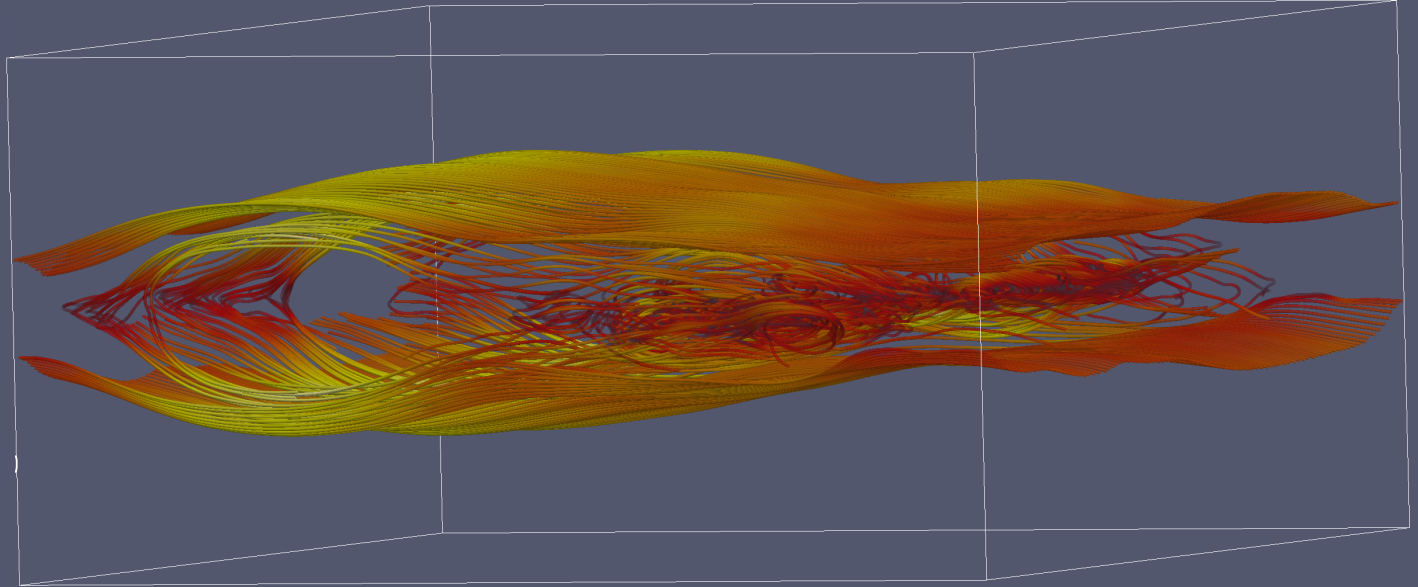


# Relativistic Reconnection: radiative and 3D effects

Relativistic  
reconnection  
in pair plasma

Using PIC codes:  
Zeltron  
Tristan-MP



Greg Werner  
Dmitri Uzdensky  
Vladimir Zhdankin  
Mitch Begelman  
(CU Boulder)

Sasha Philippov  
(UC Berkeley)

How are

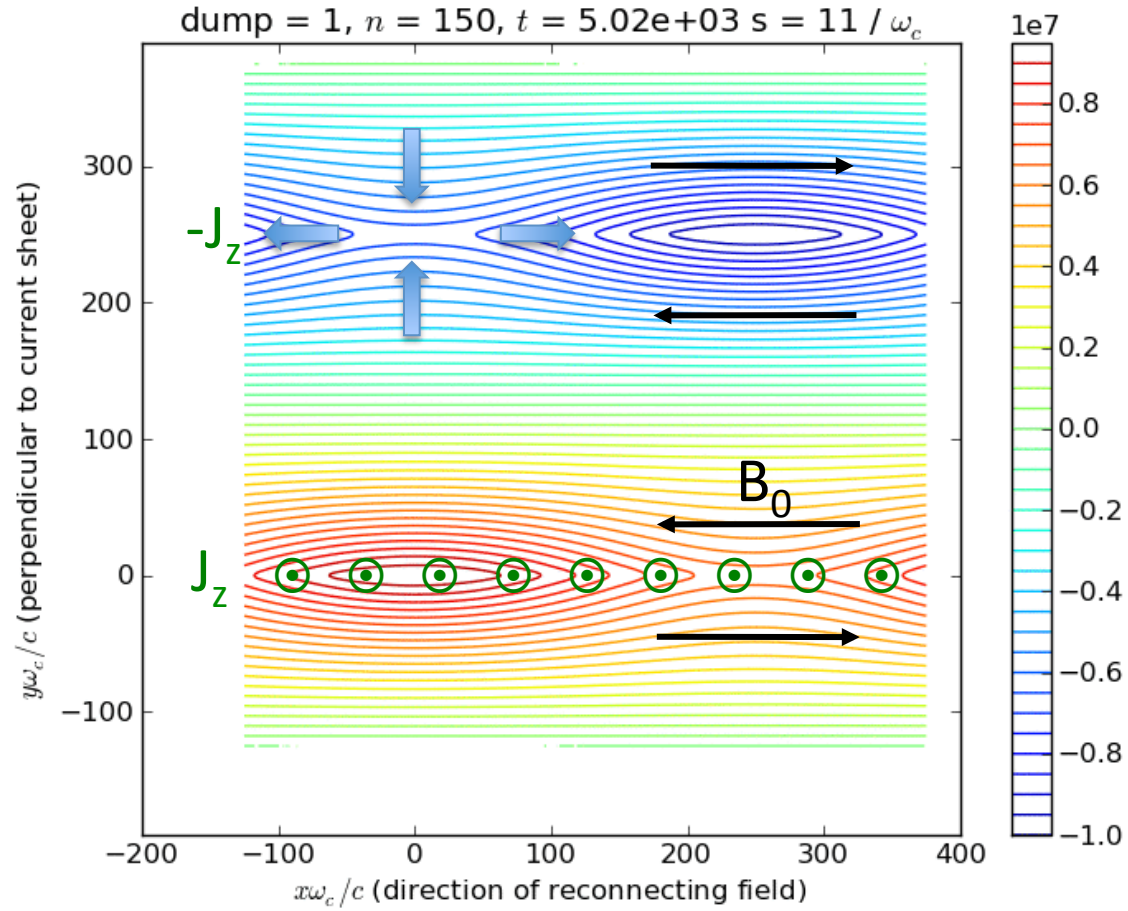
- reconnection dynamics/energetics
- and NTPA

affected by:

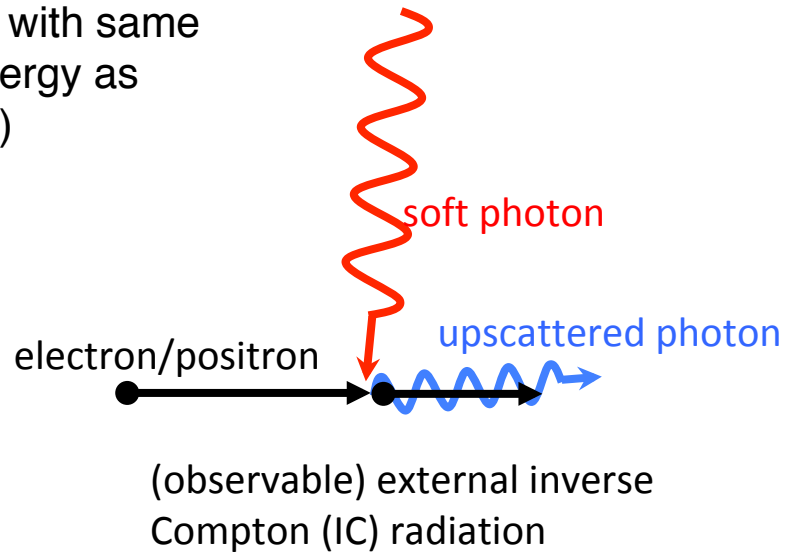
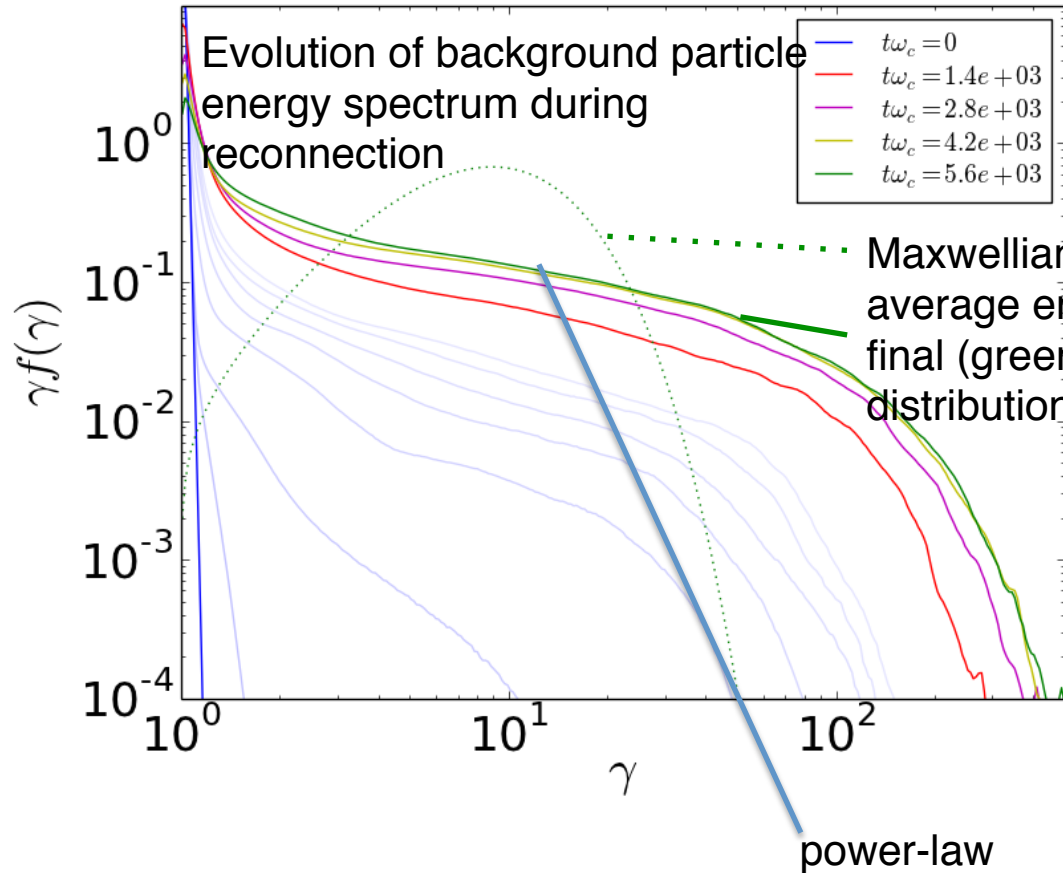
- length in the 3<sup>rd</sup> dimension (z); i.e., 3D-ness
- guide magnetic field
- external inverse Compton cooling
- ?

Bonus: Code Comparison

# Reconnection's main job: magnetic field energy $\rightarrow$ particle/plasma energy

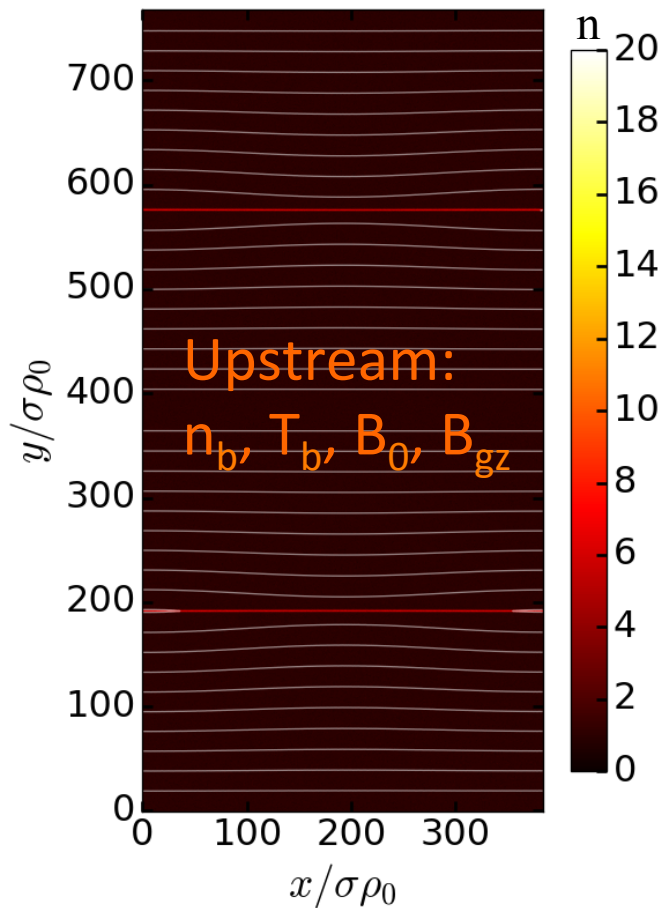


Reconnection  $\rightarrow$  particle energization/NTPA  $\rightarrow$  radiation  
 NTPA=Nonthermal Particle Acceleration  
 /  
 observational diagnostic



(The Lorentz factor  $\Upsilon$  is used interchangeably with particle energy  $\Upsilon mc^2$ .)

above from 2D simulations (e.g., Sironi&Spitkovsky 2014, Guo et al 2015, Werner et al 2016)



## Reconnection parameters

Upstream Parameters: pair plasma

$n_b, T_b, B_0, B_{gz}$  (guide field)

Dimensionless parameters:

$T_b / m_e c^2 \gg 1$  (ultrarelativistically-hot)

$$\sigma = \frac{B_0^2}{4\pi n_b m_e c^2} \gg \frac{T_b}{m_e c^2} \quad (\text{relativistic reconnection})$$

$$\sigma_h = \frac{B_0^2}{4\pi h} \approx \frac{B_0^2}{4\pi(4n_b T_b)} = 25, 100 \quad \text{relativistic reconnection: relativistic outflows, significant energization}$$

$B_{gz}/B_0 = 0, 0.25, 0.5, 1, 2$  (guide field)

Later: vary inverse Compton (IC) radiative cooling

nominal length scale:

$$\rho_0 = \frac{m_e c^2}{eB_0}, \rho_c = \sigma\rho_0$$

System size:  $L_x, L_y, L_z$

$L_x/\sigma\rho_0 = 40$  (3D) – 320 (2D)

$L_y/L_x = 2$

$L_z/L_x = \text{varying (2D} \rightarrow \text{3D)}$

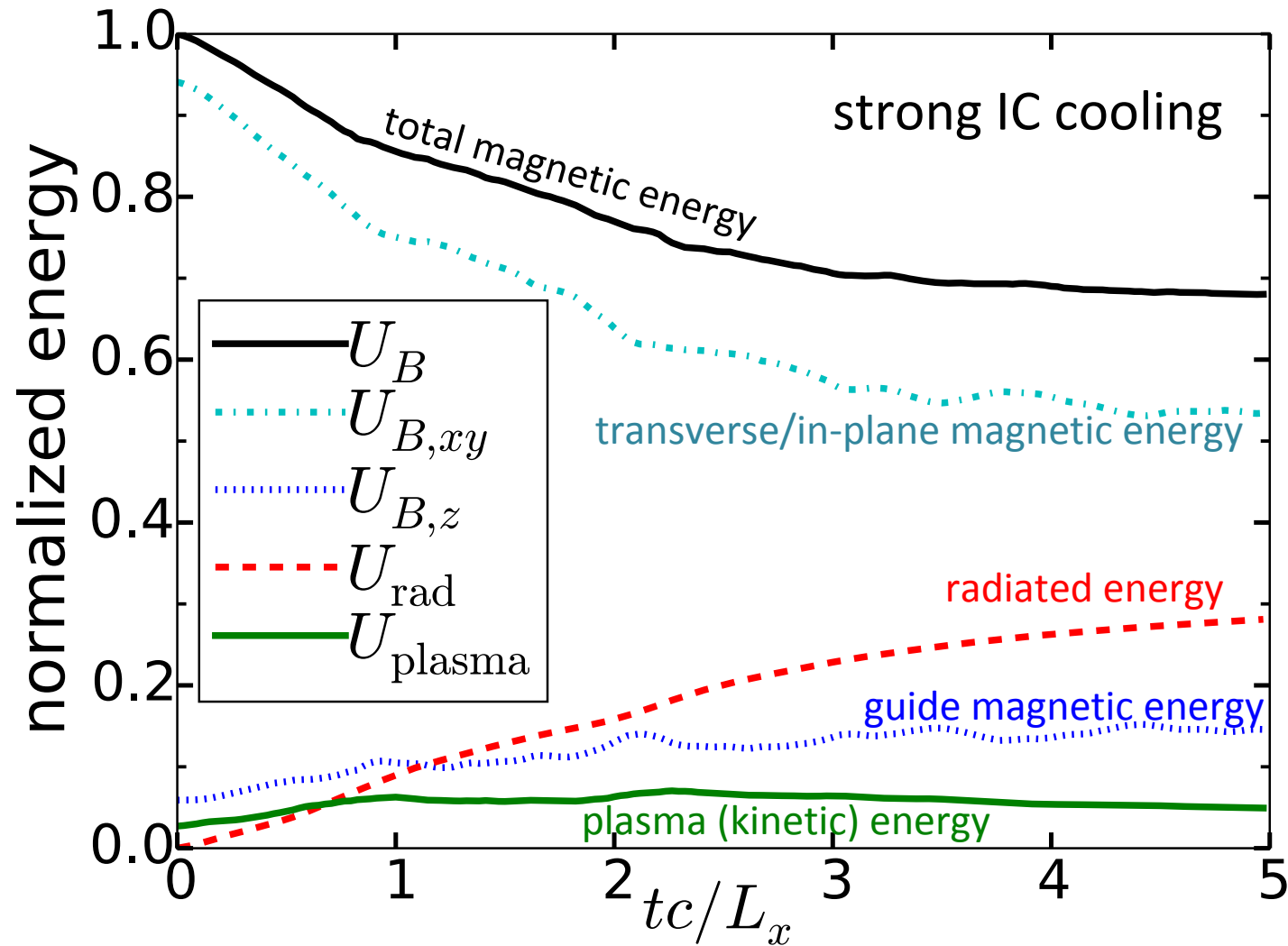
How do these parameters affect reconnection?

Specifically:

- energetics
- NTPA

Focus on two “outputs” of reconnection: **basic dynamics/energetics**, and NTPA

reconnection rate, magnetic energy dissipation, plasmoid formation, etc.

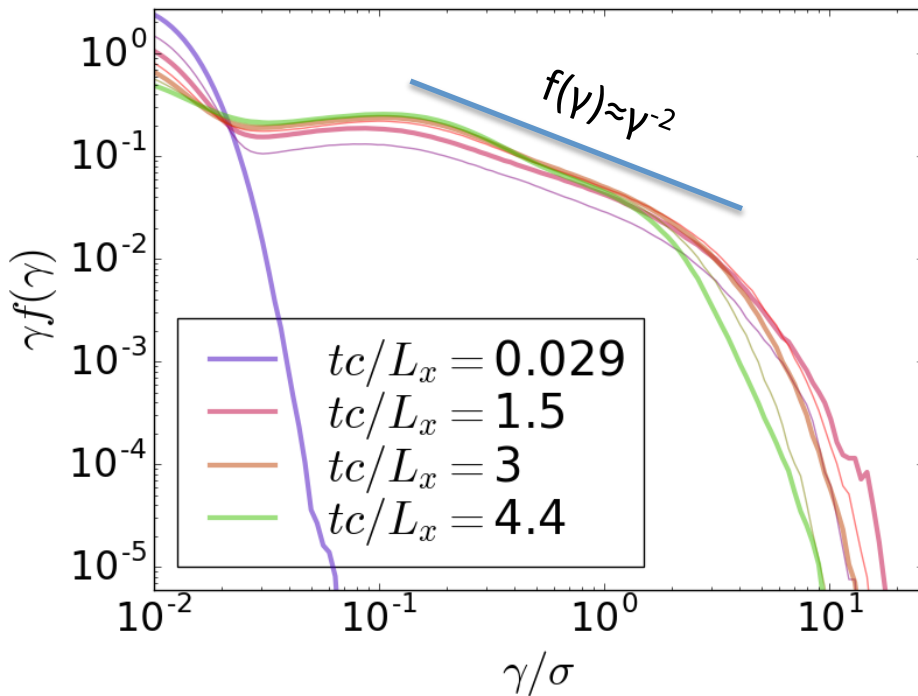


# Focus on two outputs of reconnection: basic dynamics/energetics, and **NTPA**

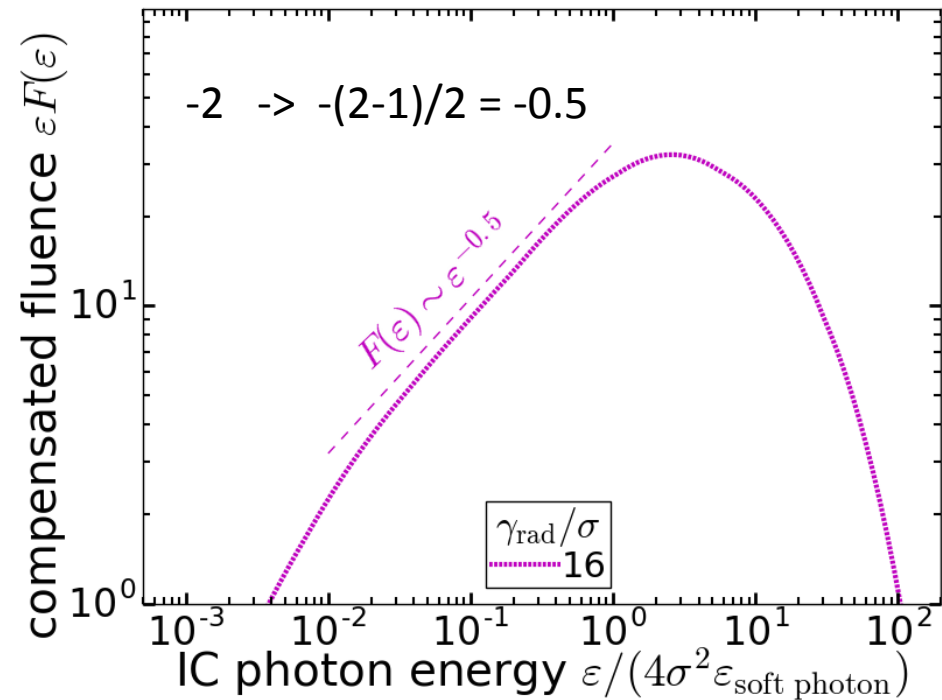
**NTPA:**

(shown here, for weak IC cooling)

particle energy spectra



photon energy spectra



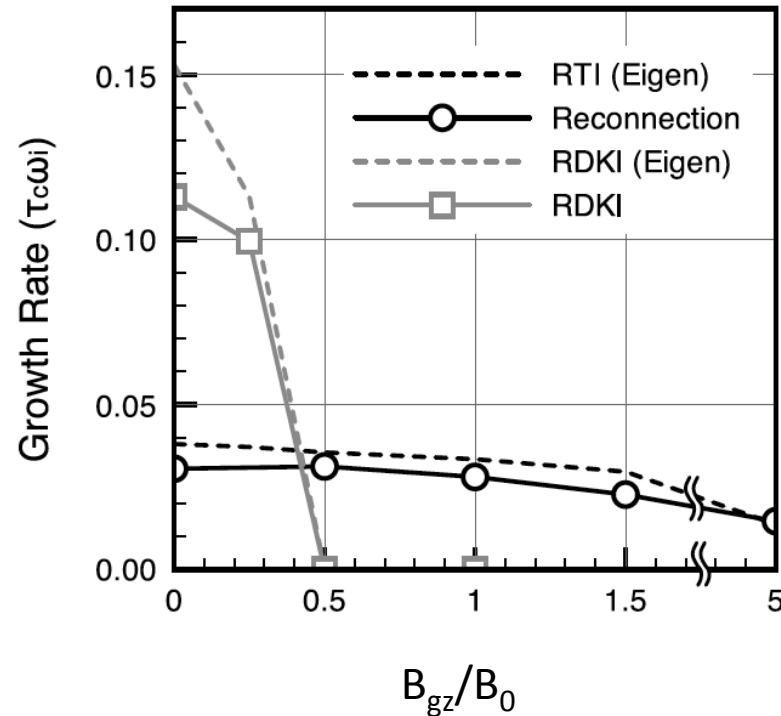
In the following, various input parameters will be varied ( $B_{gz}$ ,  $Lz/Lx$ , IC cooling) with outputs (dissipated magnetic energy and NTPA) shown.

Vary  $L_z/L_x$  (3D-ness) and  $B_{gz}/B_0$  and see what happens...

## 3D effects: does $L_z/L_x$ affect reconnection?

In particular, does the relativistic drift-kink instability (RDKI) inhibit particle acceleration? Here, guide magnetic field may be important: it inhibits RDKI.

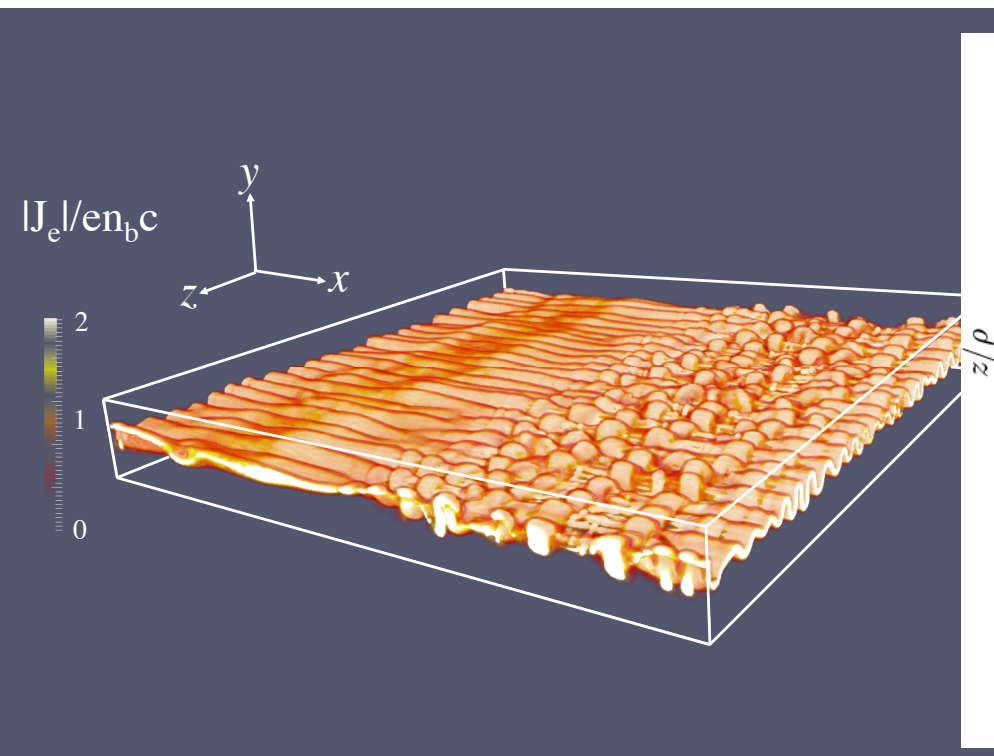
from Zenitani & Hoshino, 2008:



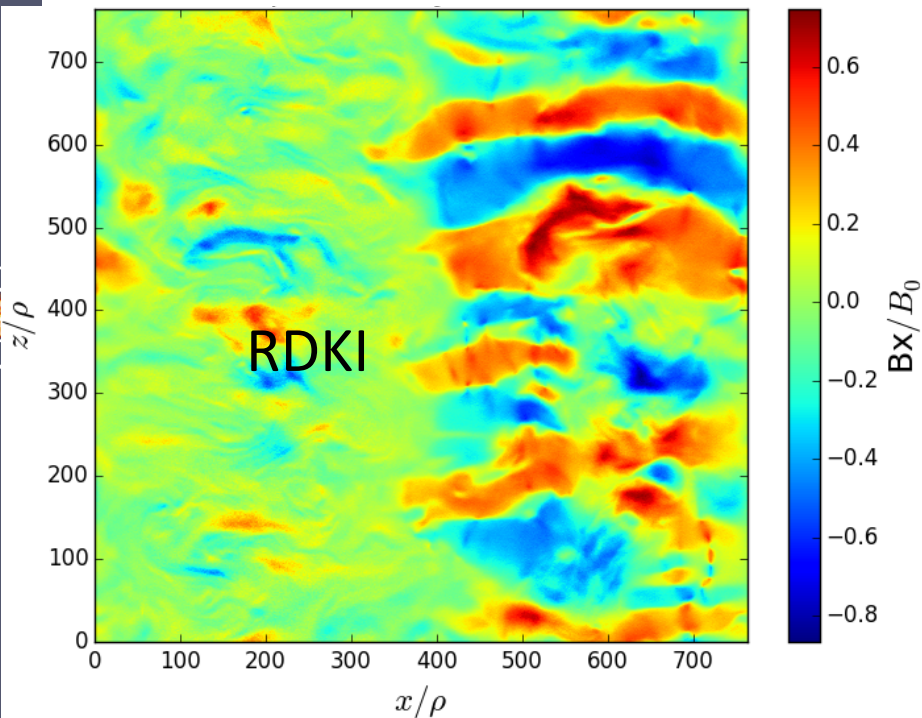
However, more recent simulations (e.g., Sironi & Spitkovsky 2014, Guo et al 2015, Werner & Uzdensky 2017) have suggested that particle acceleration is robust to 3D effects.



Despite significant RDKI, 2D and 3D reconnection have similar reconnection rates and NTPA.



$B_x$  in the x-z reconnection midplane

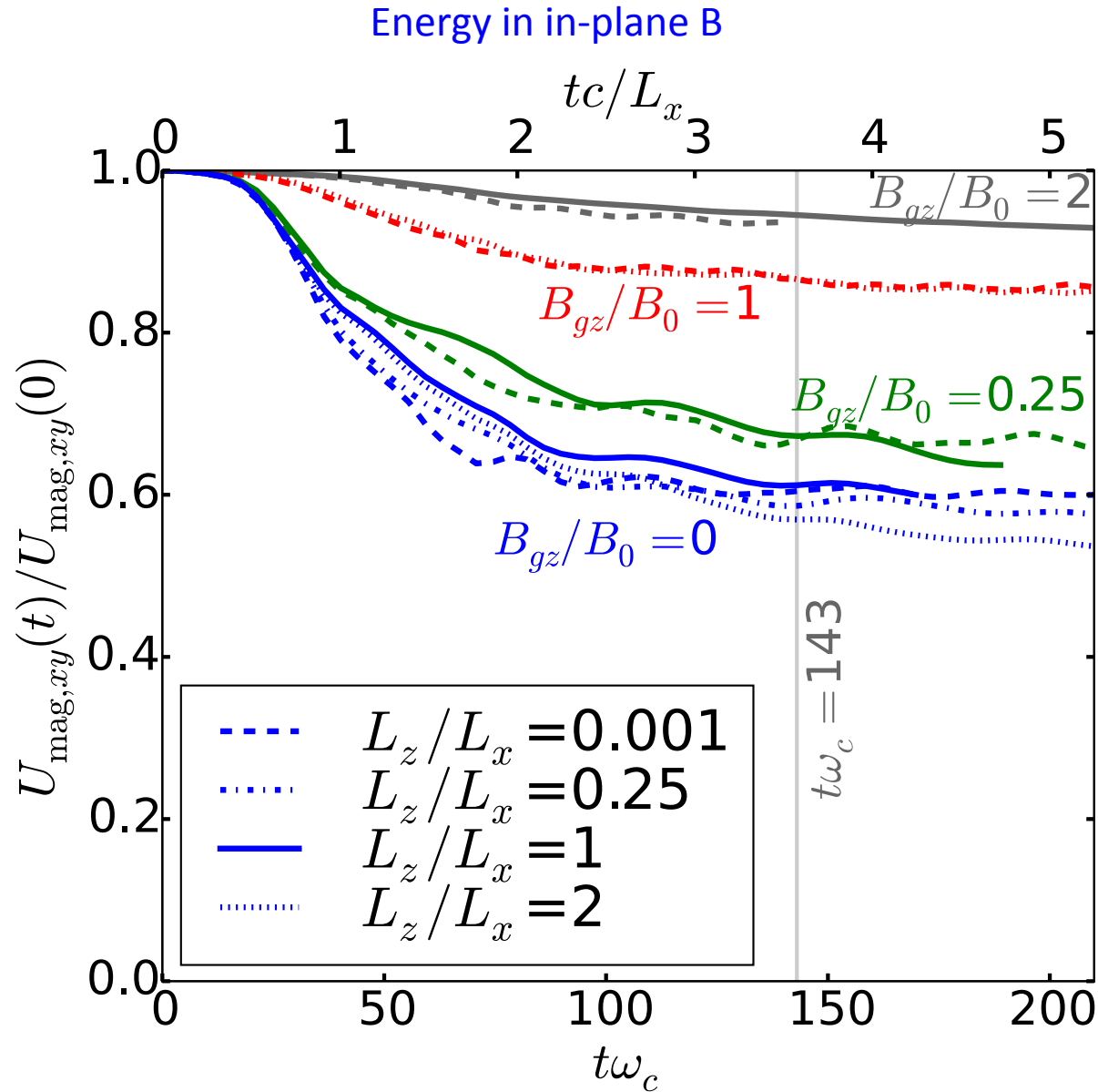


3D,  $L_z=L_x$ ,  $B_z=0$

# 3D current sheet evolution



Energetics of 2D and 3D reconnection are similar regardless of guide field  
(for later: guide field has a significant effect)



$\sigma_h = 25$

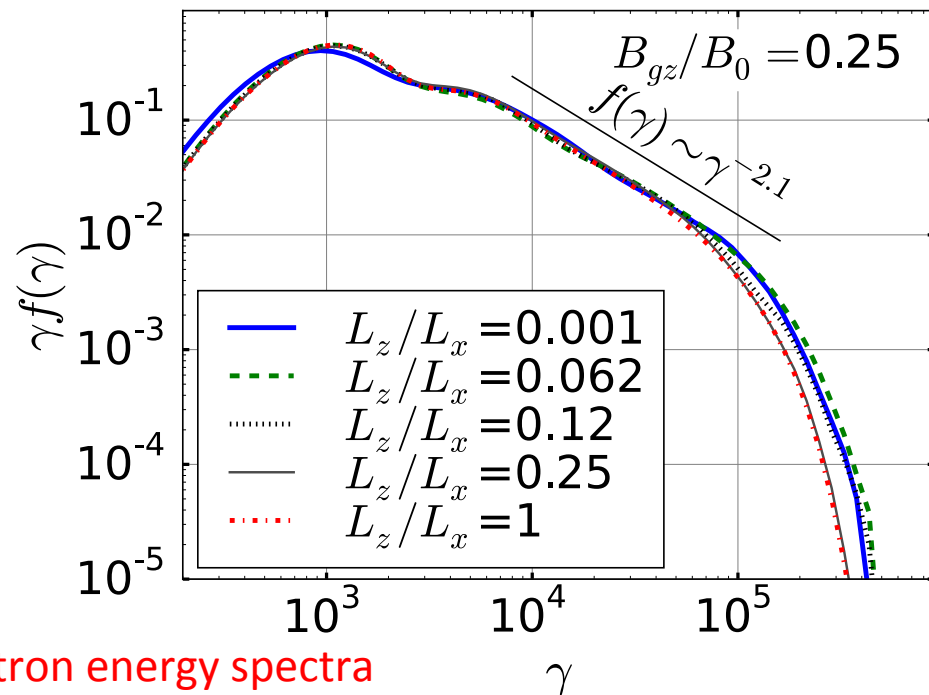
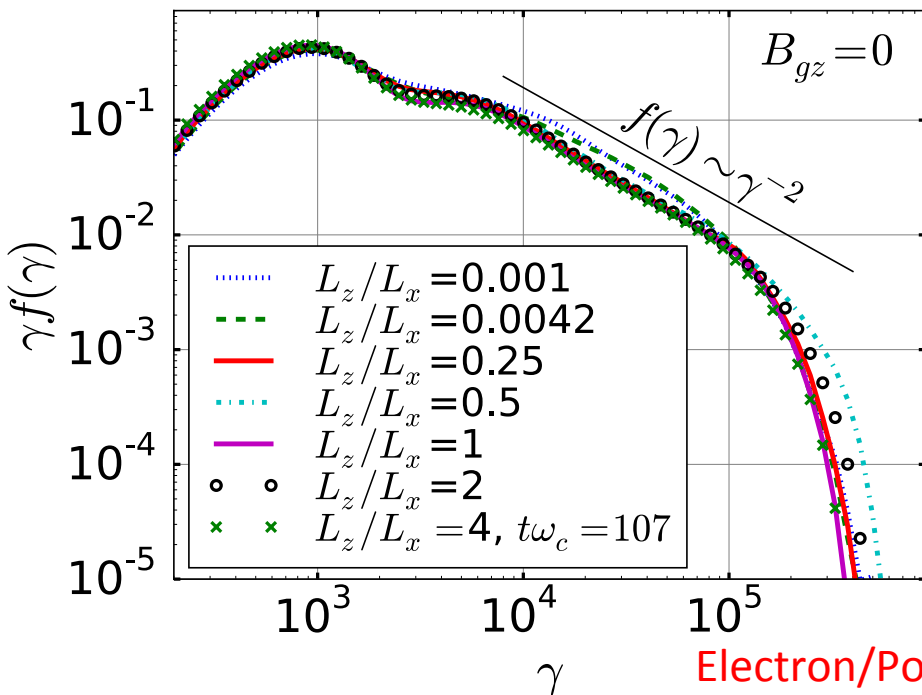
And 2D and 3D particle spectra are similar!

Nonthermal acceleration remains robust from 2D to 3D!

Also, a little guide field  $B_{gz}$  hardly disturbs acceleration.

$$L_x = 40\sigma\rho_0$$

$$L_y = 80\sigma\rho_0$$



Electron/Positron energy spectra

$$\sigma_h = 25$$

## Compressing plasmoids



$$n/n_b$$

During reconnection, the in-plane magnetic field compresses plasmoids.

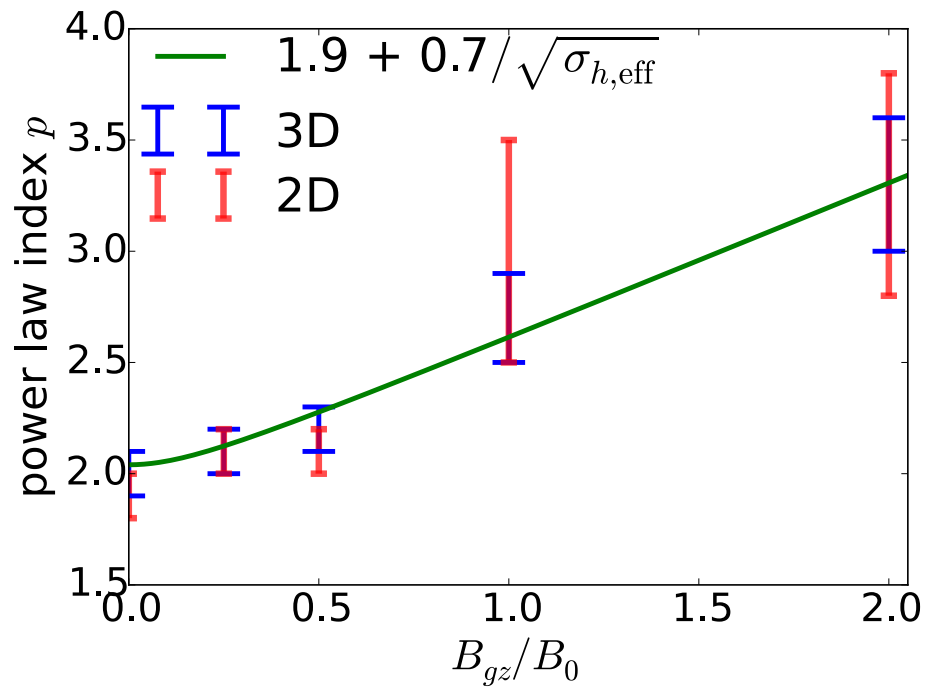
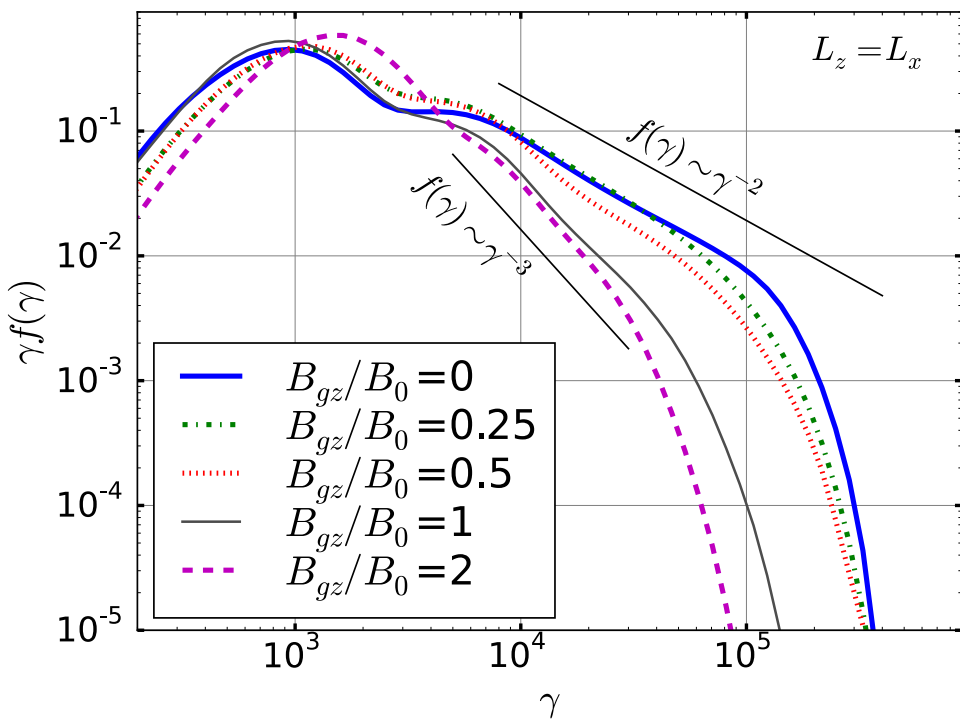
When there's a guide field, that guide field resists compression. This slows reconnection and inhibits particle acceleration.

# Guide field not only slows reconnection rate, but steepens the NTPA power law.

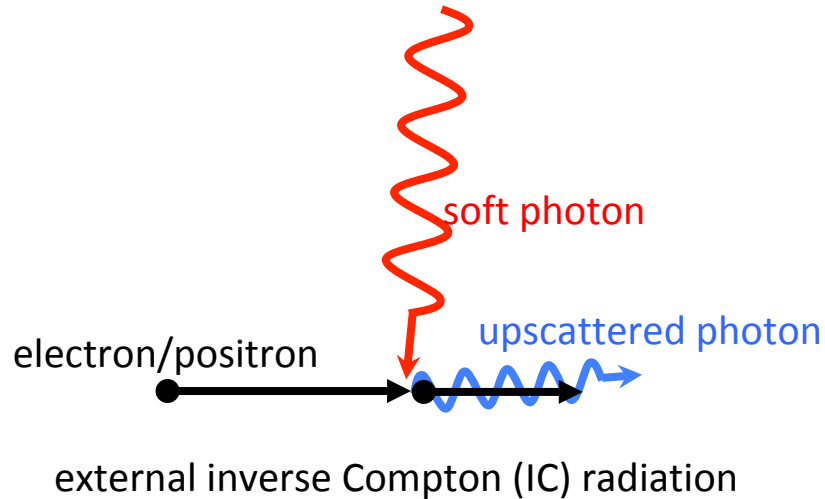
Guide field slows reconnection, dissipates less magnetic energy (guide field resists compression).

$$\sigma_h = 25$$

$$\sigma_{h,\text{eff}} = \frac{B_0^2 / 4\pi}{h_{\text{particle}} + B_z^2 / 4\pi}$$



## Reconnection in a bath of soft (low-energy) photons



High energy electrons (or positrons) scatter of photons, emitting high energy photons, and experiences radiation reaction (radiation) force.

If  $U_{ph}$  is the photon energy density, then the power loss, for an electron with  $\gamma m_e c^2$  is:

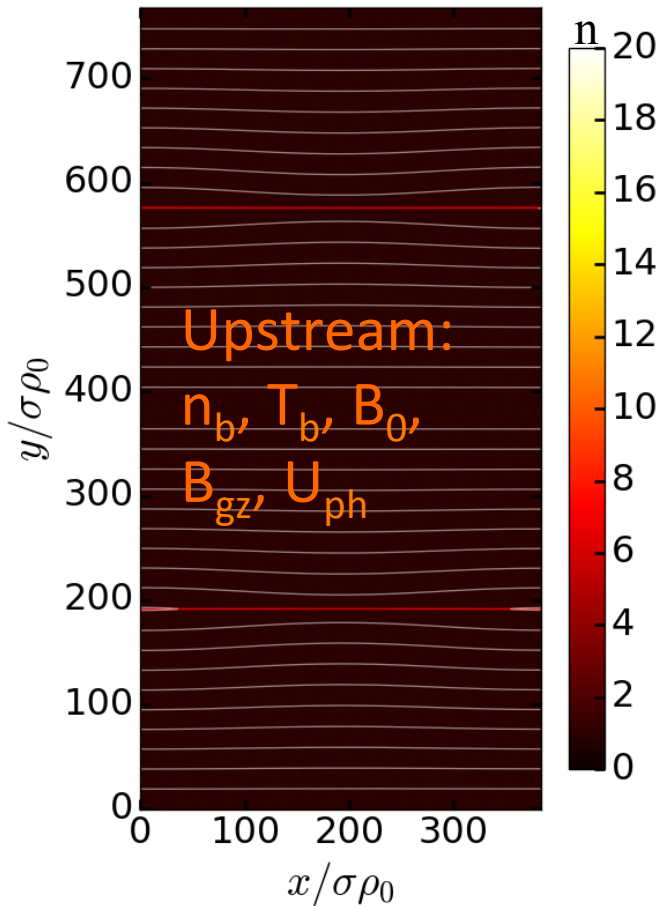
$$P_{rad} = \frac{4}{3} \sigma_T c U_{ph} \gamma^2$$

Power gain (accel.) in the reconnection electric field  $E=0.1B_0$ :  $P_{acc} = (0.1)eB_0 c$

These 2 forces (powers) balance for  $\gamma = \gamma_{rad}$ : 
$$\gamma_{rad} = \sqrt{\frac{3(0.1)eB_0}{4\sigma_T U_{ph}}}$$

Particles can't gain much more energy than this.

## Reconnection setup with photon bath



Upstream Parameters: pair plasma  
 $n_b, T_b, B_0, B_{gz}$  (guide field),  
 $U_{ph}$  (soft radiation bath energy density)

Dimensionless parameters:

$T_b / m_e c^2 \gg 1$  (ultrarelativistically-hot)

$$\sigma = \frac{B_0^2}{4\pi n_b m_e c^2} \gg \frac{T_b}{m_e c^2} \quad (\text{relativistic reconnection})$$

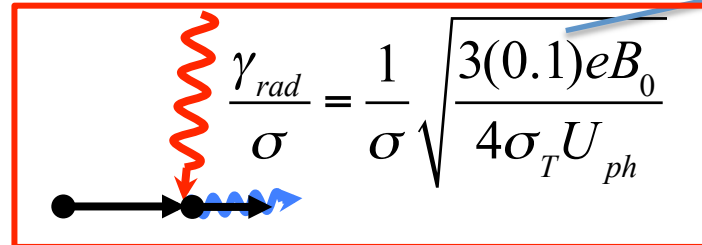
$$\sigma_h = \frac{B_0^2}{4\pi h} \approx \frac{B_0^2}{4\pi(4n_b T_b)} = 25, 100 \quad \text{relativistic reconnection: relativistic outflows, significant energization}$$

$B_{gz}/B_0 = 0, 0.25, 0.5, 1, 2$

nominal length scale:

$$\rho_0 = \frac{m_e c^2}{eB_0}, \rho_c = \sigma\rho_0$$

System size:  $L_x, L_y, L_z$



nominal reconnection rate

$\gamma_{rad} m_e c^2$  is the energy at which acceleration by reconnection  $E$  equals deceleration by IC radiation reaction.

$L_x/\sigma\rho_0 = 40$  (3D) –  $320$  (2D)

$L_y/L_x = 2$

$L_z/L_x = \text{varying}$  (2D→3D)

How do these parameters affect reconnection?



IC scattering doesn't affect basic reconnection dynamics very much

$\gamma_{\text{rad}} = \infty$  (no cooling)

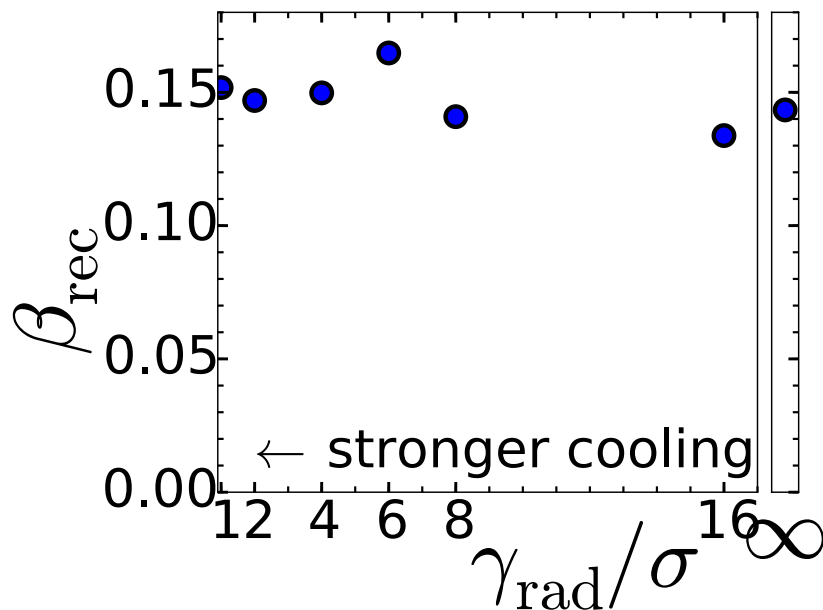
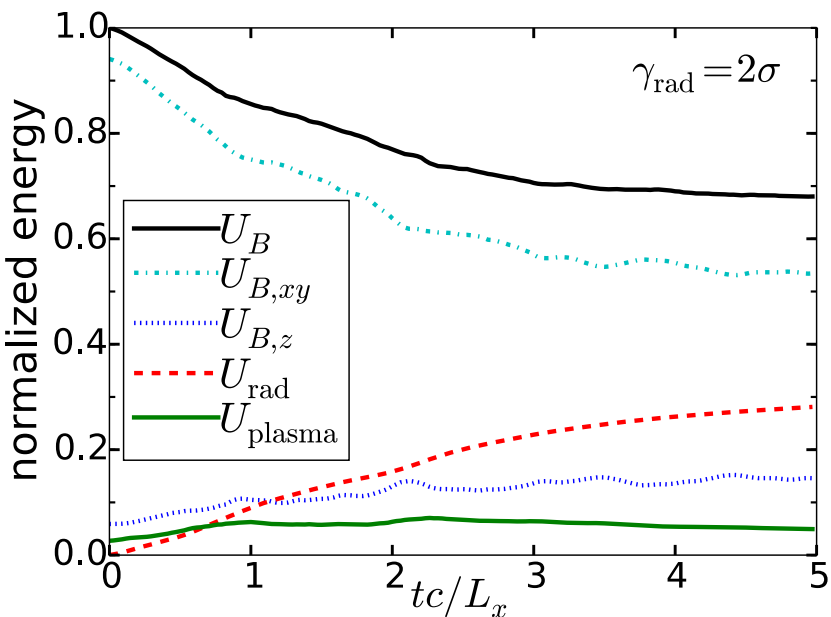
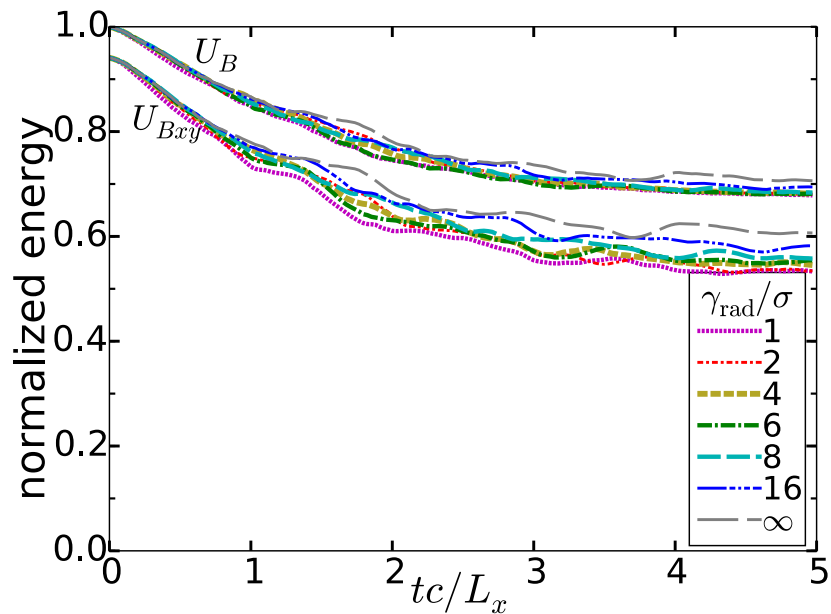
$\gamma_{\text{rad}} = 2\sigma$  (strong cooling)



color=plasma density (normalized to  $n_b$ )

IC cooling has little effect on magnetic energy dissipation, reconnection rate

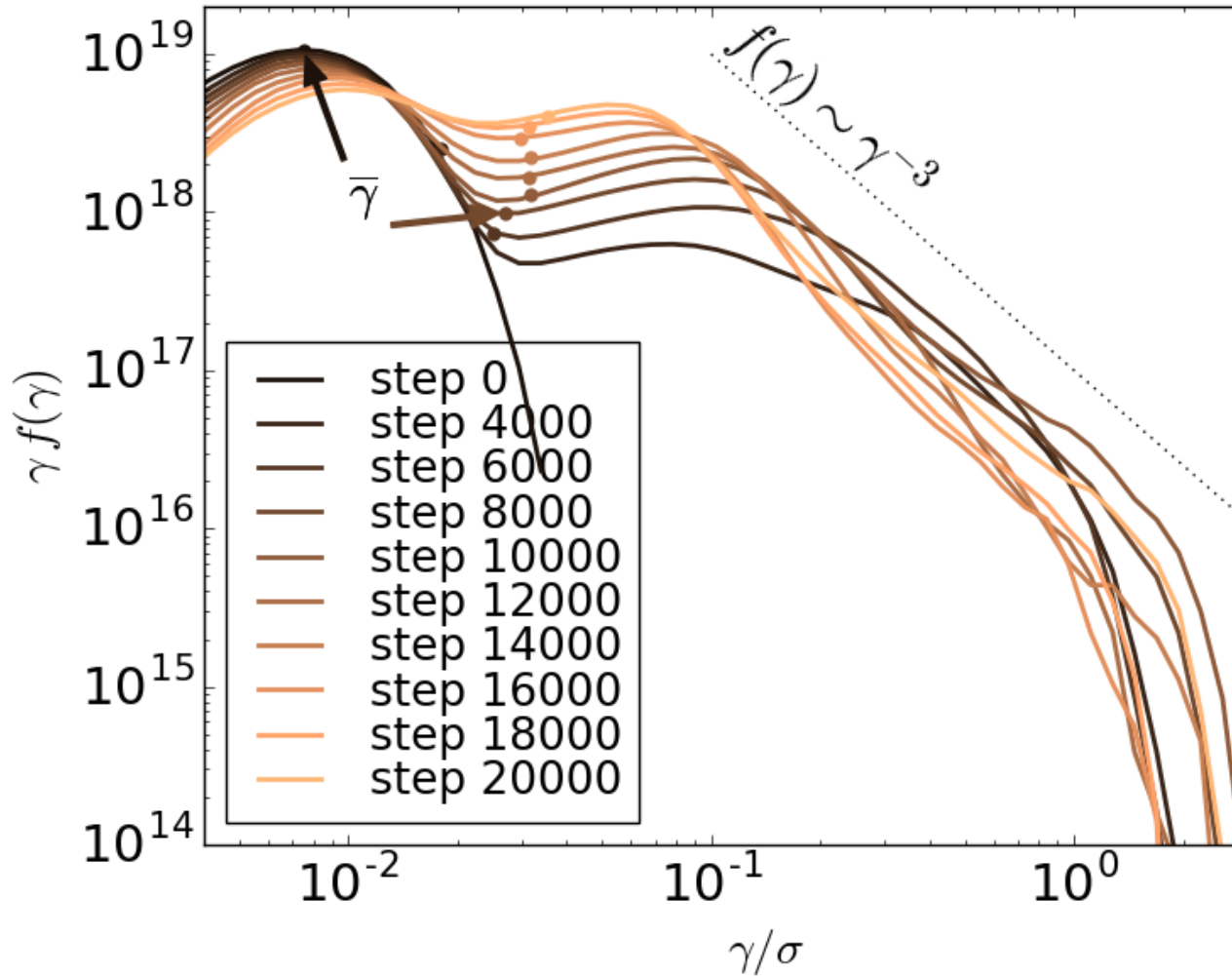
$$\sigma_h = 100, B_{gz} = B_0/4$$



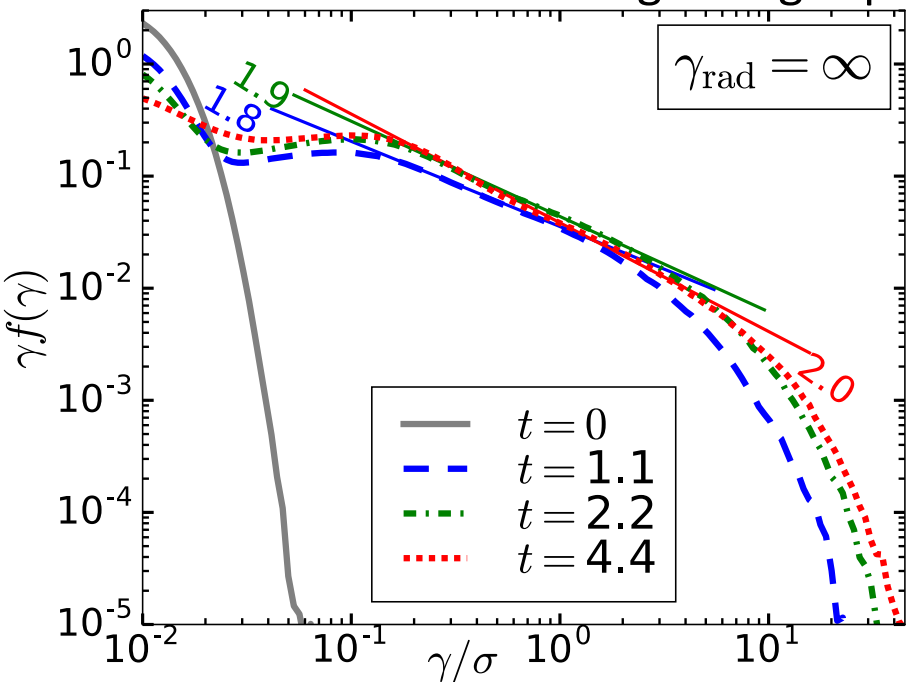
Strong cooling doesn't alter the amount of magnetic energy transferred to particles...but strong cooling means particles promptly radiate that energy.

IC cooling changes particle spectra significantly: noisy, steeper

$$\sigma_h = 100, B_{gz} = B_0/4$$

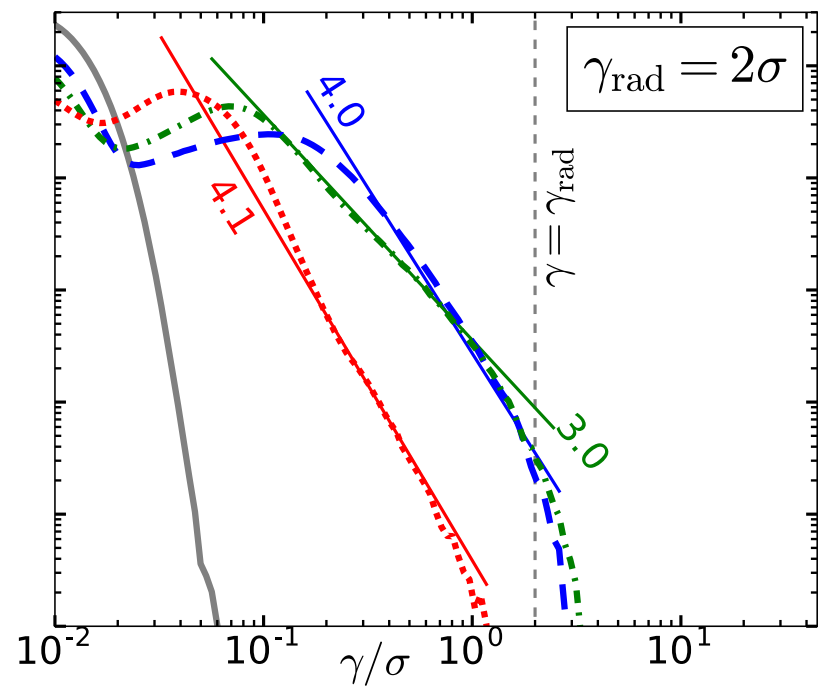
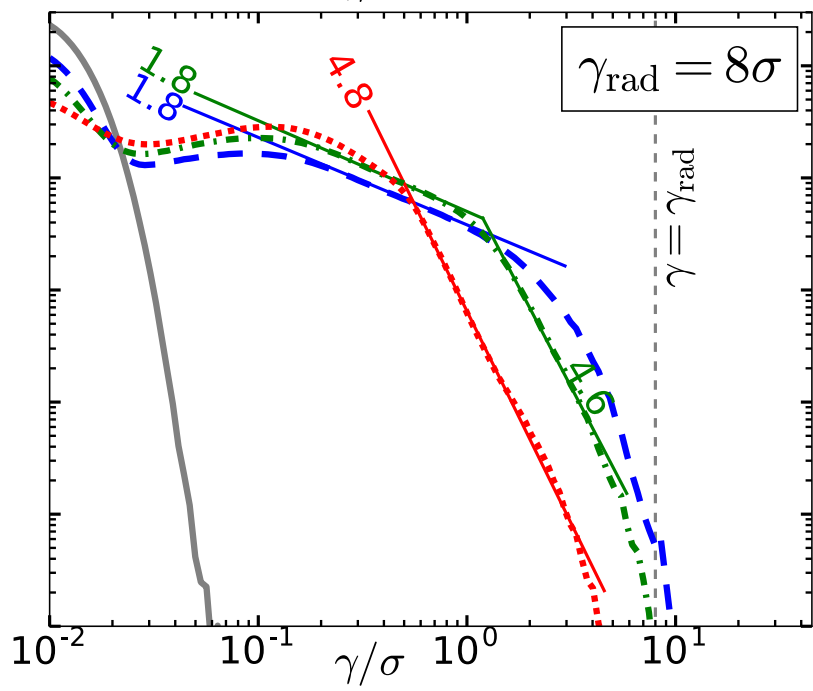


# IC cooling changes particle spectra significantly



Weak cooling: usual hard power law  
 Strong cooling: variable steep power law  
 Intermediate: both power laws

$$\sigma_h = 100, B_{gz} = B_0/4$$



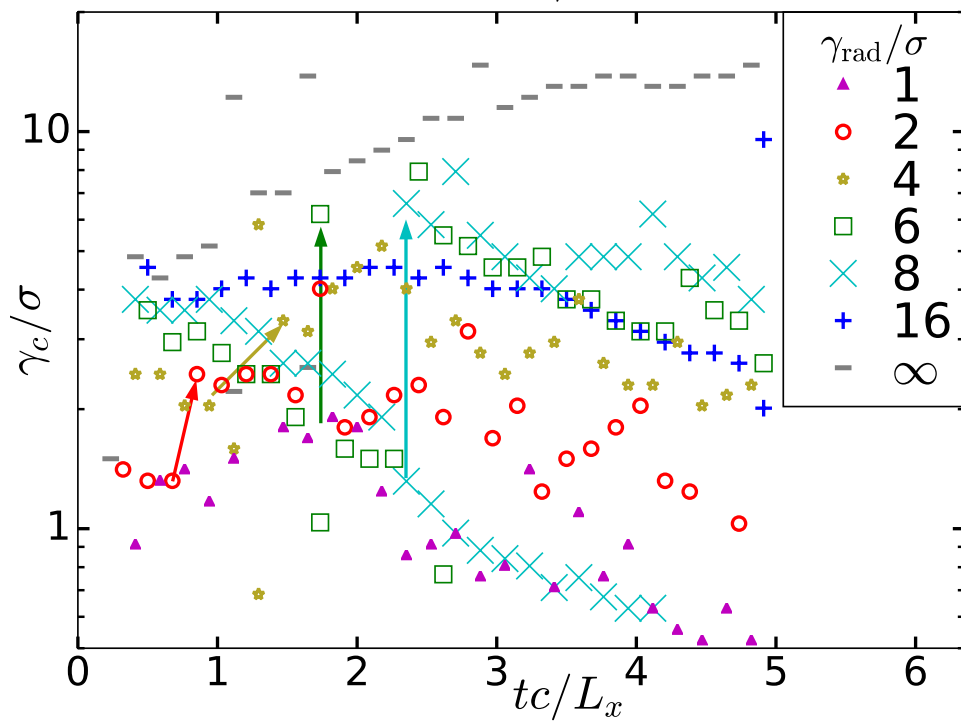
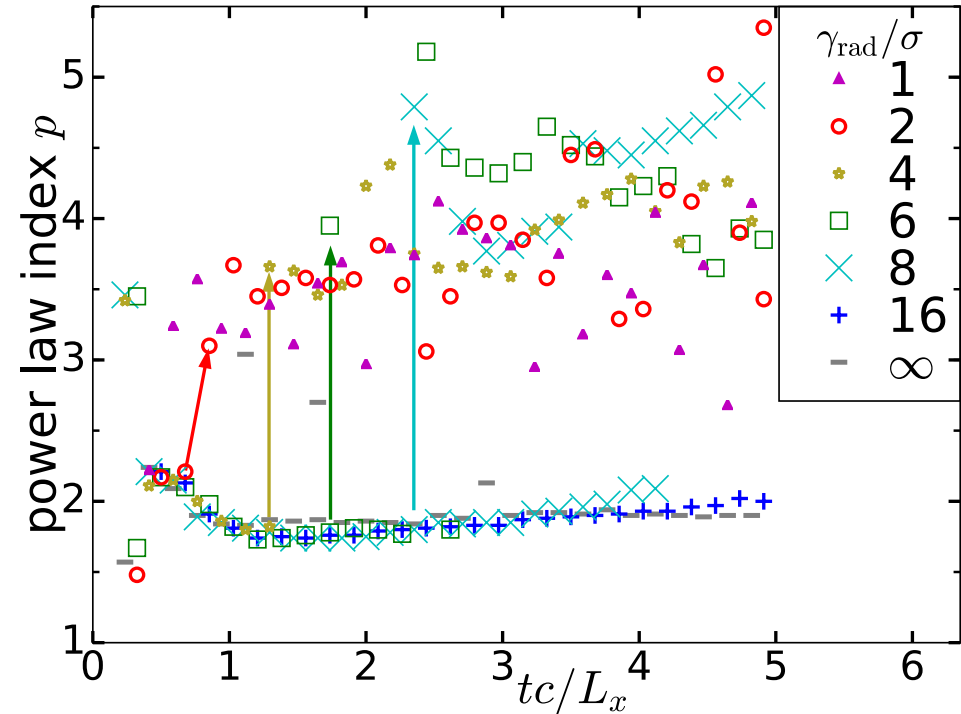
Time-dependence of power laws shows both power laws present (mostly); steep power law is highly variable

Regardless of IC cooling, (plasmoid-dominated) reconnection is bursty process with discrete acceleration episodes that yield NTPA spectra with a hard slope ( $p_h=1.9$  in this case).

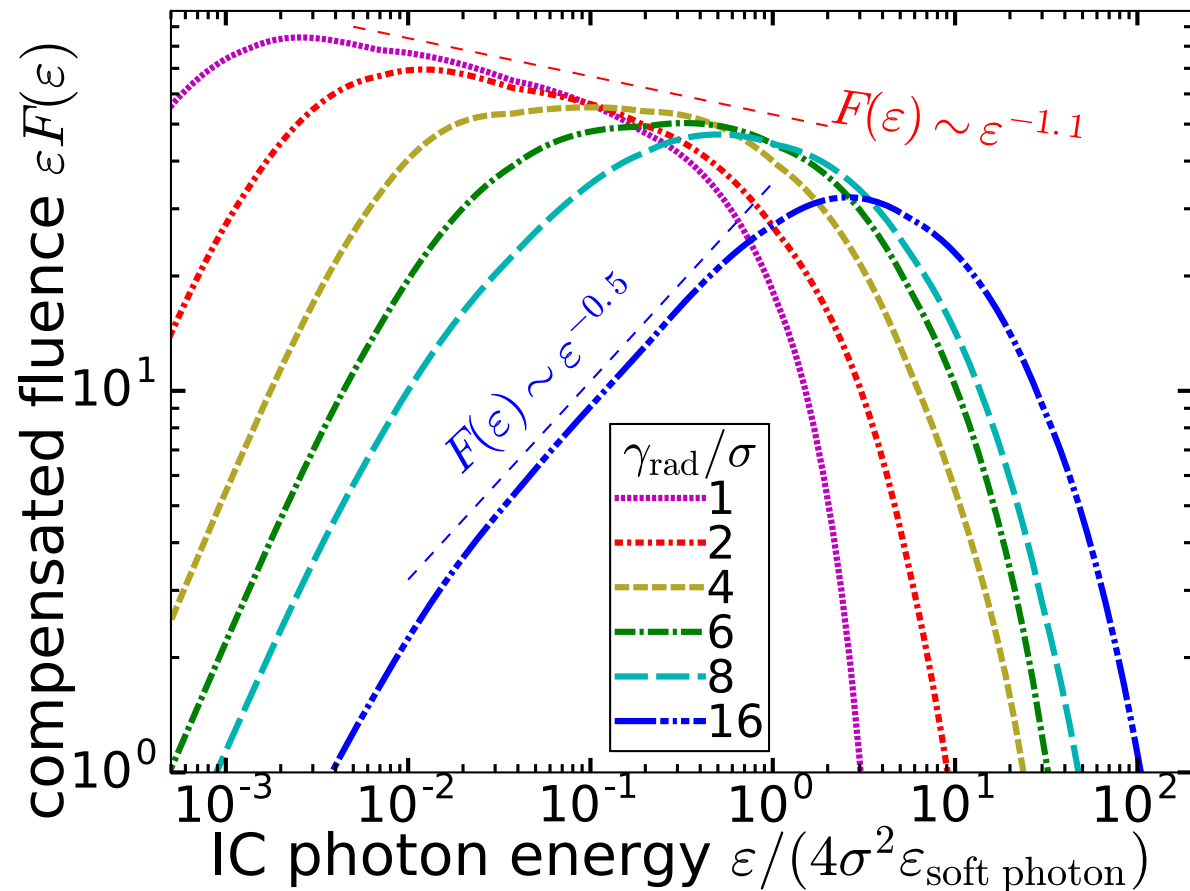
Cooling occurs between episodes, steepening the slope. Depending on episodes of acceleration and cooling, the steep slope  $p_s$  falls between 3 and 5.

Continuous acceleration/cooling would yield  $p_{s,\min}=p_h+1=3$ ...but additional cooling results in higher  $p_s$ .

$$\sigma_h=100, B_{gz}=B_0/4$$



# Time-integrated IC photon spectra



Photon power law index  $\alpha = (p-1)/2$ .

Hard slope  $p_h=1.9 \rightarrow \alpha = 0.45$  (measured 0.5)

Steep slope  $p_s=3-5 \rightarrow \alpha = 1-3$

however: a harder slope means more IC emission,

so  $\alpha$  should be dominated by the hardest  $p_{s,\text{min}}=3 \rightarrow \alpha=1$  (measured 1.1)

In this particular case (ultrarelativistic pair plasma,  $\sigma_h=100$ ,  $B_{gz}=B_0/4$ ),  
adding a soft photon bath changes index from  $\alpha=0.5$  to  $\alpha=1.1$ .

## Simulation comparison: TRISTAN-MP and Zeltron

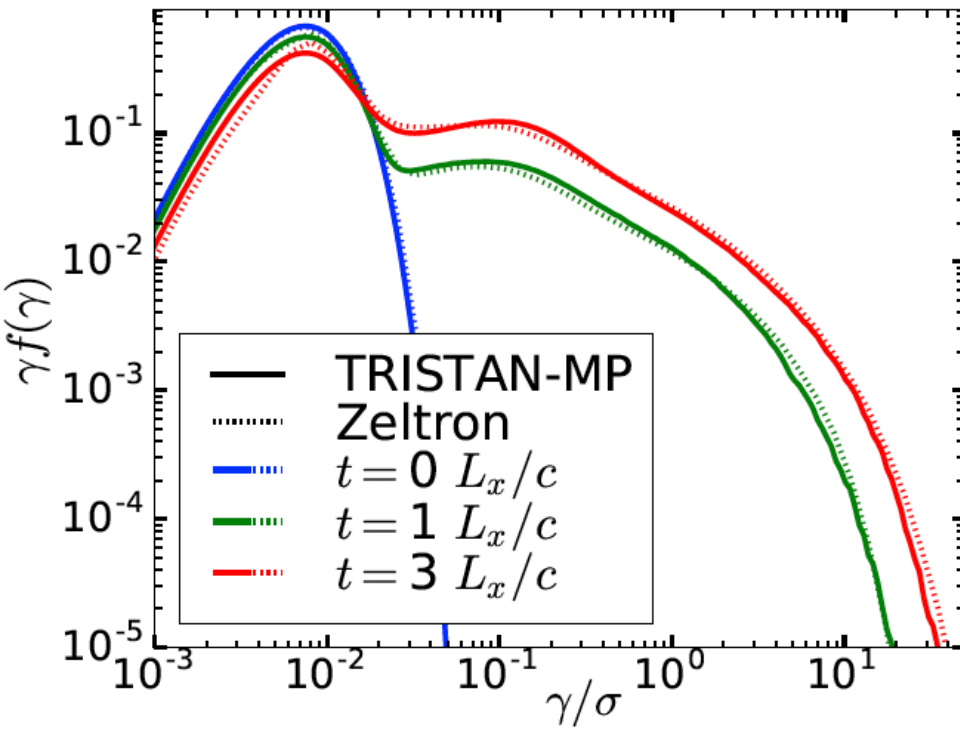
Both codes implement same fundamental algorithms: explicit EM-PIC with minor variants.

Both (as of this year) use charge-conserving current deposition (div E = rho is automatically maintained to high precision), though different variants.

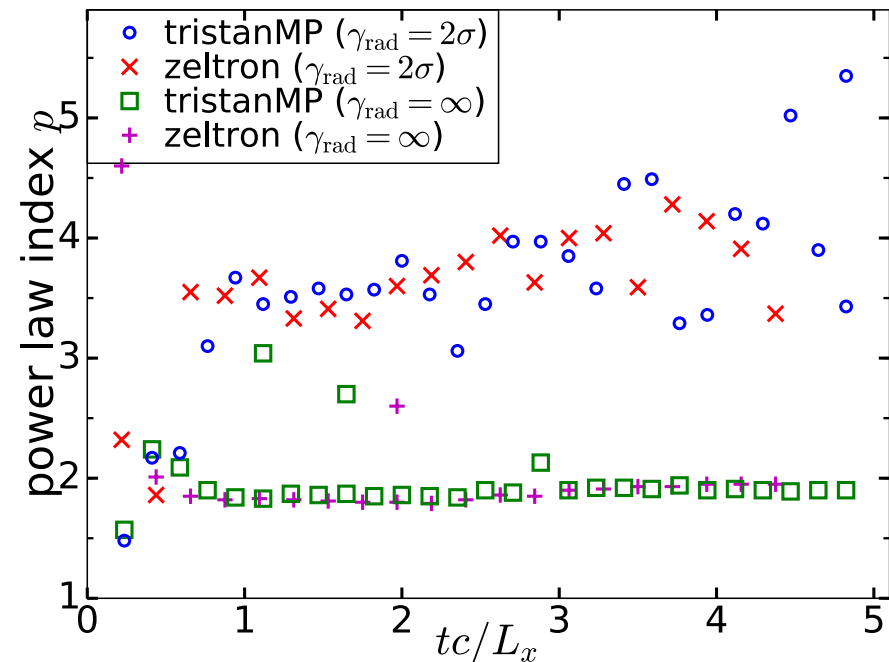
Both codes implement IC radiation reaction force (in somewhat different ways).

The implementations are entirely independent. Do they agree? Yes; very well.

Particle spectra (no IC cooling)



Spectra power law indices



## Conclusions

- 3Dness ( $L_z/L_x$ ) has little effect – despite significant RDKI – on reconnection rate and NTPA
- Guide field slows reconnection and inhibits NTPA
  - magnitude of effect depends on guide field enthalpy vs. particle enthalpy
    - if guide field enthalpy is large, the guide field resists compression and slows reconnection
    - if guide field enthalpy is small (compared to particle enthalpy), not much effect
- IC cooling (drag due to radiation reaction) has little effect on reconnection rate
- IC cooling significantly affects NTPA
  - Particle spectrum forms a broken power law, with
    - a hard slope  $p_h$  (independent of IC cooling strength)
    - a highly-variable steep slope  $p_s$ , with  $p_s > p_h + 1$  (also independent of cooling)
      - $p_s = p_h + 1$  would mean continuous acceleration and cooling
      - $p_s > p_h + 1$  for episodes of acceleration followed by further cooling
    - a break that decreases in energy as cooling strength increases
  - For very weak cooling, the break is above the reconnection-high-energy-cutoff and only the hard power law appears;
  - for intermediate cooling, both power laws are visible;
  - for strong cooling, the hard power law appears only at the very beginning before being overwhelmed by the steep power law
- The IC radiation spectrum varies with the particle spectrum.
  - For weak cooling,  $p_h = 1.9 \rightarrow \alpha = 0.5$
  - For strong cooling,  $p_s$  varies,  $p_s \geq 3 = p_h + 1$ , but the hardest component dominates so the photon spectrum corresponds roughly to  $p_s = 3$ , or  $\alpha = 1$  (measured 1.1).



