

Chapter 12 Nuclear Structure & Radioactivity

The type of nuclei is specified as ${}^A_Z X_N$ or A_X
 ${}^{14}_6 C_6$ or ${}^{14}_7 C$ ${}^4_2 He_2$ or ${}^4_3 He$ ${}^{235}_{92} U_{143}$ or ${}^{235}_{92} U$

The density of nucleons is roughly constant (Chap 12 slide 1)

$$\frac{A}{\frac{4\pi}{3} R^3} \approx \text{const} \rightarrow R \approx R_0 A^{1/3} \quad R_0 \approx 1.2 \times 10^{-15} \text{ m} \pm 0.2 \times 10^{-15} \text{ m}$$

Combining protons and neutrons into a nucleus can give off a lot of energy (Chap 12 slide 2)

$$M_p = 1.672622 \times 10^{-27} \text{ kg} \quad M_n = 1.674927 \times 10^{-27} \text{ kg}$$

$$M_p + M_n = 3.347549 \times 10^{-27} \text{ kg} \quad M_d = 3.343584 \times 10^{-27} \text{ kg}$$

$$(M_p + M_n - M_d) c^2 = 3.965 \times 10^{-30} \text{ kg} (3 \times 10^8 \text{ m/s})^2 = 3.57 \times 10^{-13} \text{ J} = 2.23 \text{ MeV}$$

This is the amount of energy you get when $p + n \rightarrow d$

In general, the binding energy of the nucleus

$$B = (N m_n + Z m_p - M_x) c^2$$

\uparrow nuclear mass

Sometimes, you have the mass of the atom (Table D)

To a good approximation

$$B \approx (N m_n + Z m({}_1 H_1) - M({}^A_Z X_N)) c^2$$

\uparrow hydrogen atom \uparrow atom mass

Nuclei with the same Z are isotopes

${}^{12}_6 C_6$ and ${}^{14}_6 C_8$
 ${}^1_1 H_1$ and ${}^3_2 He_1$,
 ${}^3_1 H_2$ and ${}^3_3 He_1$

" " " " N " isotones

" " " " A " isobars

Isotope interesting because chemistry mostly from Z
 Isobar interesting because some nuclear decay $p \rightarrow n$ or $n \rightarrow p$
 keeps A the same

The calculation of binding energy from first principles is very difficult but can be done for small nuclei. However, a phenomenologically motivated formula gives pretty decent estimate. (Chap 12 slides 3-6)

$$B = \frac{a_v A}{\text{Volume}} - \frac{a_s Z^{2/3}}{\text{Surface}} - \frac{a_c Z^2}{A^{1/3}} - \frac{a_A (N-Z)^2}{A} + a_p \quad (\text{Chap 12 slide 7})$$

Volume each nucleon has ~same nearby nucleons, each contribute

Surface nuclei at surface are missing neighbors

Coulomb from electric repulsion between nuclei

Asymmetry from Pauli exclusion princ.

Pairing nucleons have energy cost when going from even \rightarrow odd

(see Table of calculated binding energies on web page + chap 12 slide 8)

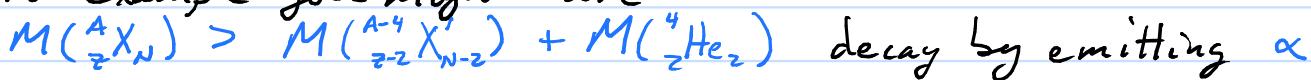
Can find the Z that gives the strongest binding for a fixed A

$$A \Rightarrow \frac{\partial B}{\partial Z} = 0$$

$$Z = A / (Z + \frac{a_c}{2 a_A} A^{2/3}) = A / (Z + 0.015 A^{2/3})$$

Just because $B > 0$ does not mean that nucleus is stable.

For example you might have



Could also emit a proton or neutron.

Emitting a proton or α , need to tunnel!



Lower energy α tunnel farther
 \Rightarrow longer half-life

$Q = [M(^A_X) - M(^{A-4}X') - M(^4He)]c^2$ is the energy released in the decay.

Decay because of the weak force

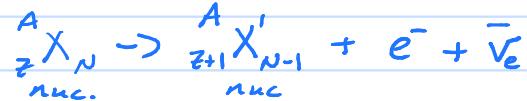
$$n \rightarrow p + e^- + \bar{\nu}_e$$

very little mass

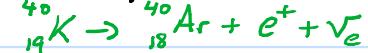
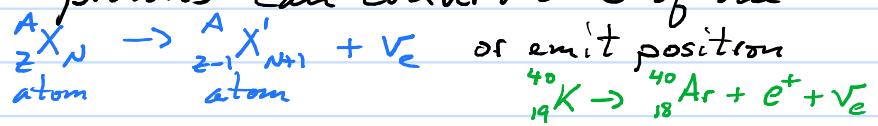
Return to neutrino below

$$Q = (m_n - m_p - m_e) c^2$$

For nuclei overly rich in neutrons



" " " " protons can convert one of the electrons and protons



Nuclei that decay behaves like spontaneous photon emission from atoms. All quantum systems with continuous final energy decay the same way: exponential decrease in population and Lorentzian distribution of final energies.

$N(t)$ = # of radioactive nuclei at time t

$$\frac{dN(t)}{dt} = -\lambda N(t) \Rightarrow N(t) = N(0) e^{-\lambda t} = N_0 e^{-\lambda t}$$

$$\text{Activity } a(t) = \lambda N(t) = a_0 e^{-\lambda t}$$

$$\lambda = \text{decay constant} = \frac{1}{(\text{mean lifetime})} = \frac{\ln(2)}{t_{1/2}} = \frac{0.693}{t_{1/2}}$$

(Chap 12 slides 9-11)

Radioactivity can be used to date various rocks, fossils, ...
 Basic idea: 1) (example Uranium-Lead dating) know that it starts with 100% ${}^A X$ and 0% ${}^A X'$ but radioactivity gives $(1-f) {}^A X$ and $f {}^A X'$
 2) (example ${}^{14}\text{C}$ dating) method to estimate initial fractions

Some basics of nuclear energy production (Chap 12 slides 12-15)