for spin  $\frac{1}{2}$



Can show the spin  $\frac{1}{2}$  for any direction

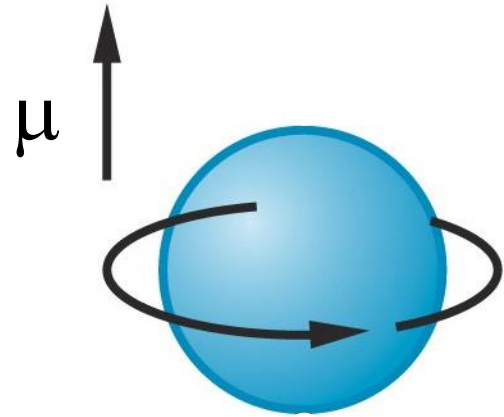


# Trapping Neutral Particles

$$PE = - \vec{\mu} \bullet \vec{B}$$

$$\vec{\mu} \bullet \hat{B} \cong \text{constant}$$

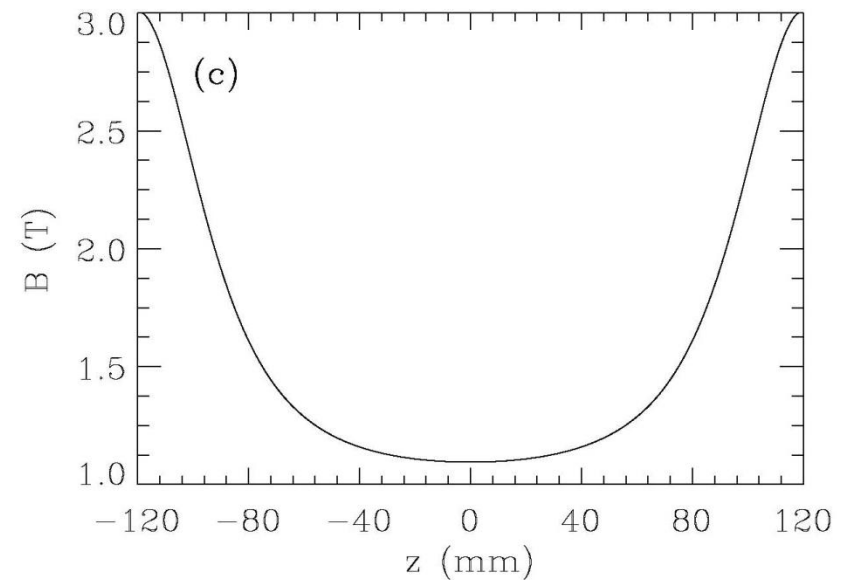
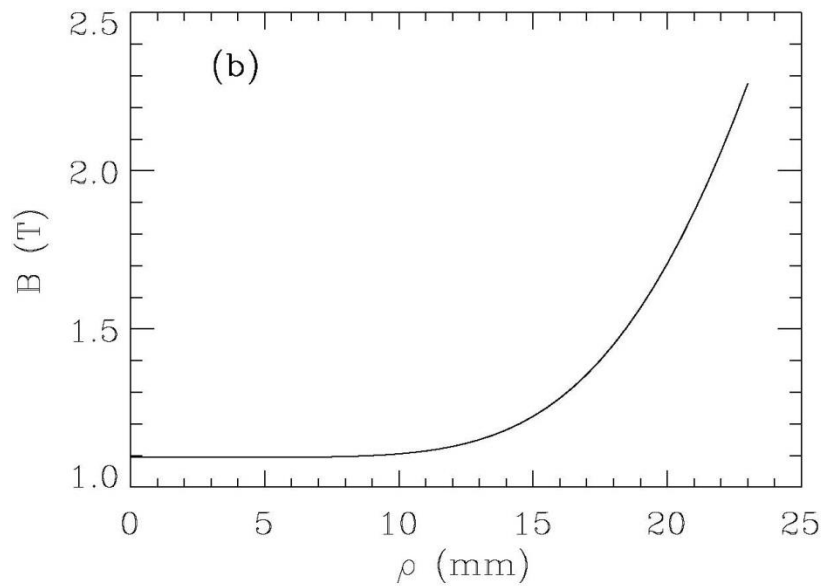
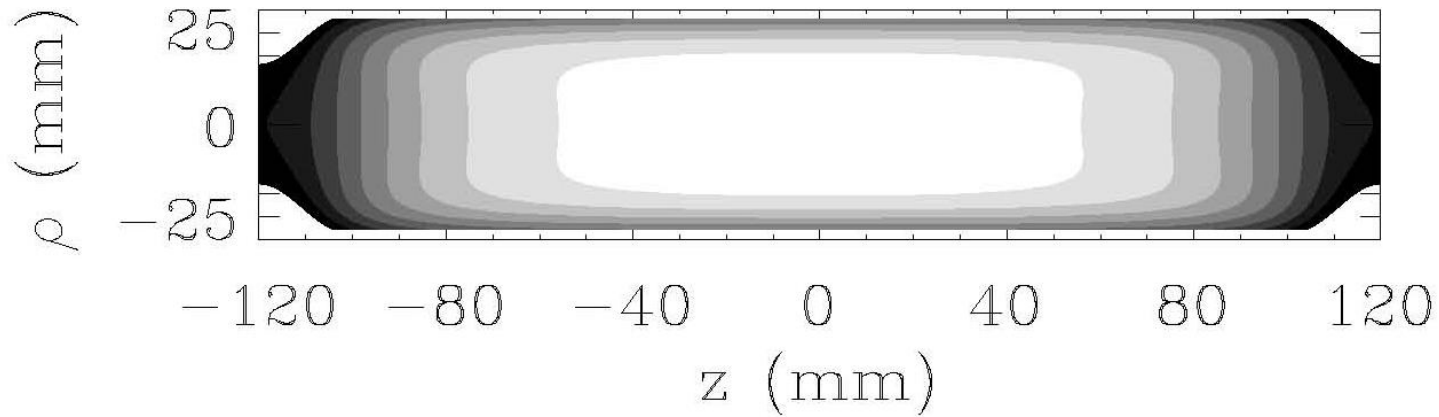
$$PE = - \left( \vec{\mu} \bullet \hat{B} \right) B(x, y, z)$$



μ for ground state ~2/3 K/T

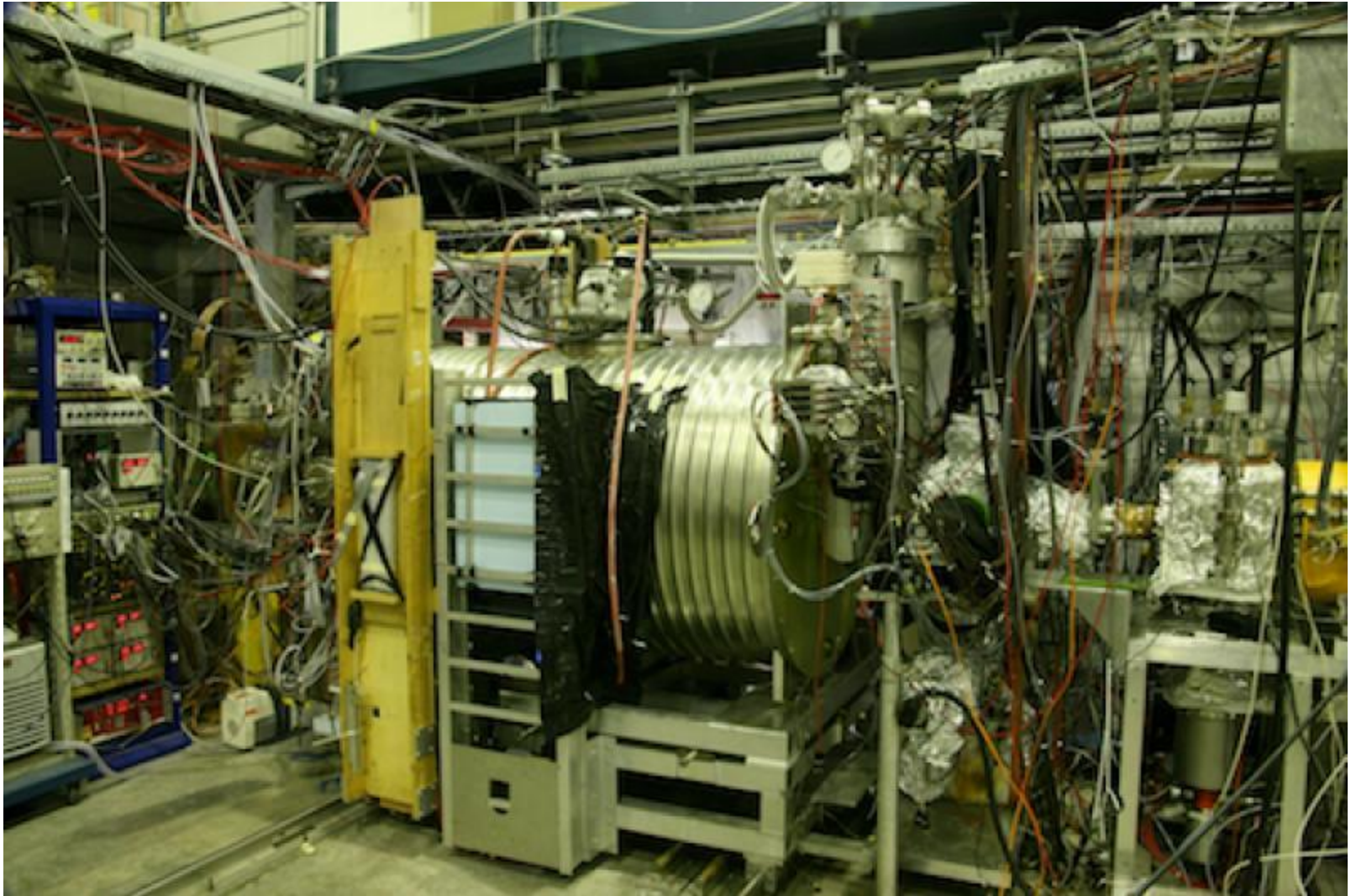
Actually 4 states: 2 trapped and 2 ejected

# AntiHydrogen Trap

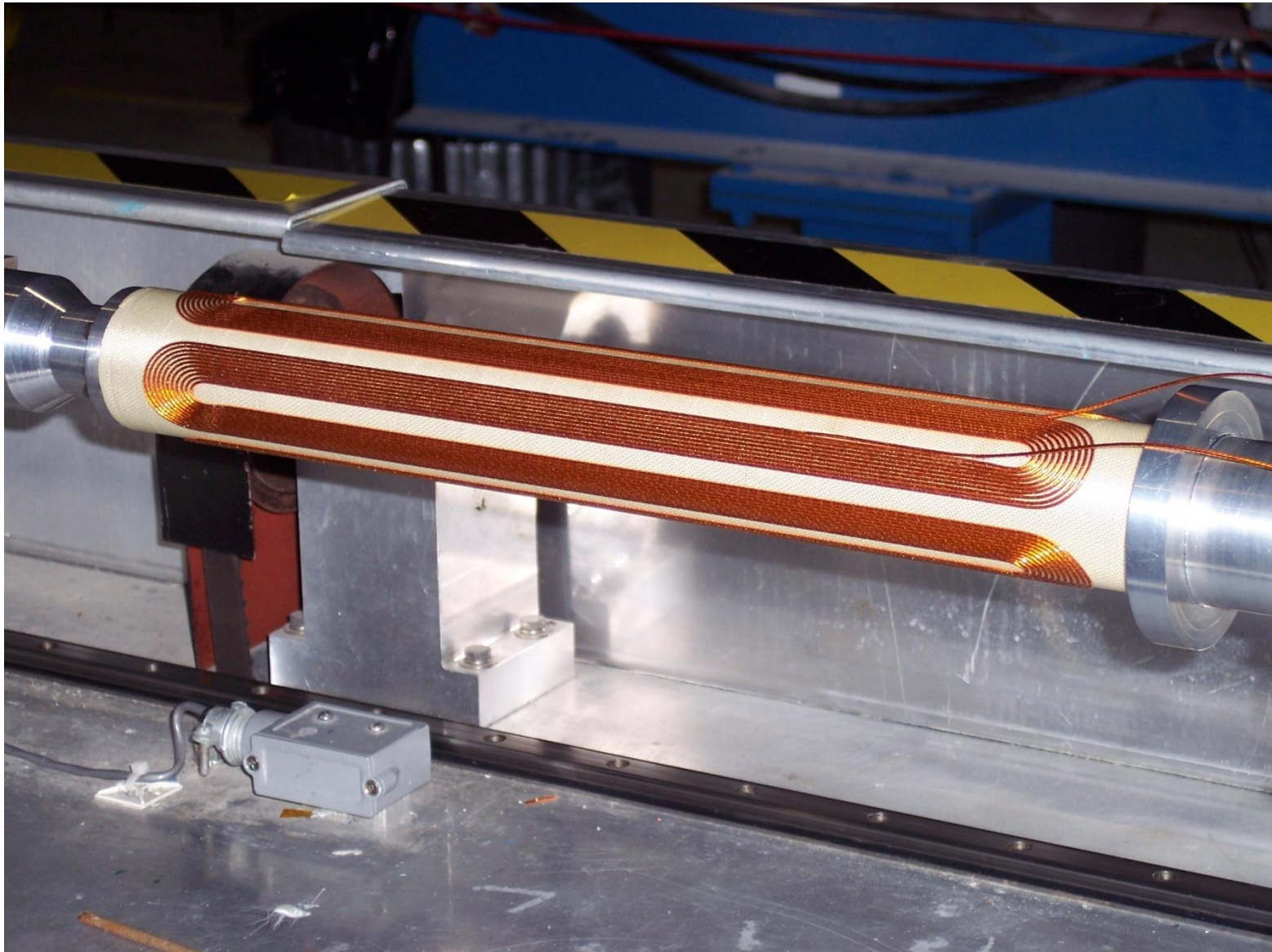




# External View



# Actual Magnets: Octupole





# Highly Excited (Rydberg) Atoms

Energy  $\sim 1/n^2$

Size  $\sim n^2$

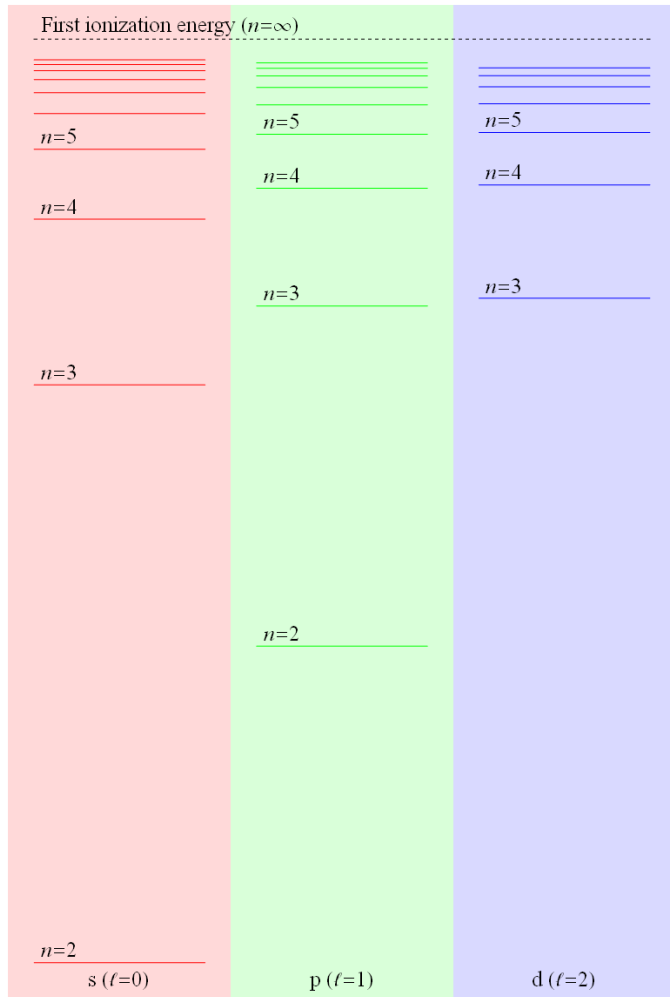
Degeneracy  $\sim n^2$

Degeneracy (cylindrical)  $\sim n$

Electric field  $\sim 1/n^4$

Magnetic field  $\sim 1/n$

Microwave ( $n=50$  to  $51$ )  $\sim 51$  GHz



Useful for sensing fields, long range interaction between pairs, strong interaction with a cold ground state atom, ..<sup>16</sup>



# Trilobite Molecules

Because so many states at nearly same energy, rearrange to get a lot of wave function at a spot

Can bind a cold atom to the Rydberg atom

<https://journals.aps.org/prl/abstract/10.1103/PhysRevLett.85.2458>

<https://www.nature.com/articles/nature07945>

<https://www.nature.com/articles/s41467-018-04135-6/>

