

Pulsar as plasma generator

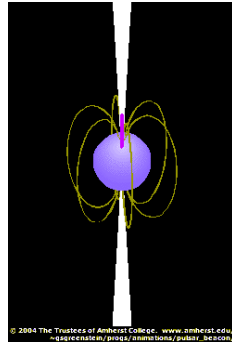
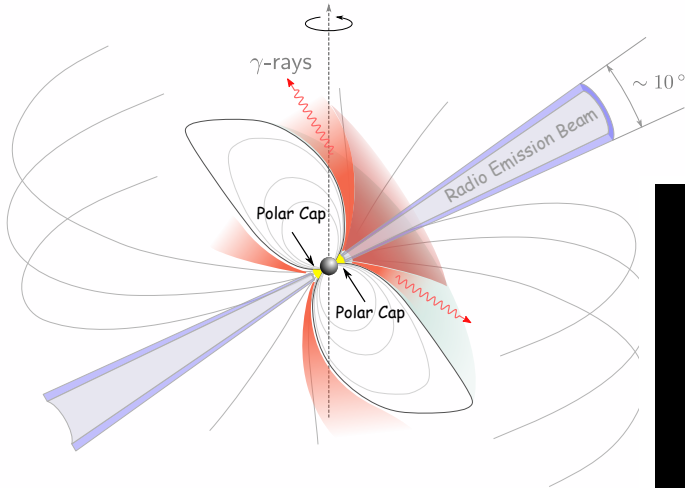
Andrey Timokhin

NASA Goddard Space Flight Center

Workshop on Relativistic Plasma Astrophysics,
Purdue University, May 11-15, 2014

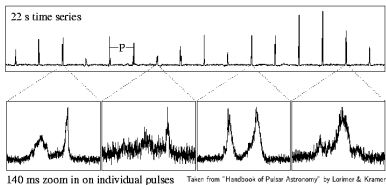
Pulsar: rapidly rotating magnetized neutron star

“Electric lighthouse”

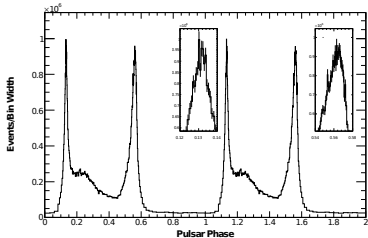


Pulsars: What we see

radio:



gamma:



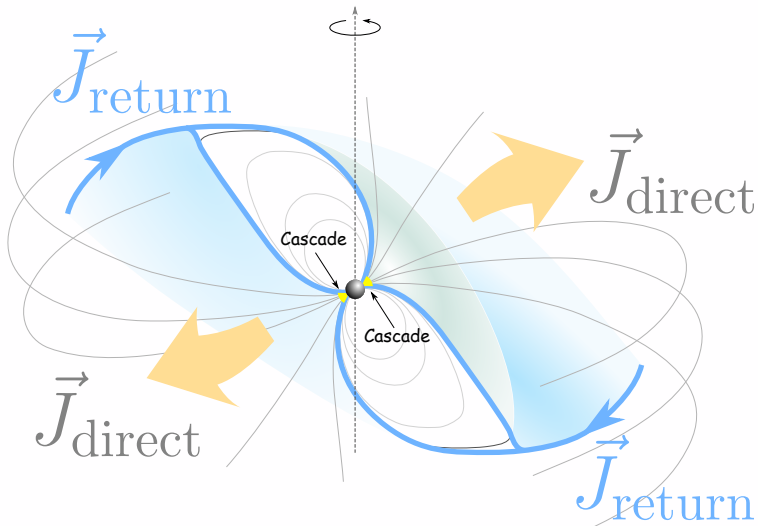
- Pulse peaks are narrow
- Negligible energy budget



- PWNe feed by dense plasma
- Energy goes there

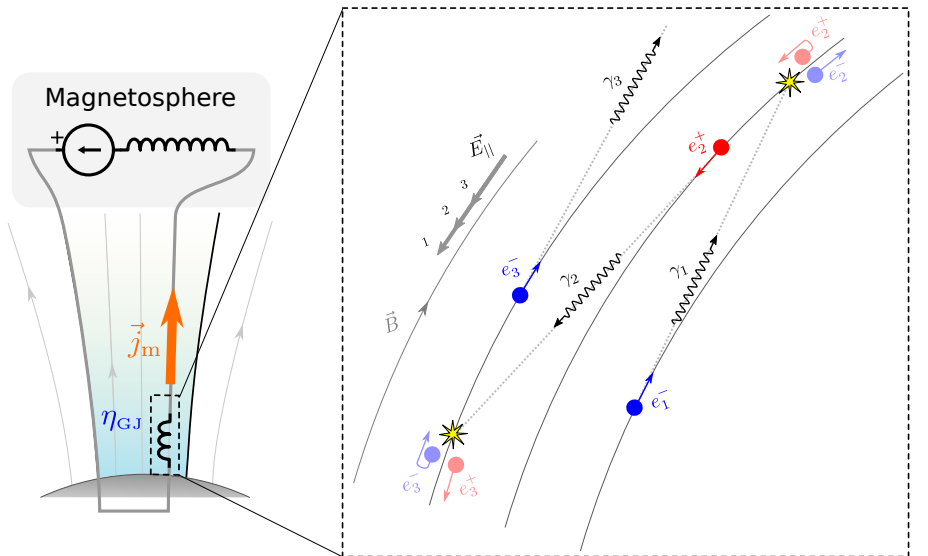
Pulsar Magnetosphere: Theorist view

The global circuit in the electrical generator



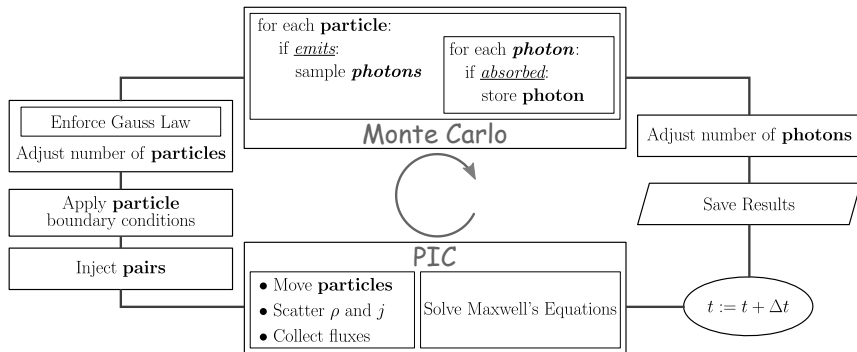
Plasma creation in the polar cap

Cascades are electromagnetically driven



Numerical code for cascade modeling *PAMINA*

PIC And *Monte-Carlo* code for cascades *IN* Astrophysics



[AT 2010, AT & Arons 2013]

Modeling from the first principles:

Particle acceleration $\leftrightarrow$ Electric field

Particles $\rightarrow$ Photons $\rightarrow$ Particles(Pairs)

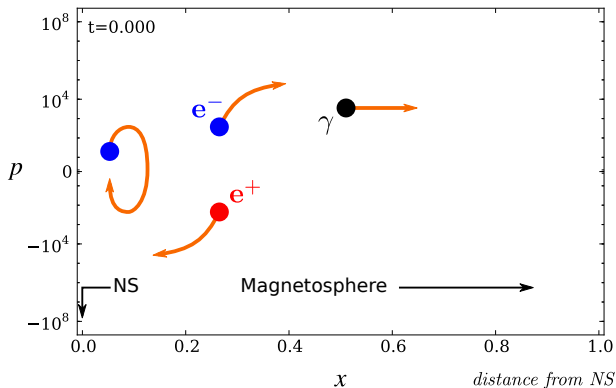
Particle-In-Cell

Monte Carlo

Limit cycle: series of discharges

No particles extraction from the NS

particles' momenta $p \equiv \frac{v}{c}\gamma$

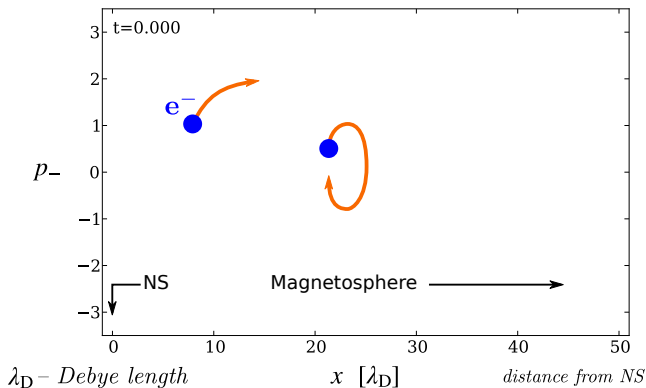


• electrons • positrons • γ -rays

Formation of a low energetic flow for $j/j_{\text{GJ}} < 1$

Free particle extraction from the NS

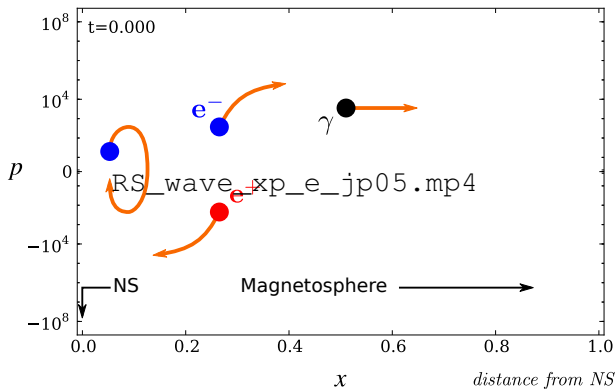
particles' momenta $p \equiv \frac{v}{c}\gamma$



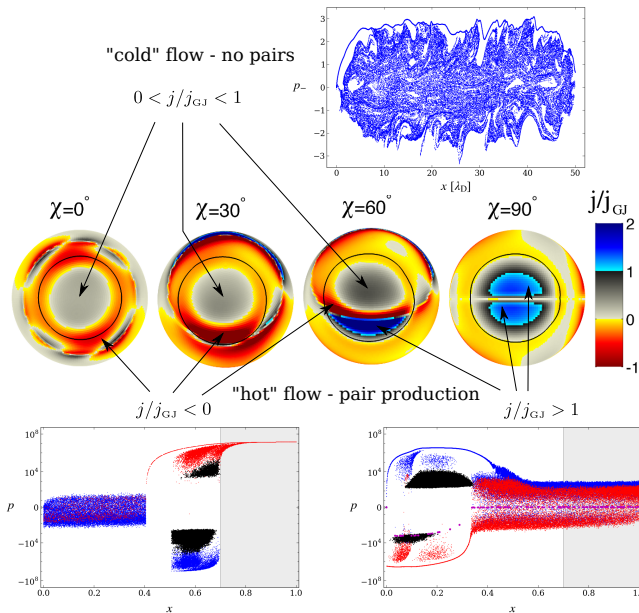
Waves during discharge

It did not escape our attention...

particles' momenta $p \equiv \frac{v}{c}\gamma$



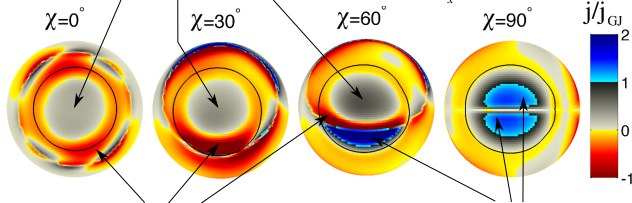
Free particle extraction from the NS



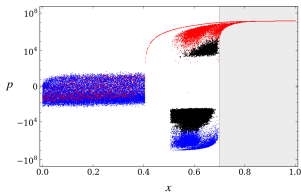
Particle acceleration with no particle extraction

"hot" flow everywhere,
pair production:

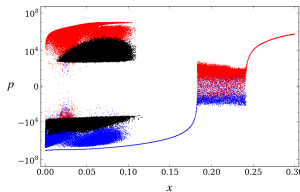
near NS $0 < j/j_{\text{GJ}} < 1$



far from NS $j/j_{\text{GJ}} < 0$



near NS $j/j_{\text{GJ}} > 1$



Particles'
momenta:
 $p \sim [0, 10^7]$

electrons
positrons
photons

Particles'
momenta:
 $p \sim [0, 10^7]$

electrons
positrons
photons

Electron-positron cascade

Cascade – splitting of primary particle energy into energy of pairs

$$\kappa_{\max} \simeq 2 \frac{\epsilon_p}{\epsilon_{\text{th}}} .$$

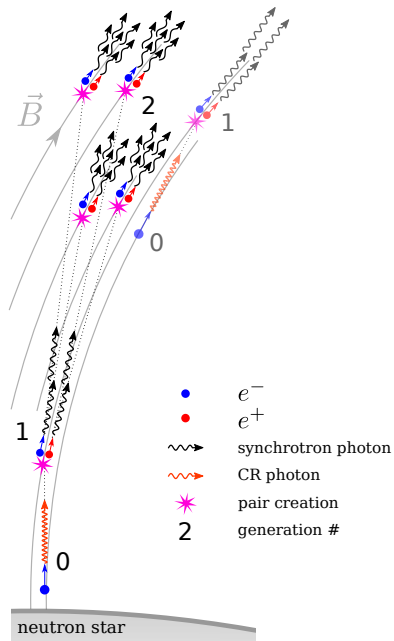
Cascade efficiency

- flux of primary particles
- energy of primary particles
- efficiency of splitting the energy

Full cascade

Synchrotron cascade

Curvature Radiation



Photon absorption in magnetic field

$$\alpha_B(\epsilon_\gamma, \psi) = 0.23 \frac{\alpha_f}{\lambda_c} \frac{B}{B_q} \sin \psi \exp \left(-\frac{4}{3\chi} \right)$$

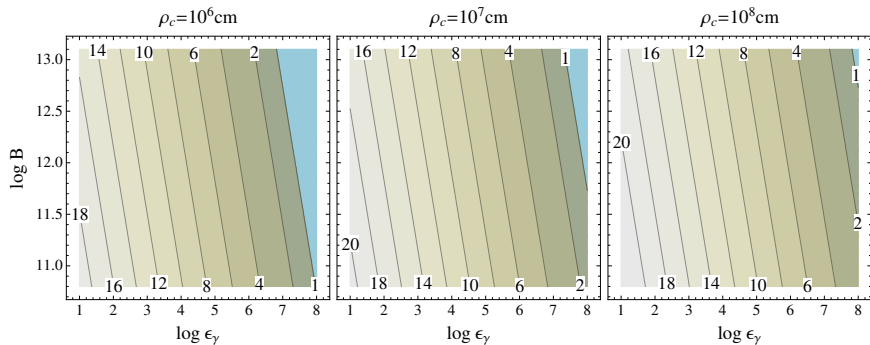
$$\chi \equiv \frac{1}{2} \epsilon_\gamma \frac{B}{B_q} \sin \psi ,$$

$$\tau(\epsilon_\gamma, l) = \int_0^l \alpha_B(\epsilon_\gamma, \psi(s)) ds ,$$

In the approximation of constant magnetic, $B \approx \text{const}$ ($r \ll R_{\text{NS}}$)

$$\tau(\chi) = A_B \frac{\rho_c}{\epsilon_\gamma^2 (B/B_q)} \left[\frac{1}{6} \chi (3\chi - 4) \exp \left(-\frac{4}{3\chi} \right) - \frac{8}{9} \text{Ei} \left(-\frac{4}{3\chi} \right) \right] .$$

$$1/\chi_a: \tau(\chi_a) = 1$$

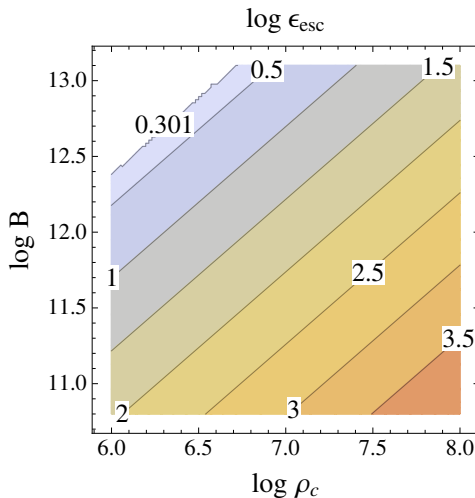


Energy of escaping photons

$$\rho_c \psi_a = s_{\text{esc}} R_{\text{NS}}$$

$$\epsilon_{\text{esc}} = \frac{2\rho_c}{s_{\text{esc}} R_{\text{NS}}} \frac{\chi_a(\epsilon_{\text{esc}}, B/B_q, \rho_c)}{B/B_q}$$

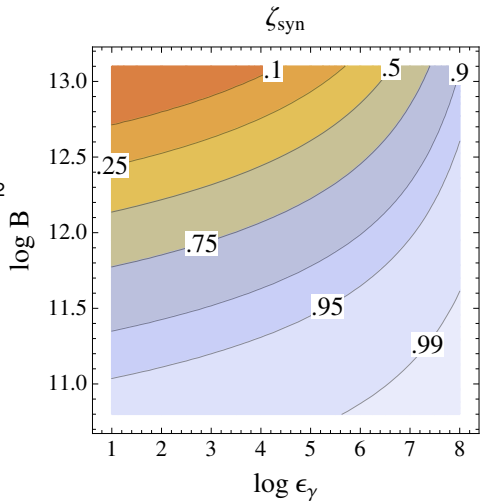
$s_{\text{esc}} \sim 1$



Fraction of photon energy going into synchrotron radiation: ζ_{syn}

$$\begin{aligned}\epsilon_{\pm, F} &= (1 - \beta_{\parallel}^2)^{-1/2} \\ &\approx \epsilon_{\pm, I} [1 + (\psi_a \epsilon_{\pm, I})^2]^{-1/2} \\ &= \epsilon_{\pm, I} \left[1 + \left(\frac{\chi_a}{B/B_q} \right)^2 \right]^{-1/2}\end{aligned}$$

$$\zeta_{\text{syn}} = \frac{2(\epsilon_{\pm, I} - \epsilon_{\pm, F})}{\epsilon_{\gamma}}$$



Effect of magnetic field on Synchrotron cascade

Decrease in multiplicity:

- high $B \rightarrow$ small photon MFP, lower ψ_a , more energy in pair kinetic energy
- low $B \rightarrow$ large escaping energy

There should be an optimum magnetic field: $B \sim 10^{12}$ G

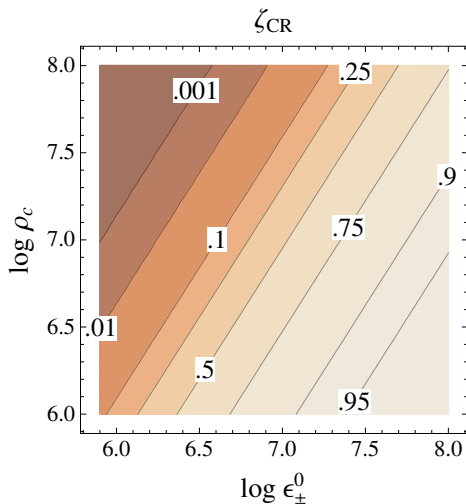
Fraction of the primary particle energy going into Curvature Radiation: ζ_{CR}

$$\frac{d\epsilon_{\pm}}{dt} = -\frac{2}{3} \frac{e^2}{m_e c^5} \left(\frac{c^2}{\rho_c} \right)^2 \epsilon_{\pm}^4$$

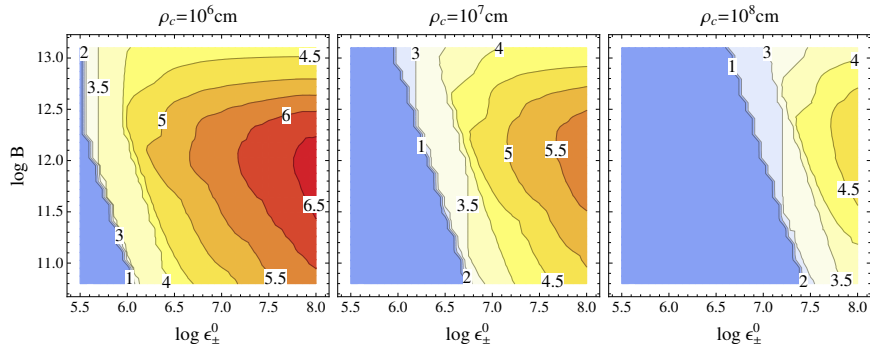
$$\epsilon_{\pm}(s) = \epsilon_{\pm}^0 \left[1 + 3H \frac{(\epsilon_{\pm}^0)^3}{\rho_c^2} s \right]^{-1/3}$$

$$\zeta_{\text{CR}}(s_{\text{CR}}) = 1 - \frac{\epsilon_{\pm}(s_{\text{CR}})}{\epsilon_{\pm}^0}$$

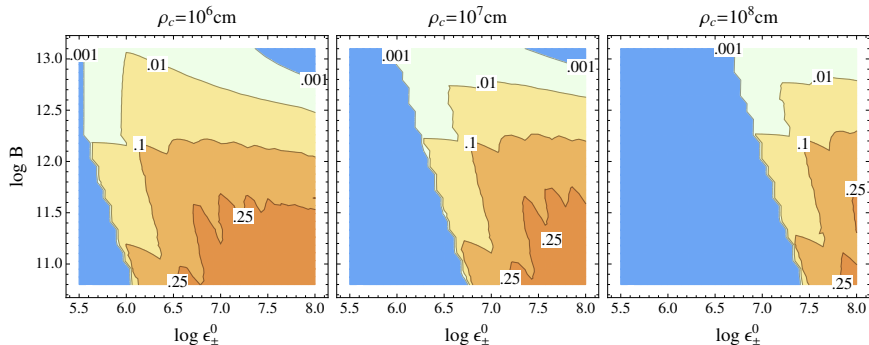
$$s_{\text{CR}} \sim 1$$



Multiplicity of CR-Synchrotron cascade: $\log \kappa$

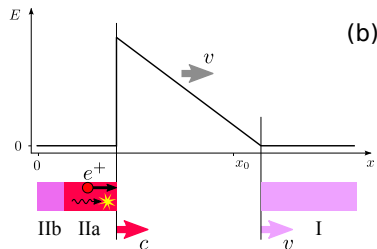
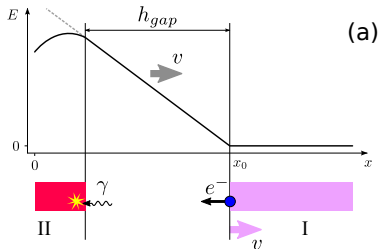


Efficiency of CR-Synchrotron cascade: $\kappa/\kappa_{\text{max}}$



$$\kappa_{\text{max}} \simeq 2 \frac{\epsilon_p}{\epsilon_{\text{esc}}} .$$

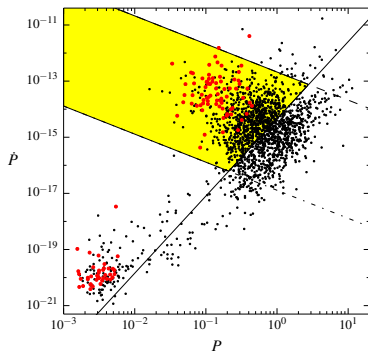
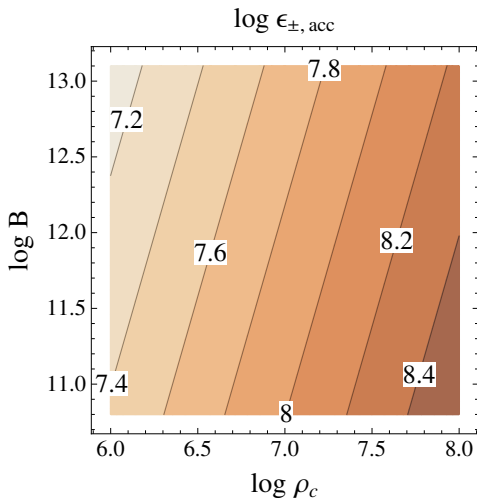
Particle acceleration



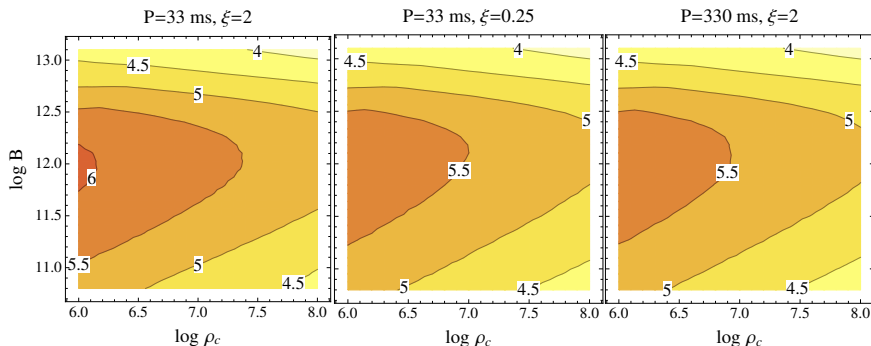
$$E(x_p) = 4\pi\eta_{\text{GJ}}^0 \xi_j (x_0 - x_p); \quad \xi_j \equiv \frac{j - j_m}{j_{\text{GJ}}^0} \left(1 + \frac{c}{v}\right);$$

$$\epsilon_{\pm, \text{acc}} \simeq 5 \times 10^7 \chi_a^{2/7} \xi_j^{1/7} \rho_{c,7}^{4/7} P^{-1/7} B_{12}^{-1/7}$$

Energy of primary particles



Multiplicity of polar cap cascade: $\kappa \sim 10^5$

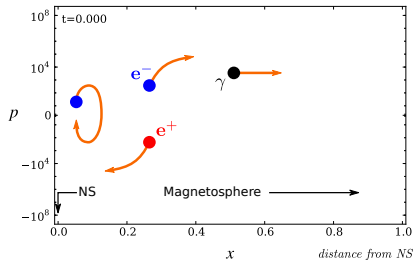


Dependence on ρ_c partially cancels out:

- small $\rho_c \rightarrow$ high splitting efficiently, but low primary particle energy
- large $\rho_c \rightarrow$ low splitting efficiently, but high primary particle energy

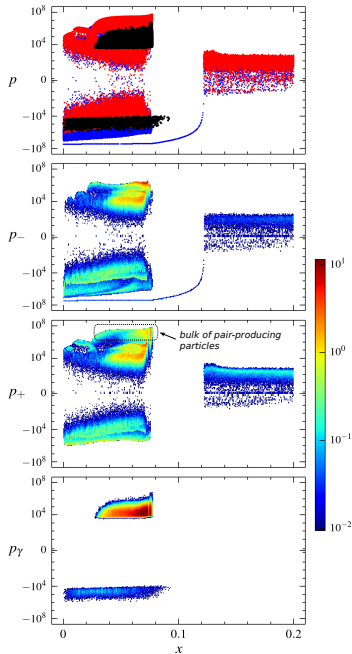
Discharge: RS flow

particles' momenta $p \equiv \frac{v}{c}\gamma$



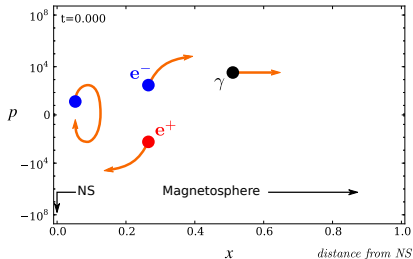
• electrons • positrons • γ -rays

- Low heating of NS surface
- Duty cycle: can be as low as $h_{\text{gap}}/R_{\text{NS}} \sim 1/100$ (for Crab)



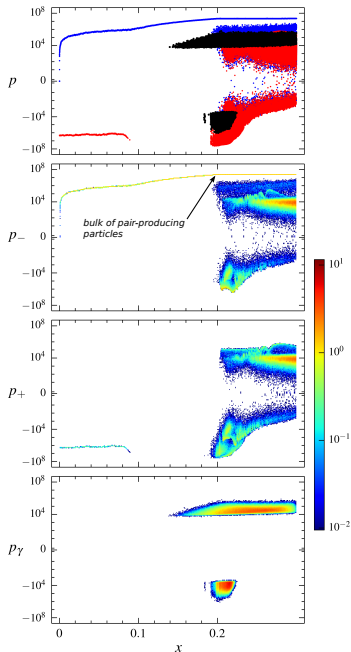
Discharge: super-GJ SCLF

particles' momenta $p \equiv \frac{v}{c}\gamma$

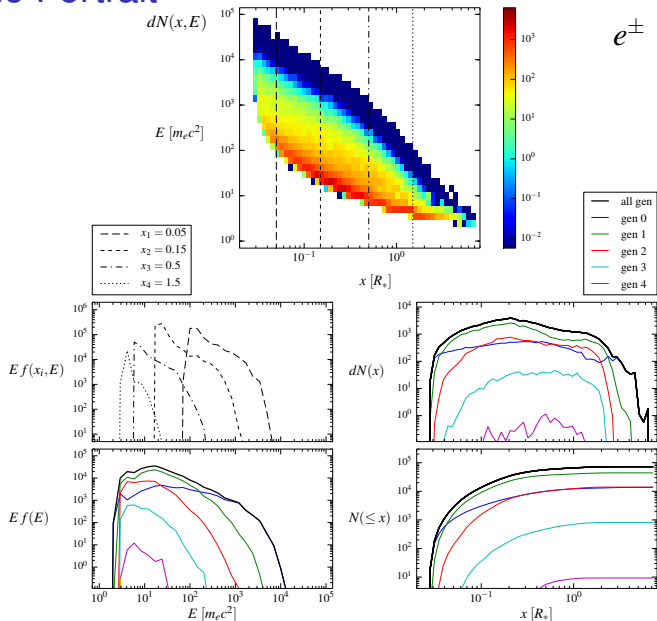


• electrons • positrons • γ -rays

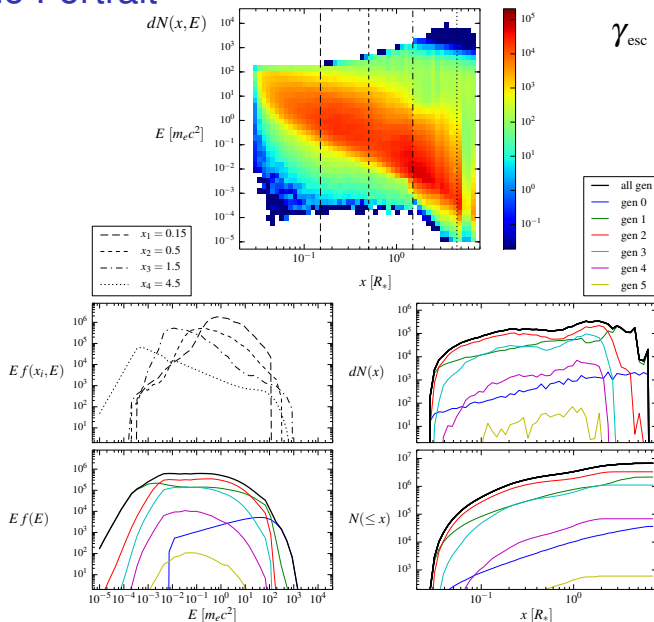
- Low heating of NS surface
- Duty cycle: $\sim 1/\text{few}$



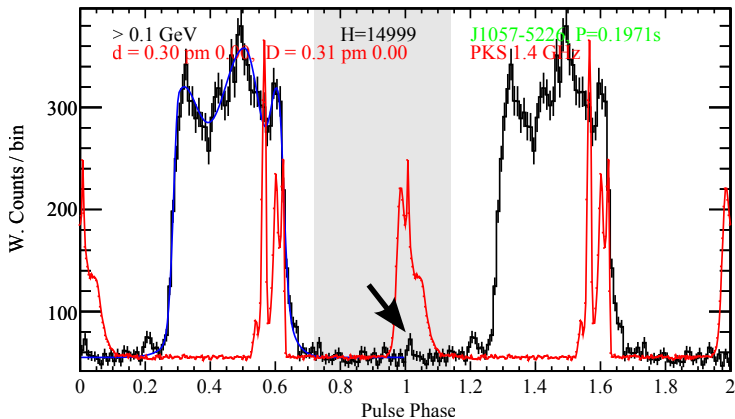
Cascade Portrait



Cascade Portrait



PSR J1057-5226: Polar cap emission?



Conclusions

- Particles can be accelerated faster and at lower altitudes
- Maximum multiplicity of polar cap cascades $\kappa \sim 10^5$
- Maximum multiplicity is not sensitive to pulsar parameters
- Plasma distribution is non-uniform
- Pair yield is mostly determined by primary particle flux:
super-GJ space charge limited flow zones should produce the highest pair yield
- Gamma-ray emission from polar caps is at lower energies
($\sim 10 - 100$ MeV)

Quest for Unicorns



Unicorns DO EXISTS!

