

Radiative Magnetic Reconnection in Astrophysical Plasmas

Dmitri Uzdensky

(University of Colorado, Boulder)

collaborators:

- B. Cerutti*, G. Werner, K. Nalewajko, M. Begelman (Univ. Colorado)
- A. Spitkovsky (Princeton)
- J. McKinney (Univ. Maryland)

*now @ Princeton

Workshop on Relativistic Plasma Astrophysics
Purdue, March 13, 2014

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OUTLINE

- **Introduction:** *Radiative* magnetic reconnection in astrophysics.
- **Astrophysical examples** of radiative reconnection:
 - Strong **synchrotron cooling** in **pulsar magnetosphere** reconnection.
 - Relativistic pair plasma reconnection and **Crab Nebula γ -flares**.
 - Other examples: Magnetar Flares; Blazar TeV flares; GRBs; black-hole accretion-disk coronae, etc.
- **Sweet-Parker** model for resistive reconnection with strong radiative cooling.
- **Summary**

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- ***News Flash!! New Results on Particle Acceleration in Relativistic Pair Reconnection!***
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Particle Acceleration in Relativistic Pair Reconnection (*Werner et al. 2014*)

- 2D simulations using Vorpal (and Zeltron) PIC codes.
- Simplest setup: no radiation, no guide field, $L_x = L_y = L$.
- Focus on *relativistic* ($\sigma^{\text{upstream}} \gg 1$), *large-system* ($L/\rho_0 \gg \sigma \gg 1$) regime ($\rho_0 = m_e c^2 / eB_0$).

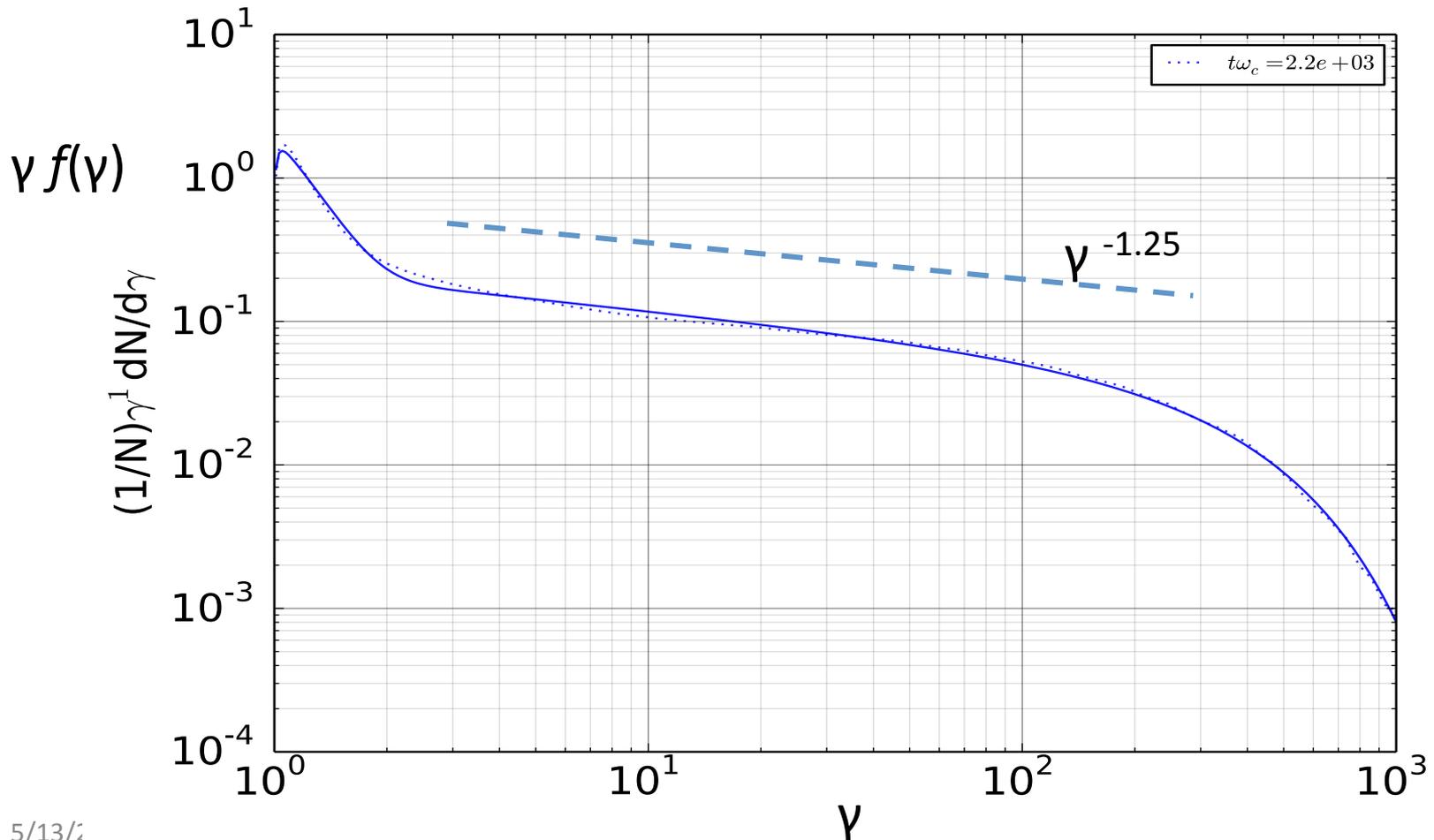
General Goal:

- completely characterize the resulting final particle energy distribution function in terms of L and σ .
- Is there a nonthermal power law tail and to what energies does it extend (how far beyond $\langle \gamma \rangle$)? What determines the high-energy cutoff?

Particle Acceleration in Relativistic Pair Reconnection (*Werner et al. 2014*)

Main Findings:

$$\sigma=100; L/\rho_0=100 \sigma$$



Particle Acceleration in Relativistic Pair Reconnection *(Werner et al. 2014)*

Main Findings:

- (Low-energy) power law:

$$f(\gamma) \sim \gamma^{-\alpha}$$

$$\alpha \rightarrow \text{const} \approx 1.3 \pm 0.2$$

for large enough L and $\sigma > 10$.

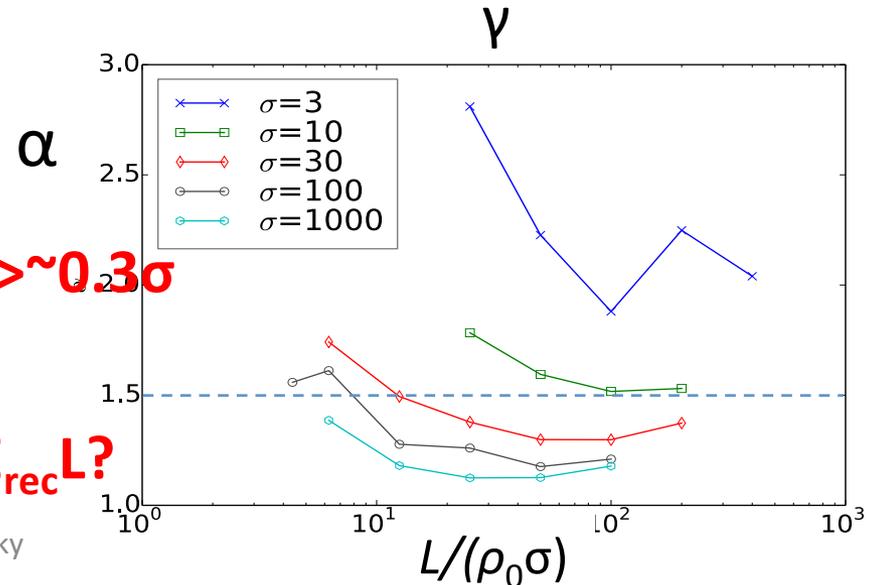
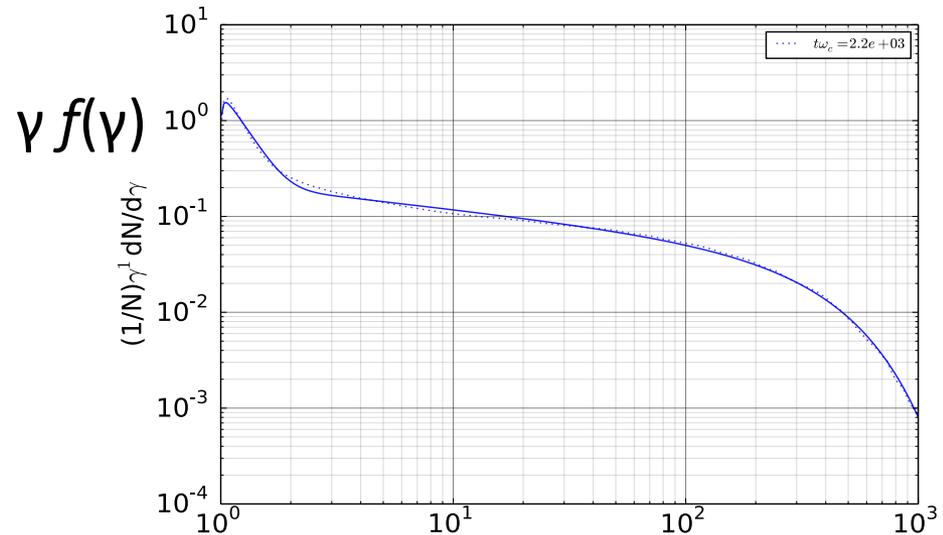
- Two high-energy cutoffs:

- $\exp(-\gamma/\gamma_{c1})$; $\gamma_{c1} \sim 3\sigma$,
weak L -dependence;

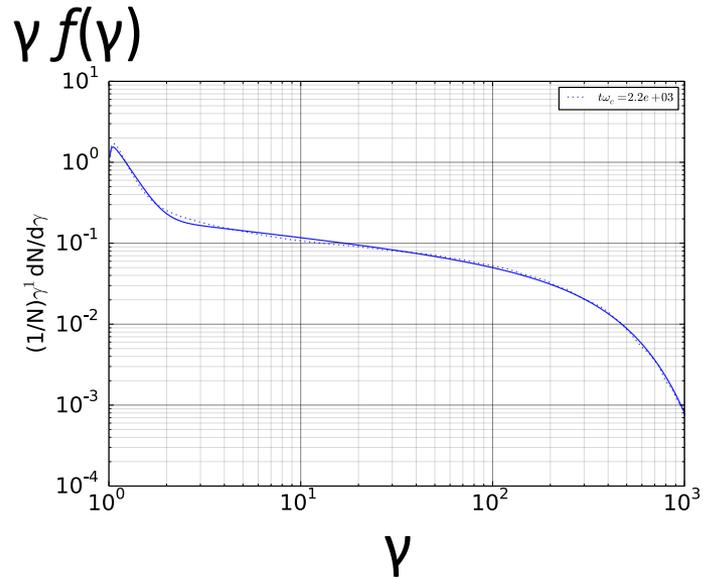
- $\exp(-\gamma/\gamma_{c2})^2$; $\gamma_{c2} \sim L$;
weak σ -dependence

$\gamma_{\text{max}} \approx (L/\rho_0)/2$ "extreme" accel.: $\epsilon_{\text{max}} = eE_{\text{rec}}L?$

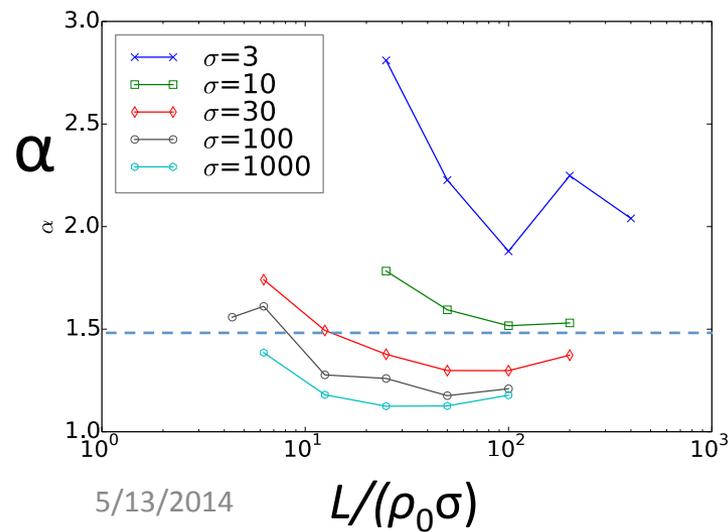
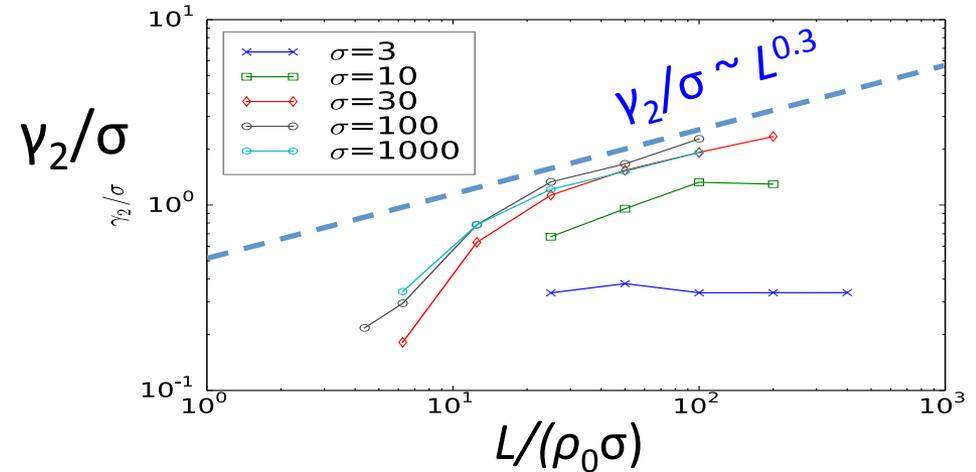
heating: $\epsilon_{\text{max}} \approx \langle \epsilon \rangle \sim 0.3\sigma$



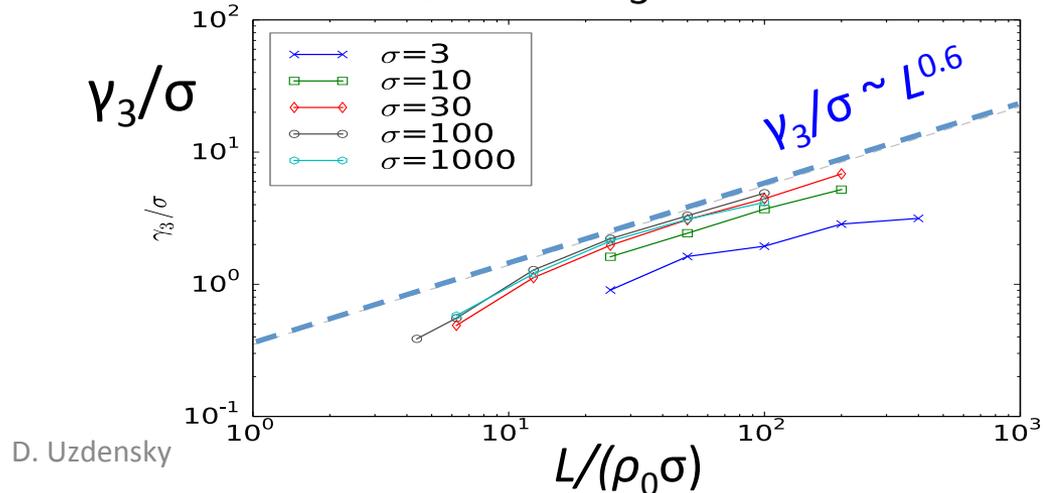
Particle Acceleration in Relativistic Pair Reconnection *(Werner et al. 2014)*



Energy-dominating scale γ_2 : max of $\gamma^2 f(\gamma)$



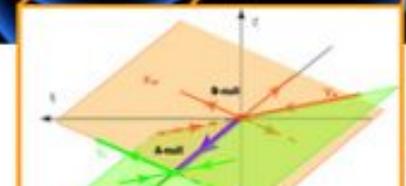
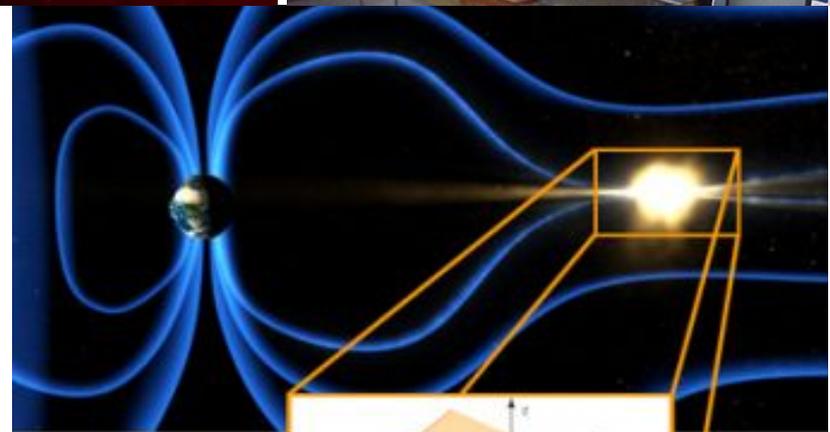
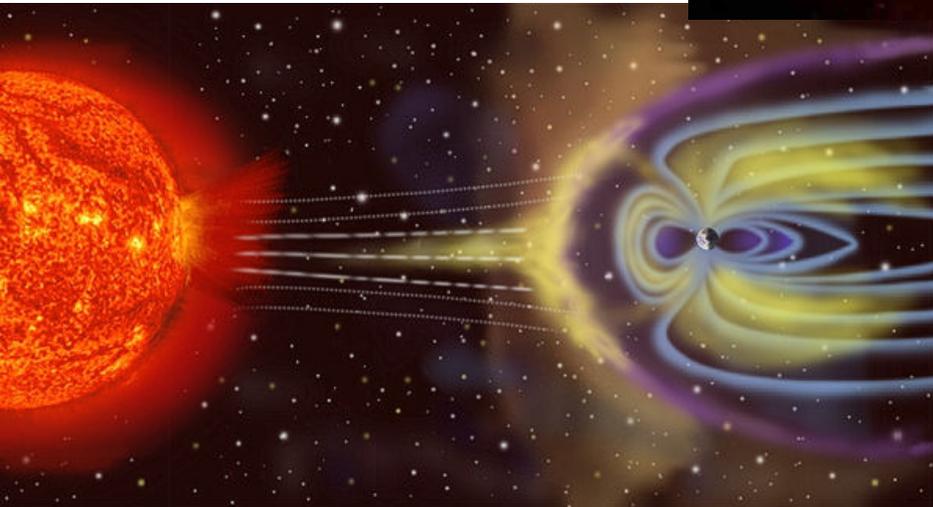
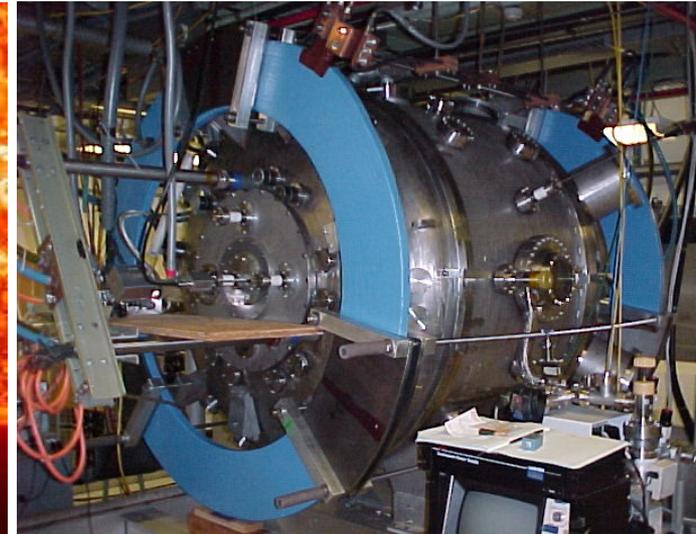
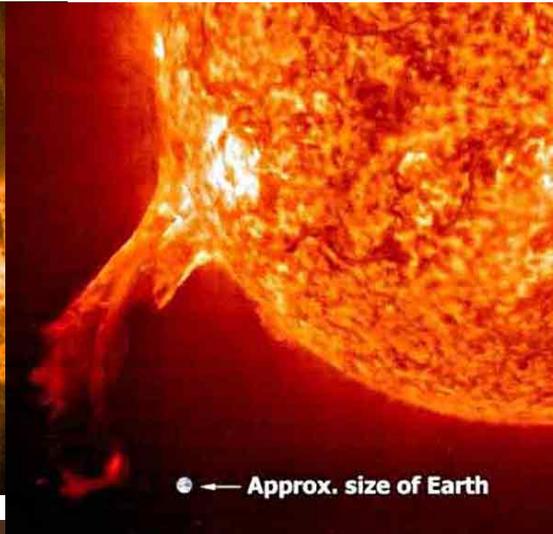
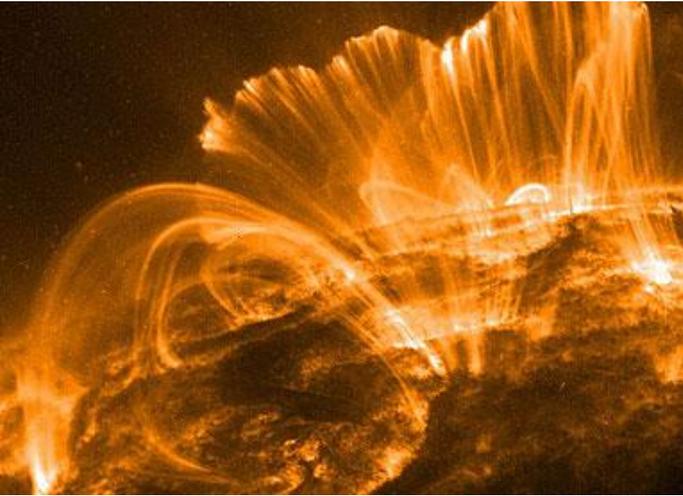
Radiation-dominating scale γ_3 : max of $\gamma^3 f(\gamma)$



OUTLINE

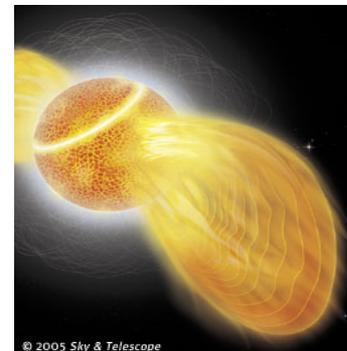
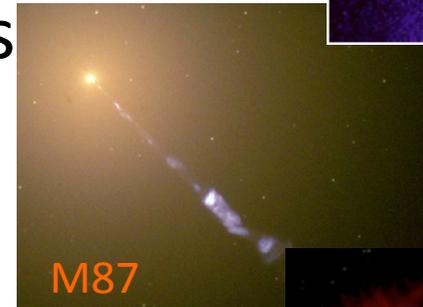
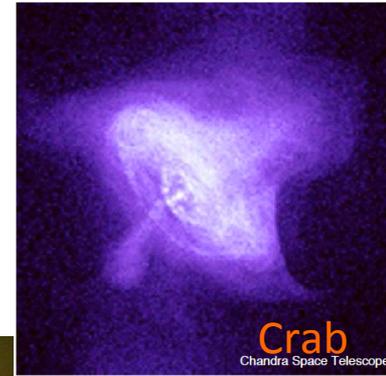
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Traditional Magnetic Reconnection in the Solar System



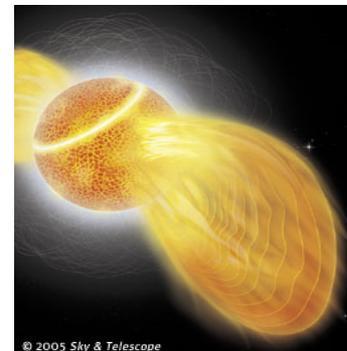
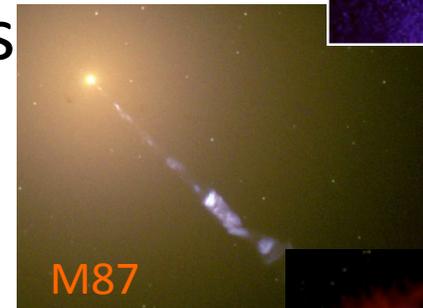
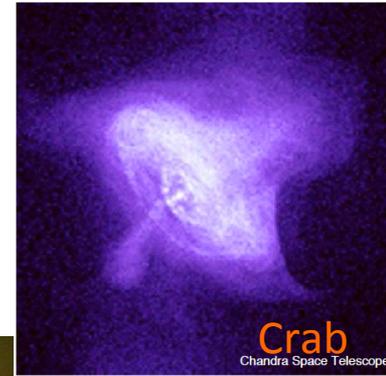
Reconnection in Astrophysics

- Pulsar magnetospheres, winds, PWNe
- AGN (e.g., blazar) jets, radio-lobes
- Gamma-Ray Bursts (GRBs)
- Magnetar (SGR) flares



Radiative Reconnection in Astrophysics

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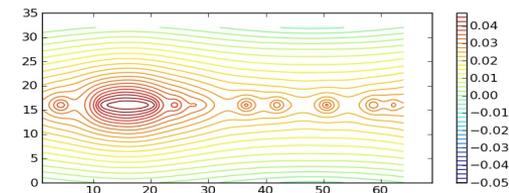
Radiation in Astrophysical Reconnection

- In conventional reconnection studies (space/solar/lab), plasma consists of electrons and ions --- **no photons!**
- In contrast, in many high-energy *astrophysical* situations **radiation** is important --- strongly affects reconnection:
 - Radiative cooling;
 - Radiative drag on reconnection outflow;
 - Radiation pressure;
 - Compton-drag resistivity.
- In addition, radiation is our only ***observational diagnostic*** into astrophysical reconnection.

How does a reconnection layer look like, literally?

- what are (prompt) **radiative signatures** (spectra, light curves) of reconnection seen by an outside observer

Radiative magnetic reconnection is a new frontier in plasma astrophysics!!



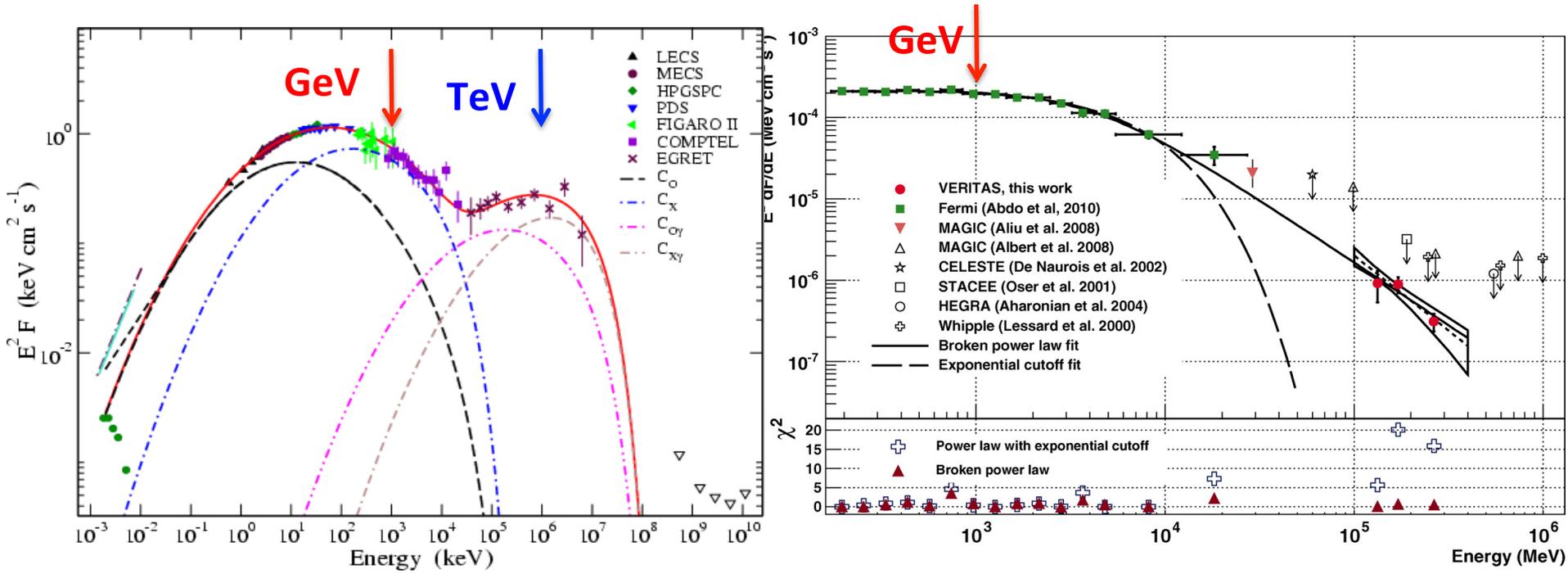
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Astrophysical Example of Reconnection with Strong Radiative Cooling: Pulsar Magnetosphere

(Uzdensky & Spitkovsky ApJ 2014)

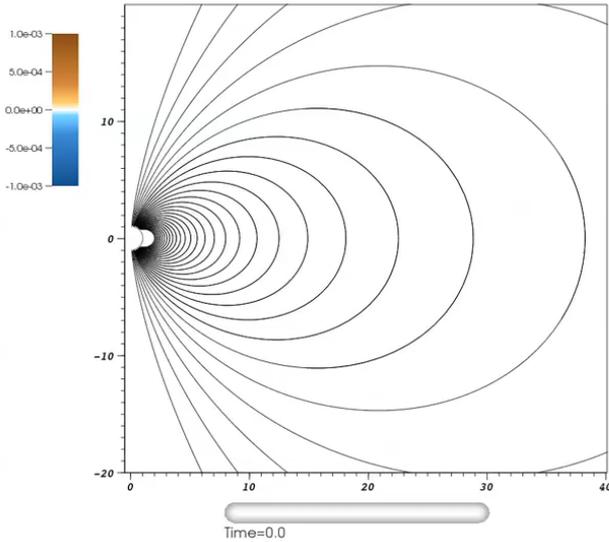
CRAB PULSED SPECTRUM



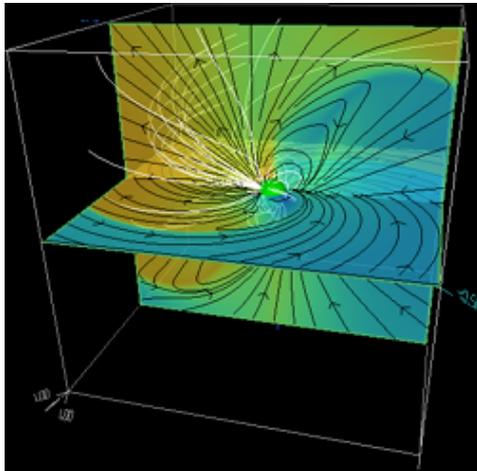
Detailed modeling of pulsed emission is affected by realistic geometry (oblique rotator), relativistic kinematics, etc. (Bai & Spitkovsky 2010).

But what are the basic plasma parameters in the emitting region?

Pulsar Magnetosphere: General Structure



Parfrey & Beloborodov 2012



Spitkovsky 2006

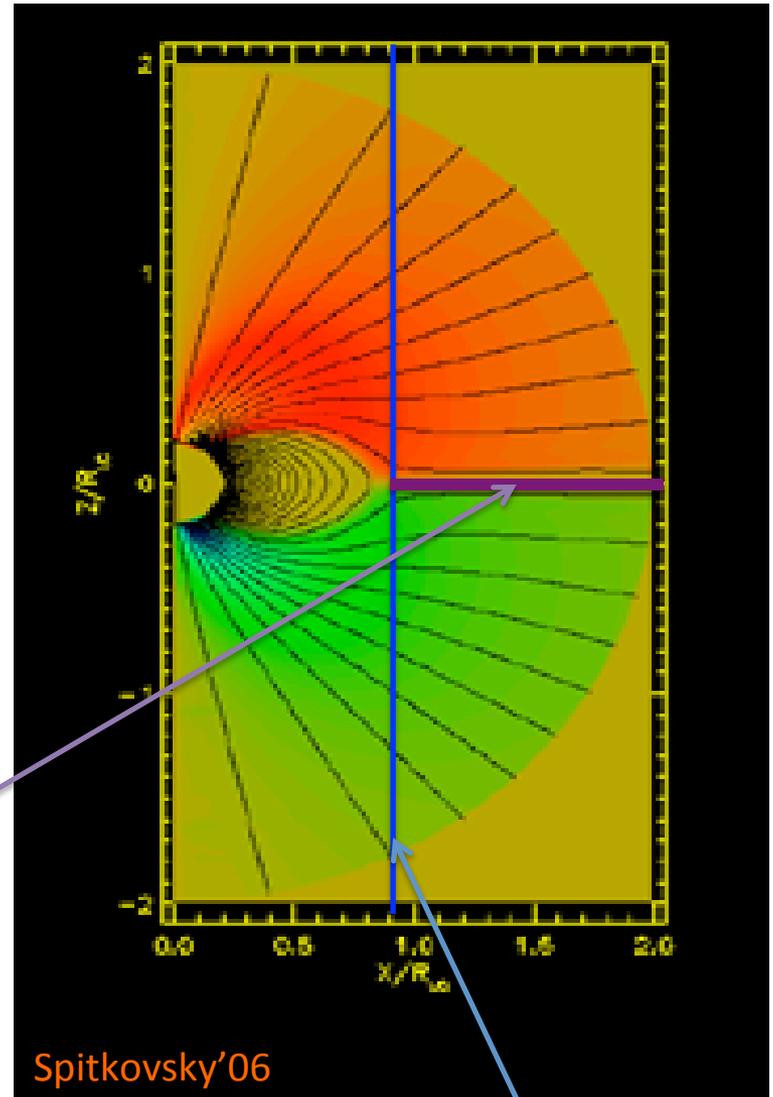
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equatorial
current sheet

$$R_{NS} \approx 10 \text{ km}$$

$$R_{LC} = c/\Omega \approx 1000 \text{ km}$$

D. Uzdensky

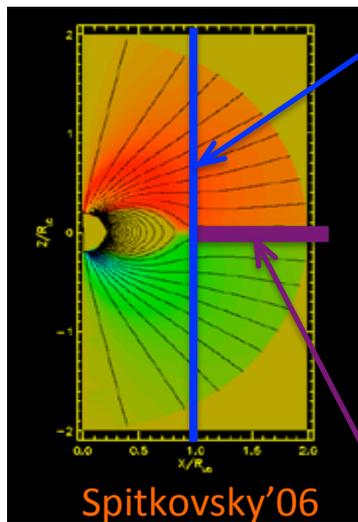


Spitkovsky'06

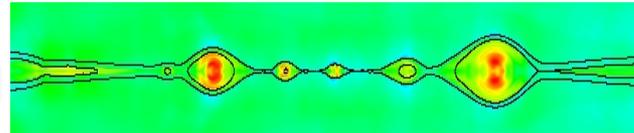
Light Cylinder

Reconnecting Pulsar Magnetosphere

- Equatorial current sheet should be tearing-unstable (*Lyubarsky' 96*), leading to a hierarchical chain of (merging) plasmoids/flux ropes.

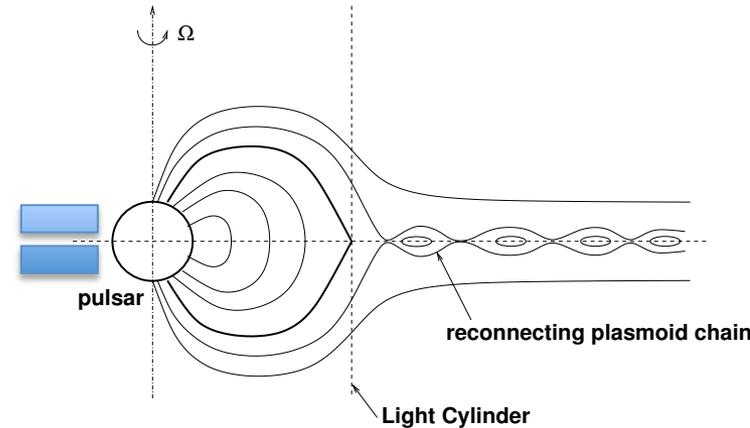


Light Cylinder



(*Loureiro et al. '07, '12*)

equatorial
current sheet



(*Bucciantini et al. 2006;*
Contopoulos & Spitkovsky 2007)

Questions:

- What are the basic plasma conditions inside inter-plasmoid current sheets (actual sites of energy dissipation) ?
- What are the observational consequences ?

Reconnection with strong synchrotron cooling near pulsar Light Cylinder

- What are T , n , and δ in pulsar-wind comoving frame ($\Gamma_W \approx 100$)?

- Three equations:

- **Pressure Balance:** $B_0^2 / 8\pi = 2nT$

- Energy balance with **strong synchrotron cooling** (c.f. Lyubarsky' 96):

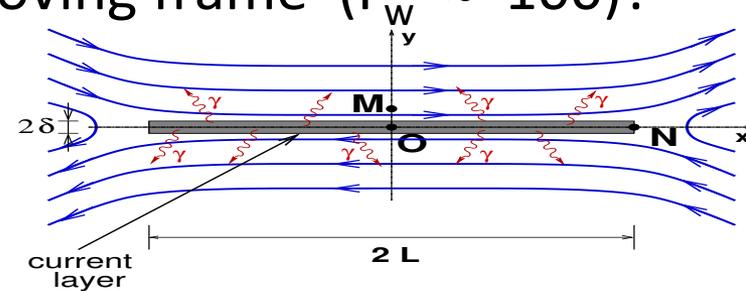
$$S_{\text{in}} = \frac{c}{4\pi} EB_0 \sim 2n \Lambda(B, \gamma_T) \delta \quad \text{where} \quad \Lambda(B, \gamma) = 2\sigma_T c \frac{B_{\perp}^2}{8\pi} \gamma^2$$

- **Ampere's Law:** $j_z = 2ne v_{\text{dr}} = 2nec \beta_{\text{dr}} \sim (c/4\pi) B_0 / \delta$

- Dimensionless parameters:

- dimensionless rec. rate: $\beta_{\text{rec}} = E/B_0 \sim 0.1$ (PIC simulations)

- e^+/e^- drift velocity: $\beta_{\text{dr}} = v_{\text{dr}} / c \approx 1 \Leftrightarrow \delta \approx \rho$



RESULTS

Obtain 3 comoving plasma parameters (T, n, δ) in terms of B_0 (~ 4 MG):

- Temperature:

$$\gamma_T = T / m_e c^2 \sim (\beta_{\text{dr}} \beta_{\text{rec}})^{1/2} b^{-1/2} \sim 4 \times 10^4 (\beta_{\text{dr}} \beta_{\text{rec}})^{1/2}$$

-- at **synchrotron radiation-reaction limit**: $e E_{\text{rec}} c = \Lambda_{\text{synch}}(\gamma, B_0)$

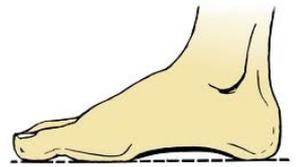
- Density:

$$n \sim (\beta_{\text{dr}} \beta_{\text{rec}})^{-1/2} (B_0^2 / 8\pi m_e c^2) b^{1/2} \sim 2 \times 10^{13} \text{ cm}^{-3} (\beta_{\text{dr}} \beta_{\text{rec}})^{-1/2}$$

-- rad. cooling \rightarrow strong plasma compression \gg ambient density
(Uzdensky & McKinney 2011)

- Layer thickness:

$$\delta \sim (\beta_{\text{rec}} / \beta_{\text{dr}})^{1/2} r_e b^{-3/2} \sim 30 \text{ cm} (\beta_{\text{rec}} / \beta_{\text{dr}})^{1/2}$$



normalized magnetic field: $b = B_0 / B_{\text{cl}} = r_e / \rho_c \ll 1$; $B_{\text{cl}} = e / r_e^2 \approx 6 \times 10^{15} \text{ G}$.

Astrophysical Implications (for Crab, etc.):

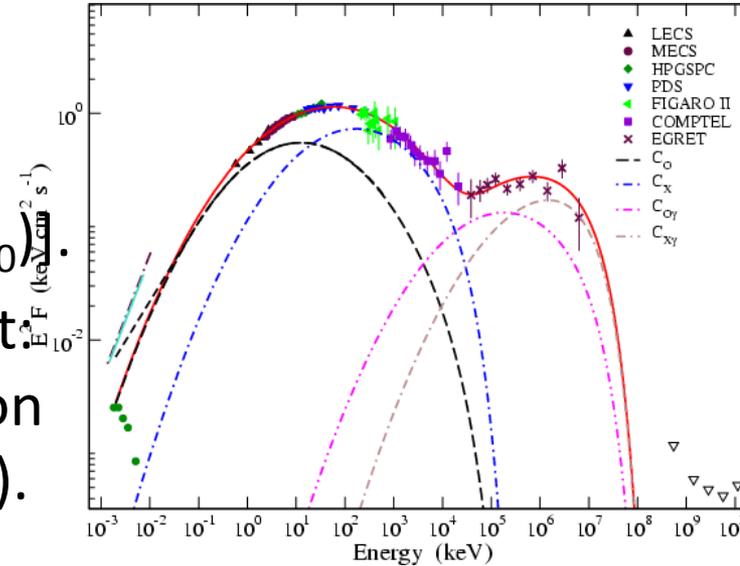
- Pulsed GeV FERMI γ -ray emission:

(*c.f. Lyubarsky '96, Petri '12, Arka & Dubus '13*)

Comoving temperature $T \approx 10$ GeV at the radiation-reaction limit $[e E_{\text{rec}} c = \Lambda_{\text{synch}}(\gamma, B_0)]$.

Synchrotron radiation at standard max. limit:

$\epsilon_{\text{ph}} \sim \beta_{\text{rec}} m_e c^2 / \alpha \approx 20$ MeV \rightarrow GeV emission after Doppler boost (due to rel. pulsar wind).



- Pulsed VHE (> 100 GeV) MAGIC/VERITAS γ -ray emission:

Inverse Compton scattering on hot pairs ($\epsilon = \Gamma T \approx 1$ TeV) in the layer.

- Pulsed radio emission:

dynamic reconnecting plasmoid chain with *cm*-scales...

SUMMARY

- (σ, L) parameter-space PIC study of relativistic pair reconnection:
 - power law $f(\gamma) \sim \gamma^{-\alpha}$ with $\alpha \approx 1.3$ and $[\exp(-\gamma/\gamma_{c1}) \times \exp(-\gamma/\gamma_{c2})^2]$ cutoff;
 - $\gamma_{c1} \sim \sigma L^{0.3} \leq 10 \langle \gamma \rangle$; $\gamma_{c2} \sim L$ – total voltage drop.
- Radiation is often important in high-energy astro reconnection!
- Radiation is our only direct diagnostic of astrophysical reconnection.
- Reconnection with strong synchrotron cooling in pulsar (e.g., Crab) magnetosphere outside LC:
 - Reconnection layer parameters (comoving T, n, δ) depend only on B_0 .
 - Can potentially explain FERMI observations of pulsed gamma-ray (GeV) emission for Crab and other pulsars.
- Sweet-Parker resistive MHD reconnection with strong rad. cooling:
 - heating/cooling balance determines layer temperature;
 - cooling leads to plasma compression and faster reconnection.