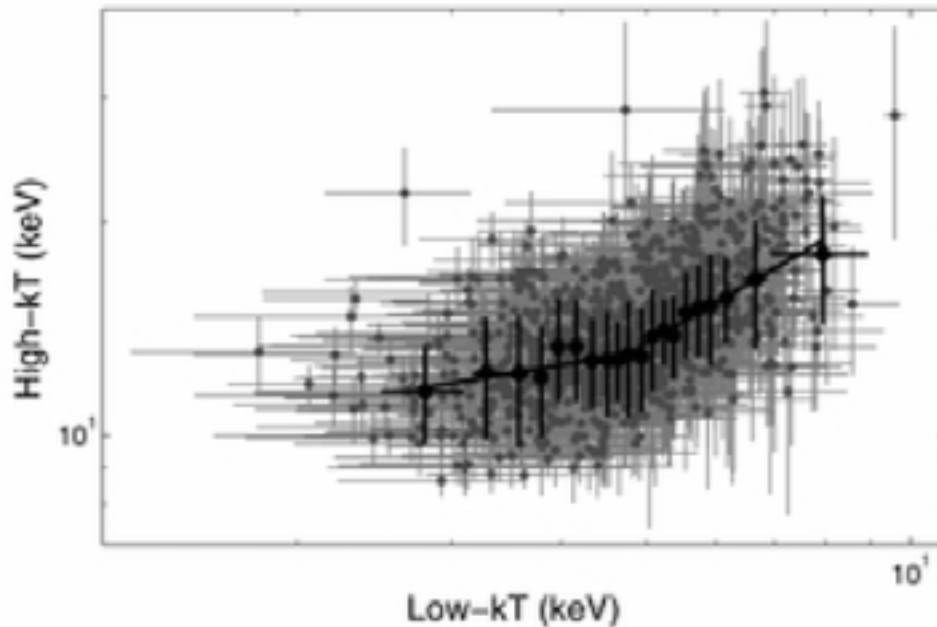


# Radiation in relativistic plasmas

**Maxim Lyutikov (Purdue)**

# Magnetar bursts - ver low T (few keV)



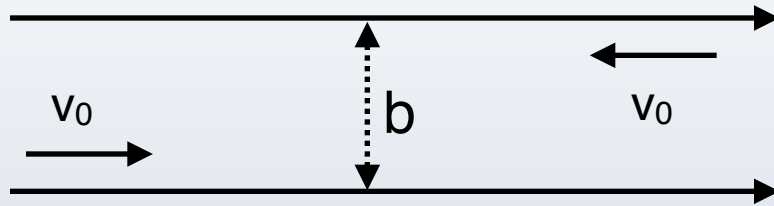
Younes +, 2014

- SGR J1550-5418, with RXTE bursts ~5 keV
- 1E1841-045, with NuSTAR bursts ~3.3 - 5 keV
- 1E1048.1-5937, with NuSTAR 6-8 keV
- 1E 1048.1-5937 with RXTE, ~3keV

Trapped pair plasma fireball: equilibrium pair density at few keV is minuscule  
Radiation processes frozen out - cannot cool

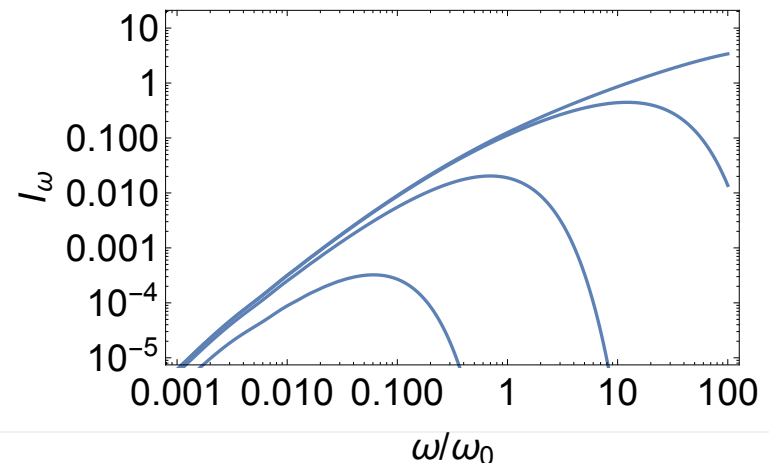
Something is wrong

# Free-free emission in $T \ll B$ (classical)



$$v = v_0 \sqrt{1 \pm \frac{e^2}{m_e v_0^2 \sqrt{b^2 + (2z)^2}}}$$

- 1D motion - lowest Landau level
- nearly classical, in a sense  $a_H = \sqrt{\frac{\hbar}{eB}} \rightarrow 0$
- Acceleration/emission usually  $\mathbf{a}_\perp$ , but now only  $\mathbf{a}_\parallel$
- Free-free emission:
  - $\omega \rightarrow 0, j_\omega \propto |\mathbf{v}_{in} - \mathbf{v}_{out}|^2$
  - Constant for  $B=0$
  - = 0 for 1D
- emissivity in single collision
  - $j_\omega = \frac{1}{\pi^2 c^3} a_\omega^2 \sin^2 \theta d\omega d\Omega$
  - low freq.  $\propto \omega^2$

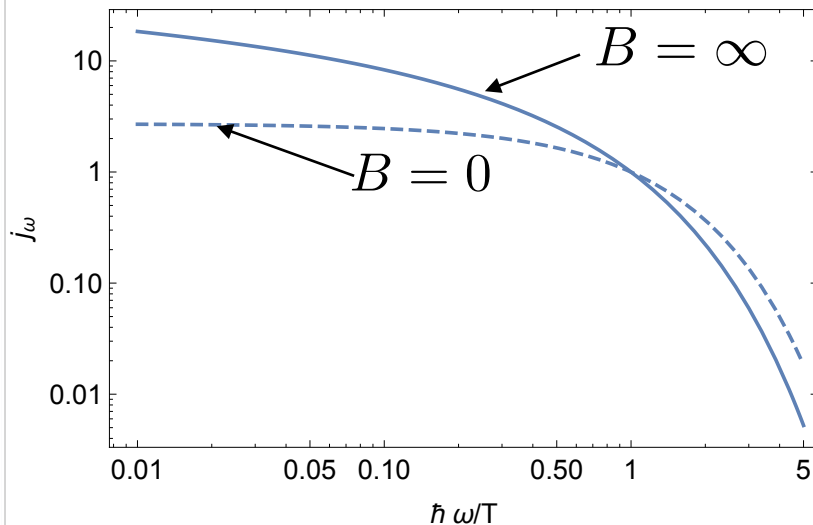


# Free-free in infinite B: only e<sup>-</sup>-e<sup>+</sup>

- Averaged over thermal

$$j_\omega = \frac{n_\pm^2}{32\sqrt{\pi}} \frac{e^6}{c^3 m_e^{3/2}} \frac{\sin^2 \theta}{\sqrt{T}} \Gamma_0 \left( \frac{\hbar\omega}{T} \right)$$

Less power at high freq.  $\propto e^{-\hbar\omega/T} \frac{T}{\hbar\omega}$



Fit magnetized free-free with B=0 formula

$$T_{\text{real}} = 12 T_{\text{fit}}$$

- Small absorption at  $> T$
- Compton redistribution

# Relaxation in 1D pair plasma

$e^-e^-$  &  $e^+e^+$



Keep the same  $|v|$   
No way to tell them apart

No relaxation due to binary collisions

Triple collisions  $\nu_3 \approx v_T n_{\pm} \pi \left( \frac{e^2}{T} \right)^2 \left[ \left( \frac{e^2}{T} \right)^3 n_{\pm} \right]$

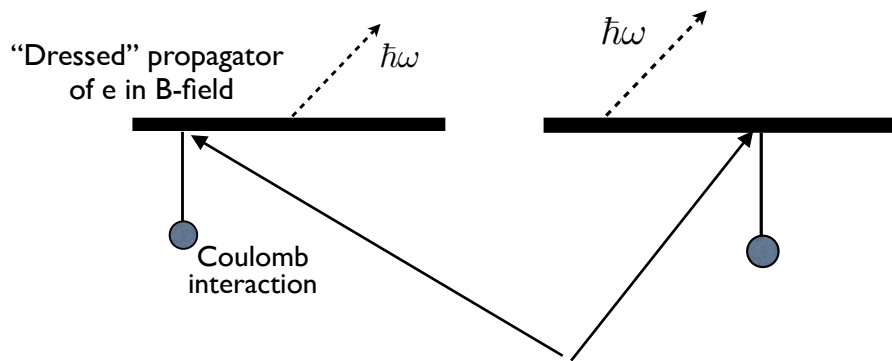
some 10 orders of magnitude smaller  
(also, no Coulomb logarithm!)

cross-section

relative # of 3 particles  
within strong interaction  
sphere

# Quantum at zero Landau

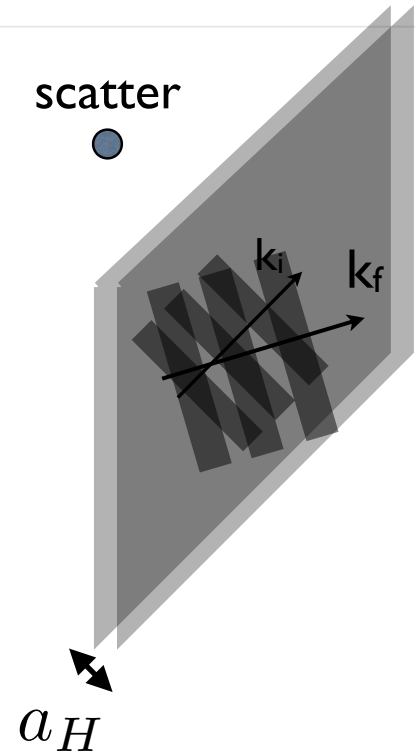
- Landau gauge
- sheets of 2D waves
- change in momentum within the 2D sheet
- Feinmann diagrams



Coulomb vertex for 0th Landau can be calculated exactly!

$$f = \frac{2\pi \exp\left(\frac{1}{4}a_H^2 \left((\kappa_{x,1} + \kappa_{x,2} + \Delta_\kappa)^2 - 2(\kappa_{x,1}^2 + \kappa_{x,2}^2)\right)\right)}{\Delta_\kappa} \rightarrow \frac{2\pi}{\Delta_\kappa}$$

Very similar to classical



scattering amplitude

# Cooling wave in pair plasma

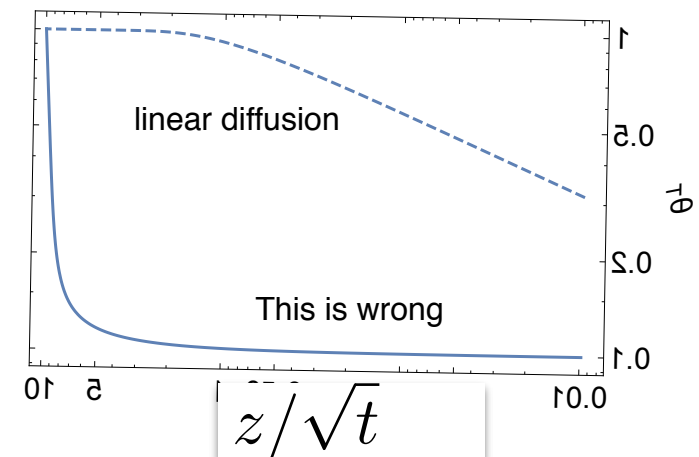
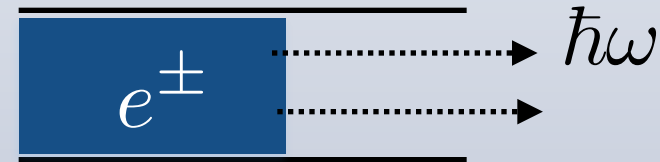
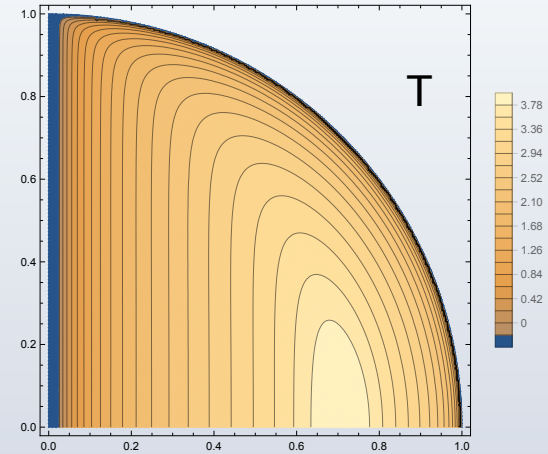
- Trapped fireball (a la spheromak/Hills vortex but with pressure)
- Cools by radiation - how?
- Half space filled with pair plasma, open the lid - how T evolves.

$$\dot{w} + F'_z = 0$$

$$w = \text{pairs} + \text{radiation}$$

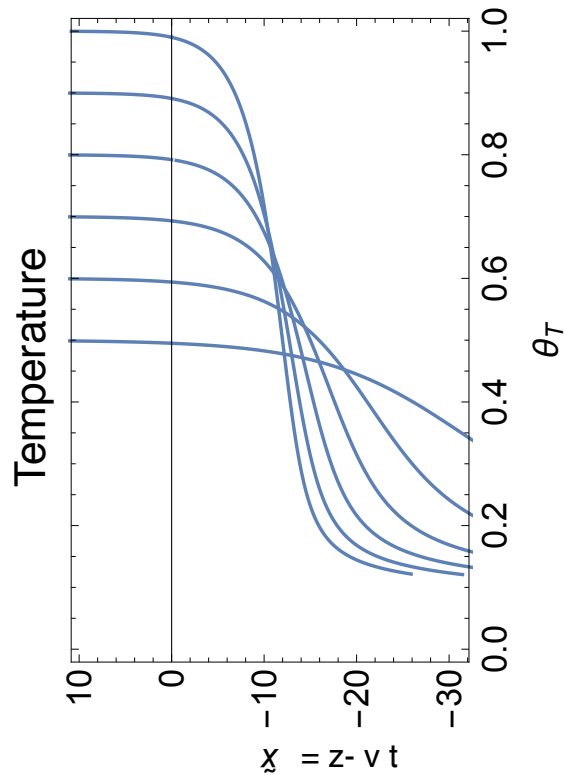
$$F = -\frac{4}{3} \frac{acT^3}{n\sigma_T} T'_z$$

- LTE (may not be good enough - photon production rate may not be able to catch up)
- Diffusion?



# Cooling wave

- Read Zeldovich & Raizer - cooling waves
- Nuclear explosions first few 100 meters
- Nonlinear diffusion: cooling wave regime  $T(z-v t)$

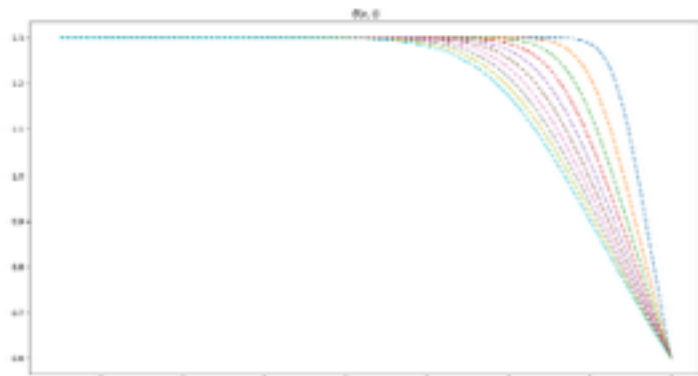


Qualitatively: non-linear diffusion

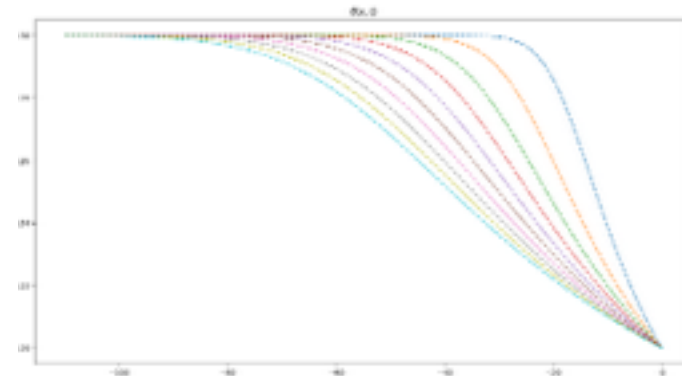
$$\frac{T''_{zz}}{T'_z} \propto \pm$$



# Full integration (with Jedidiah Riebing)



~ diffusive

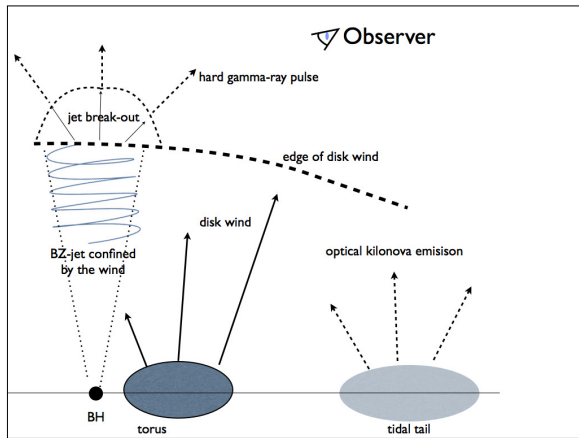
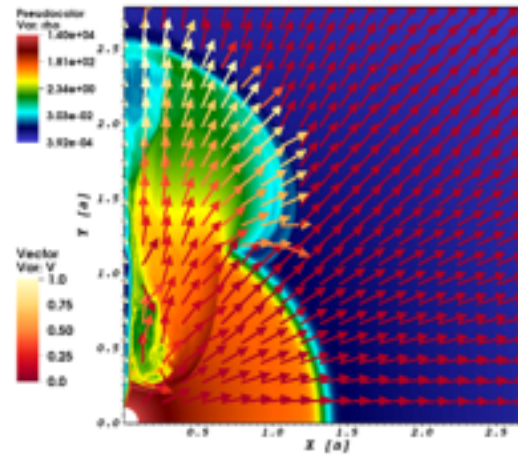
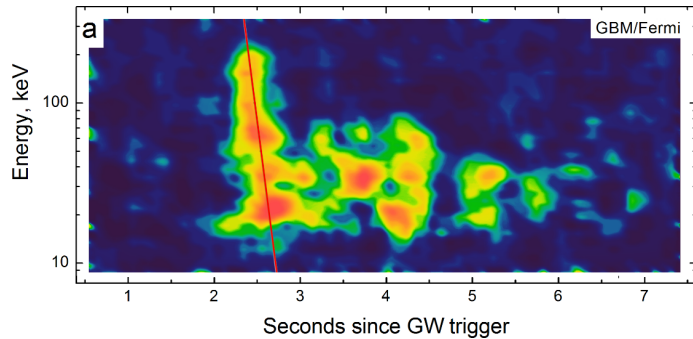


~ cooling wave

No clear separation of diffusive/cooling wave regime

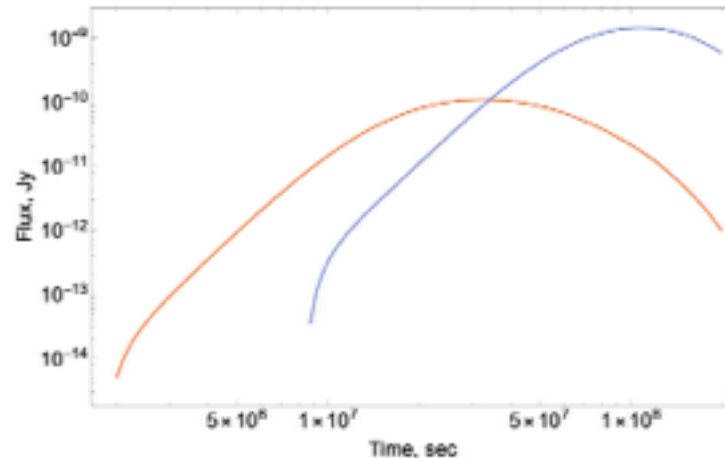
# Radiation and pairs in shocks

# Cocoon - prompt, Jet- to be seen



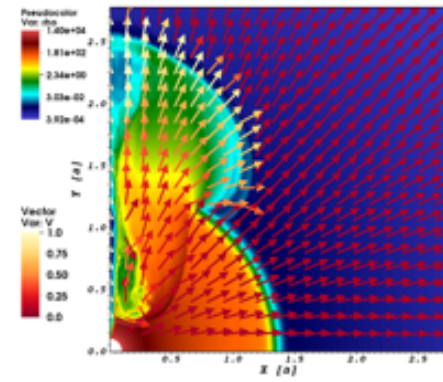
- Second peak from fast spine (not yet seen)

- NS-NS merger: hot disk
- Time to accumulate B-flux on BH  $\sim 1$  sec
- Jet plows through  $\sim 0.01 M_{\text{Sun}}$
- Breakout after  $\sim 1$  sec
- Nearly spherical break-out: prompt



Barkov, Giannios, Luo, Kathirgamaraju, Lyutikov, in prep.

Mildly relativistic shock propagating through  $\rho \sim 10^4 \text{ g cm}^{-3}$  ( $\sim 10^{-2} M_{\odot}$  over  $10^9 \text{ cm}$ )

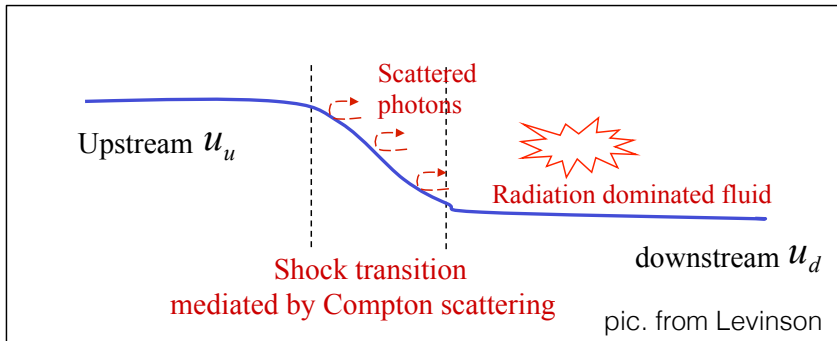


## Radiation-mediated shocks

For  $\beta \geq \mu^{-1/2} (n\lambda_C^3)^{1/6} \approx 10^{-2}$  post-shock radiation pressure > kinetic pressure

$$\mu = m_p/m_e$$

Momentum flux  $\sim$  radiation flux  $\rho v^2 \sim \sigma_{SB} T^4 / c$

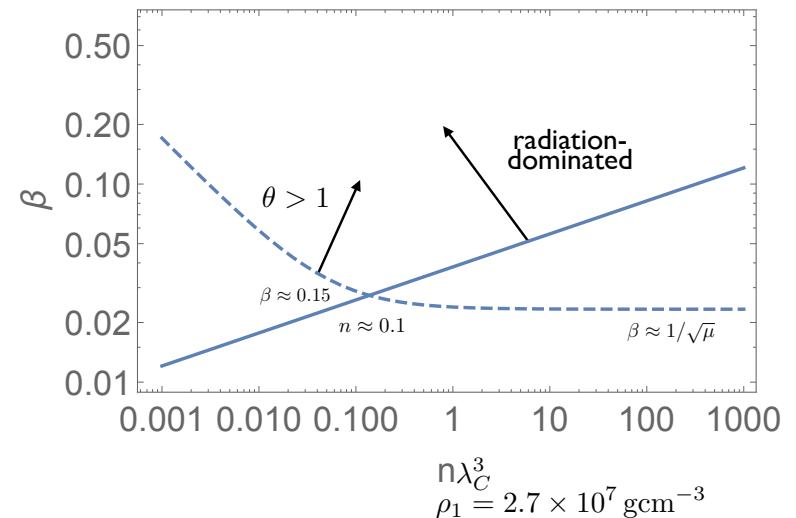


- Hot post-shock fluid emits photons - photon pressure decelerates the flow
- Mildly relativistic flows can be strongly affected by radiation
- High optical depth - LTE (?)

Highly radiation-dominated:

$$\frac{T}{m_e c^2} \sim \mu^{1/4} (n\lambda_C^3)^{1/4} \sqrt{\beta} \sim 1 \text{ for } \beta \geq 0.3$$

Pair production (and nuclear reactions) in the wind



# Resolving radiation and pair-mediated shock transitions

overall jump condition

$$\beta_1 \rho_1 = \beta \rho$$

matter flux

$$\rho_1 \beta_1^2 = p_{tot} + \rho_{tot} \beta^2$$

momentum flux

$$\rho_1 \beta_1^3 / 2 = (w_{tot} + \rho_{tot} \beta^2 / 2) \beta + F_r$$

energy flux

$$F_r = -\frac{c}{3n_{tot}\sigma_T} \nabla u_{rad}$$

Energy redistribution  
by radiation.

$$u_{rad} = \frac{4}{c} \sigma_{SB} T^4$$

Diffusive - approximation!

Pressure, enthalpy, density: sums of baryons, pairs and radiation

Even though the radiation pressure is small, it can fly far-far

Higher order diff. equation - very different structure of solutions

# Very simple case

- Radiation energy density is negligible, but efficient redistribution

$$\beta_1 \rho_1 = \beta \rho$$

$$\rho_1 \beta_1^2 = \frac{\rho}{m_p} T + \rho \beta^2$$

$$\eta = \frac{\rho_1}{\rho} = \frac{v}{v_1}$$

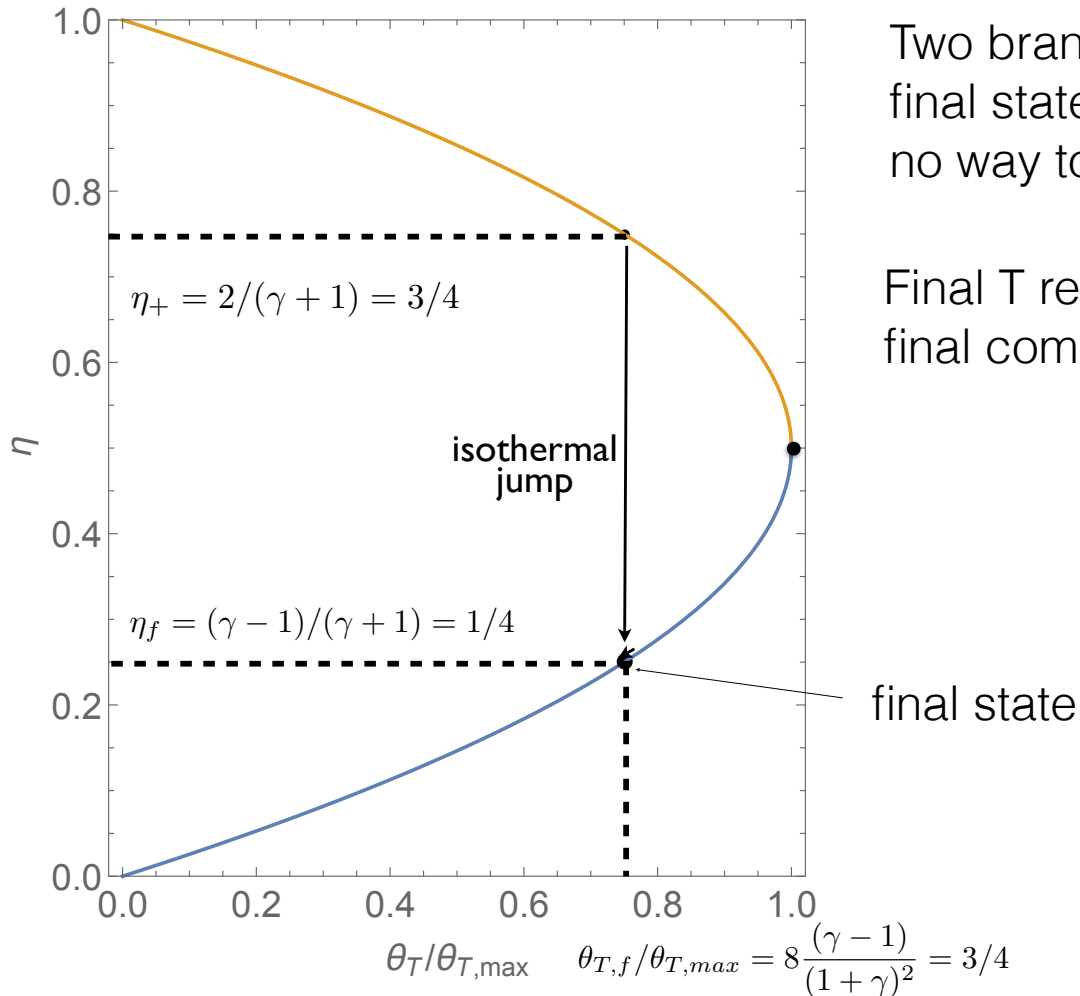
$$p = nT$$

$$\rho_1 \beta_1^3 / 2 = \left( \frac{\gamma}{\gamma - 1} \frac{\rho}{m_p} T + \rho \beta^2 / 2 \right) \beta + F_r$$

$$T = \eta(1 - \eta)m_p v^2$$

$$T = \eta(1 - \eta)m_p v^2$$

Read Zeldovich & Raizer:  
isothermal jump

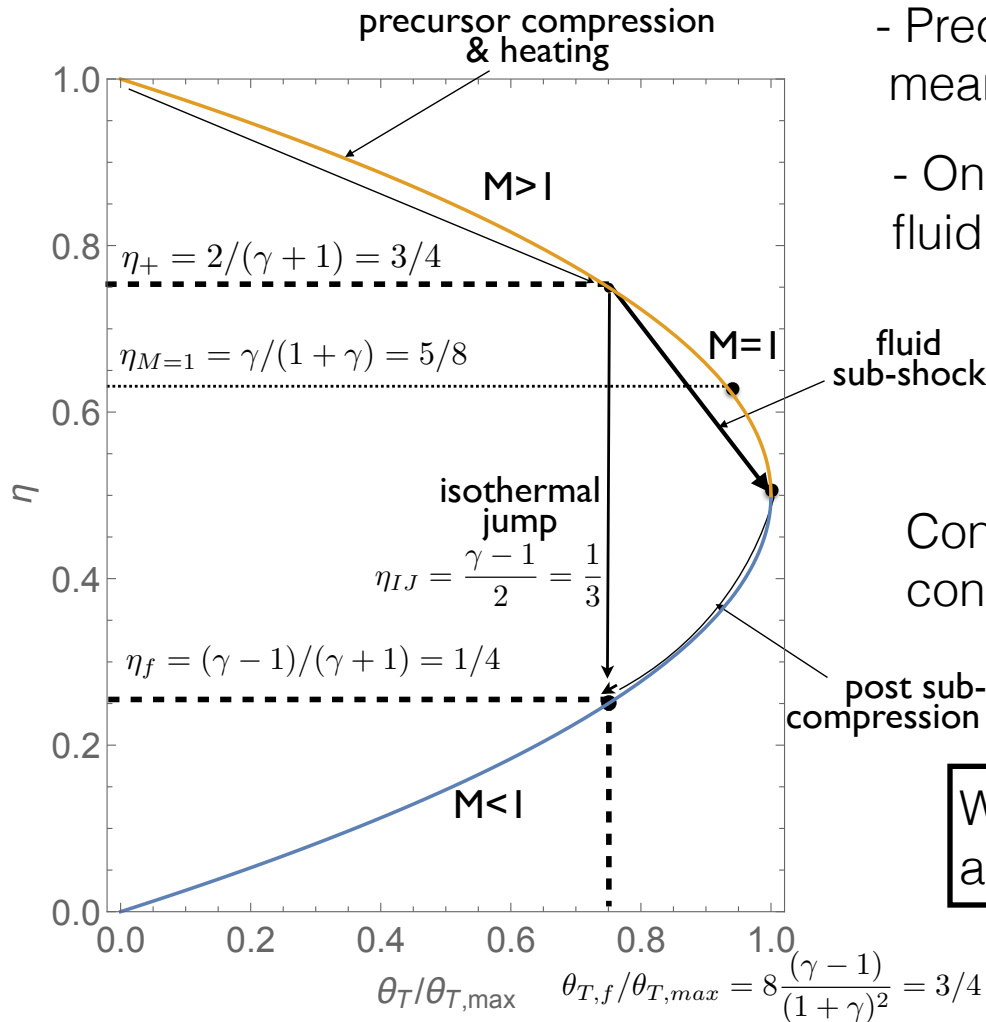


Two branches - initially on upper,  
final state on lower,  
no way to pass throughout

Final T reached before  
final compression

final state

# Fluid subshock



- Precursor: on scales  $\gg$  photon mean free path: slow down and heat-up
- On scales  $\ll$  photon mean free path: fluid subshock, radiation continuous

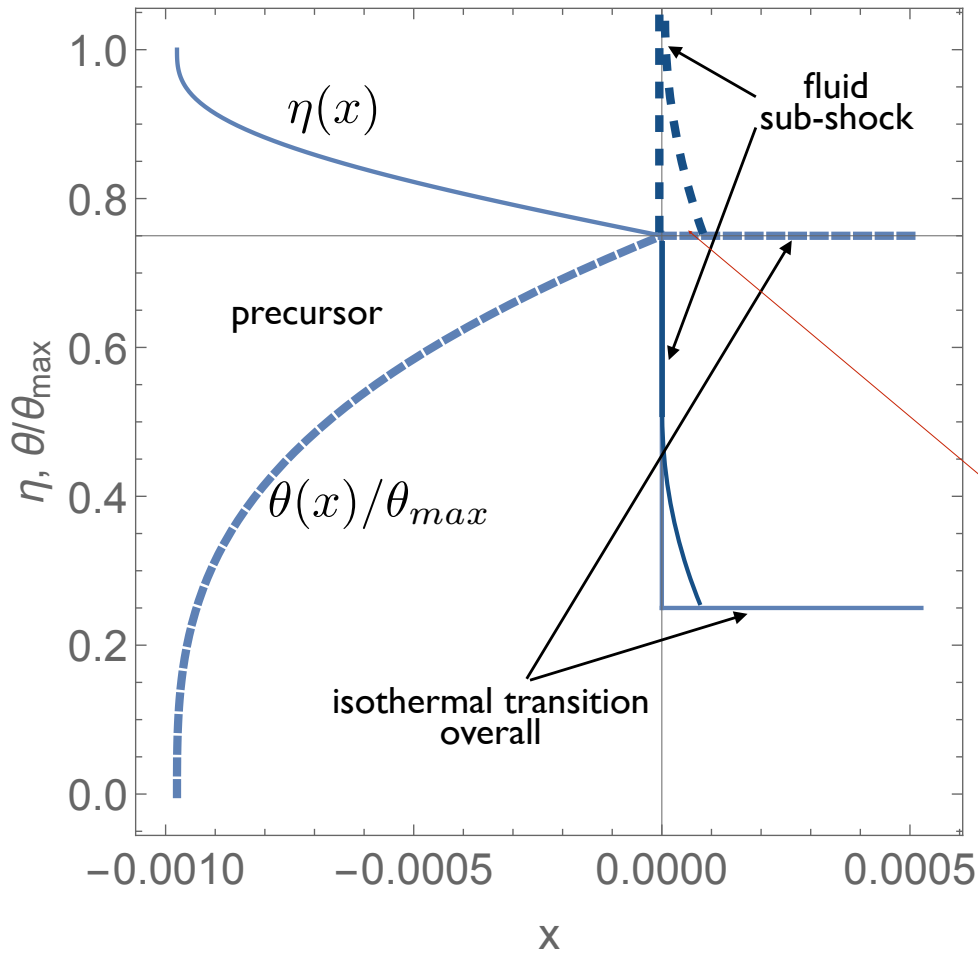
fluid sub-shock  $M_s = \sqrt{\frac{2}{\gamma(\gamma - 1)}} = \frac{3}{\sqrt{5}}$

Continue on momentum conservation curve

We did not say anything about how energy is redistributed!

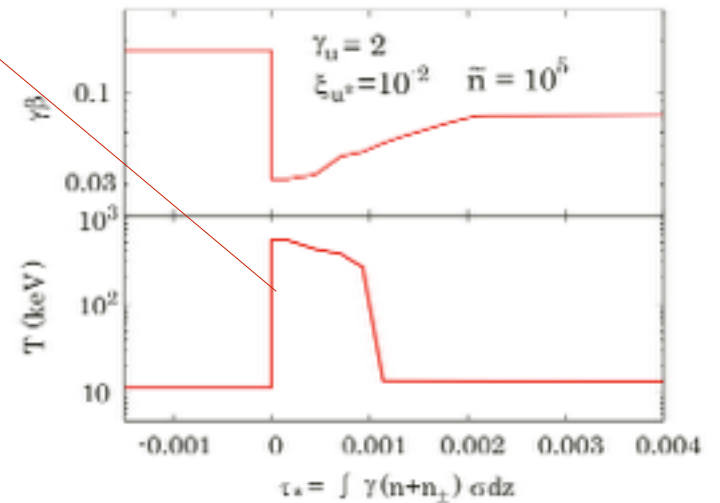


# Resolving the isothermal jump

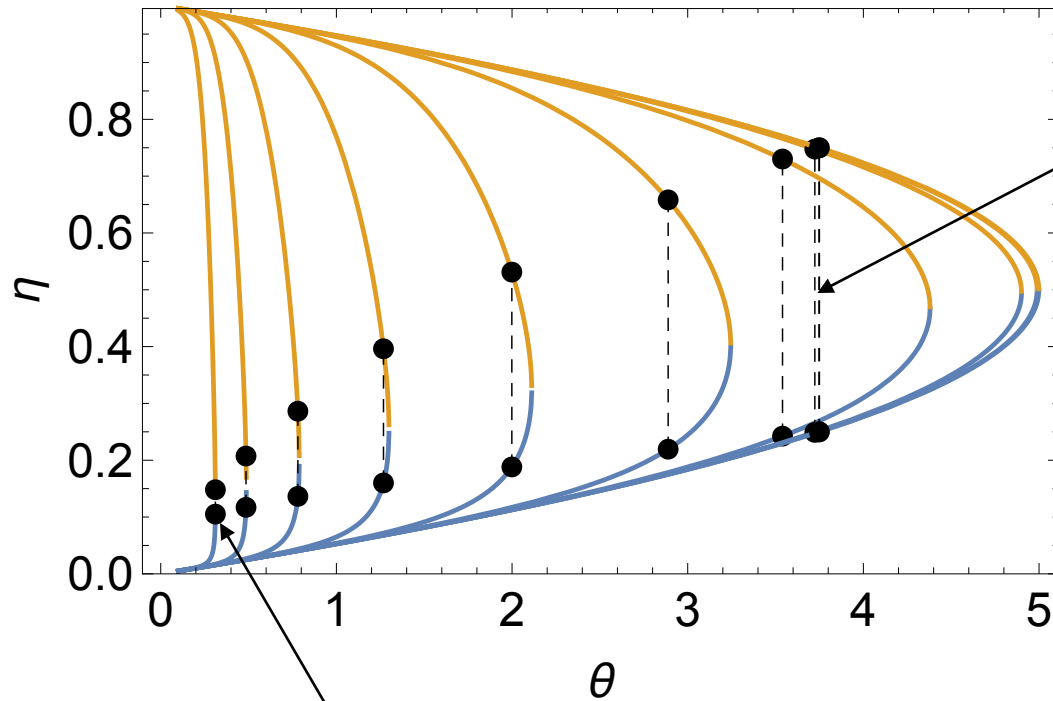


$$M_s = \sqrt{\frac{2}{\gamma(\gamma-1)}} = \frac{3}{\sqrt{5}}$$

Itoh + 2018



# Add pairs and radiation



limit of large density

Standing shocks in core collapse, high density limit: isothermal jump - post-shock  $T$  is 25% higher, but density 30% lower

- For highly radiatively dominated shocks (low density) isothermal jump disappears - no shock, continuous transition (can also be shown analytically)
- This turns out to be the regime in post NS-NS merger winds.

# Conclusion

- Shocks in NS-NS mergers evolve in new, poorly explored regime of mildly relativistic velocities, relativistic temperatures, photon and pair loading, perhaps induced nuclear reactions

