

# Mechanisms for electron acceleration in reconnection and implications for exploring energetic particle production in macro-scale systems

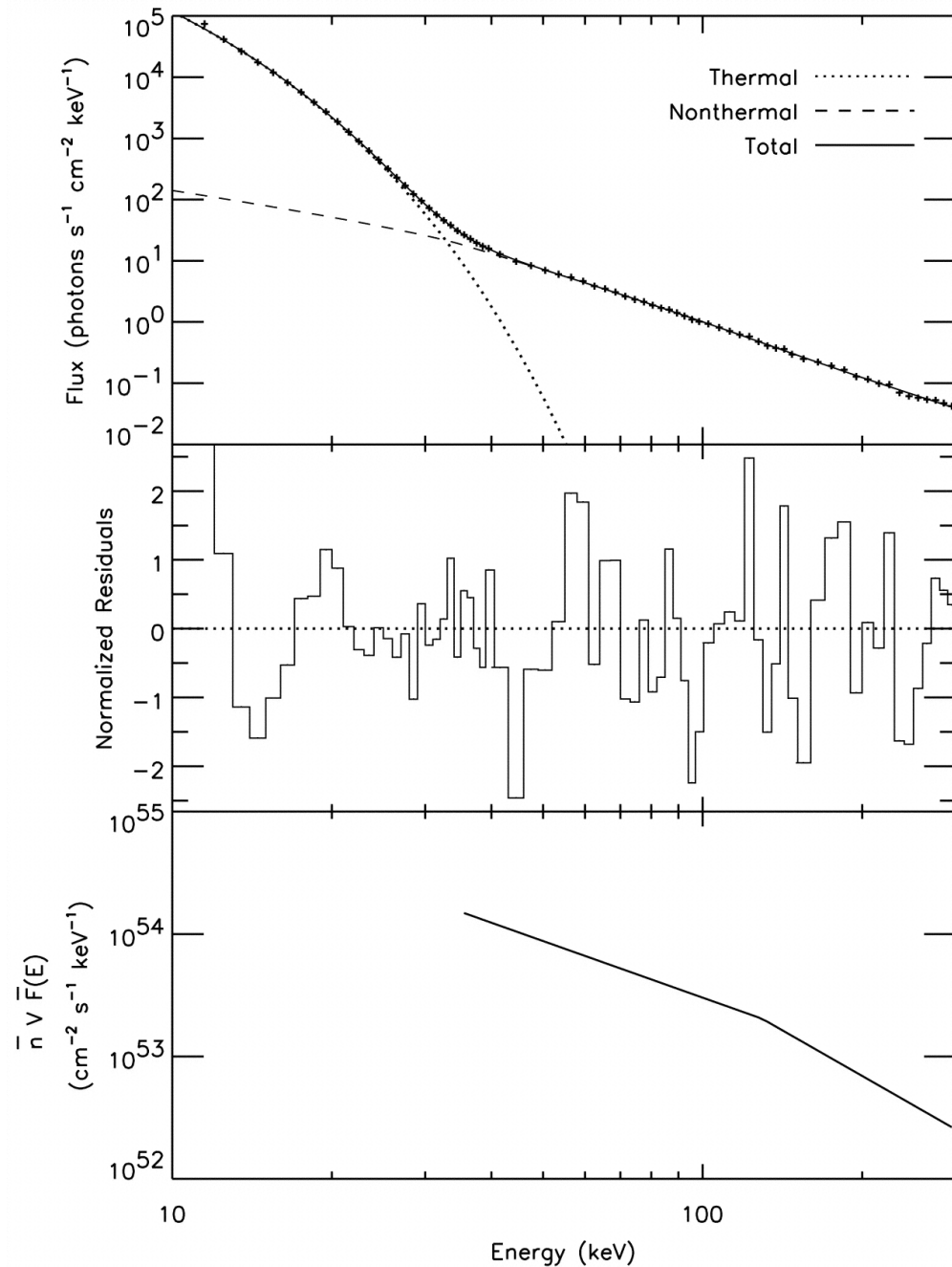
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# RHESSI observations

- July 23  $\gamma$ -ray flare  
(Holman, *et al.*, 2003)
- Double power-law fit with spectral indices:  
1.5 (34-126 keV)  
2.5 (126-300 keV)



# The computational challenge: an enormous separation of scales most astrophysical systems

- The exploration of energetic particle production during reconnection requires a kinetic treatment
  - PIC simulations must resolve a range of kinetic scales
- Solar corona characteristic scale lengths ( $T \sim 100\text{eV}$ ,  $n \sim 10^{10}/\text{cm}^3$ )
  - Debye length  $\sim 0.1\text{cm}$
  - Electron skin depth  $\sim 5\text{cm}$
  - Ion inertial length  $\sim 2.5\text{m}$
  - Coronal x-ray emission region  $\sim 10^4\text{km}$
- Separation of scales  $\sim 10^{10}$
- PIC simulations of flares in the corona are hopeless, even using a PIC code embedded in an MHD model
- Need a new computational model for exploring reconnection-driven particle acceleration in macro-scale systems

# Main Points

- Basic physics of non-relativistic reconnection
  - Heating and particle acceleration are dominated by parallel electric fields and Fermi reflection
  - Fermi reflection dominates the production of energetic particles
    - Not parallel electric fields
  - Particle acceleration is dominated by the dynamics of macro-scale magnetic islands
    - Can dump all kinetic scales
    - The classic picture of dissipation at small spatial scales is invalid
- Minimalist MHD/kinetic model for exploring energetic particles during reconnection in macro-systems
  - Order out all kinetic scales
  - An MHD backbone controls electric and magnetic fields
  - Include guiding center particles on the MHD grid
  - Include the feedback of particles on the MHD dynamics

# Energy release during reconnection

- The change in magnetic topology for reconnection takes place in the “diffusion” region
  - A very localized region around the x-line
  - This is not where significant magnetic energy is released



- Energy release primarily takes place downstream of the x-line where newly-reconnected field lines relax their tension
- Mechanisms for particle heating and energization can not be localized in the “diffusion region”

# Basic mechanisms for particle energy gain during reconnection

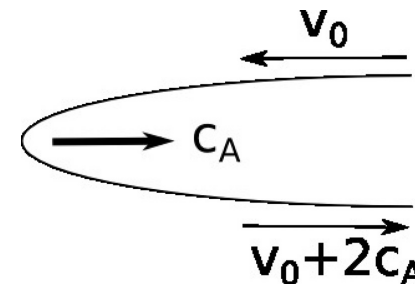
- In the guiding center limit

$$\frac{d\varepsilon}{dt} = qv_{\parallel}E_{\parallel} + q\vec{v}_c \cdot \vec{E} + \mu \frac{\partial B}{\partial t} + q\vec{v}_B \cdot \vec{E}$$

- Curvature drift

- Slingshot term (Fermi reflection) increases the parallel energy

$$v_c = \frac{v_{\parallel}^2}{\Omega} \vec{b} \times (\vec{b} \cdot \vec{\nabla} \vec{b})$$



- Grad B drift

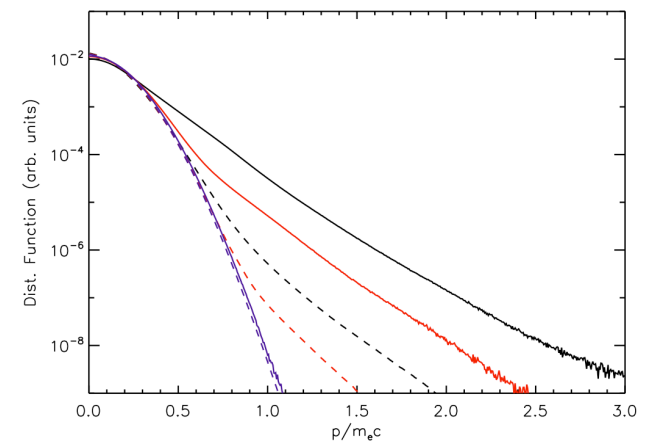
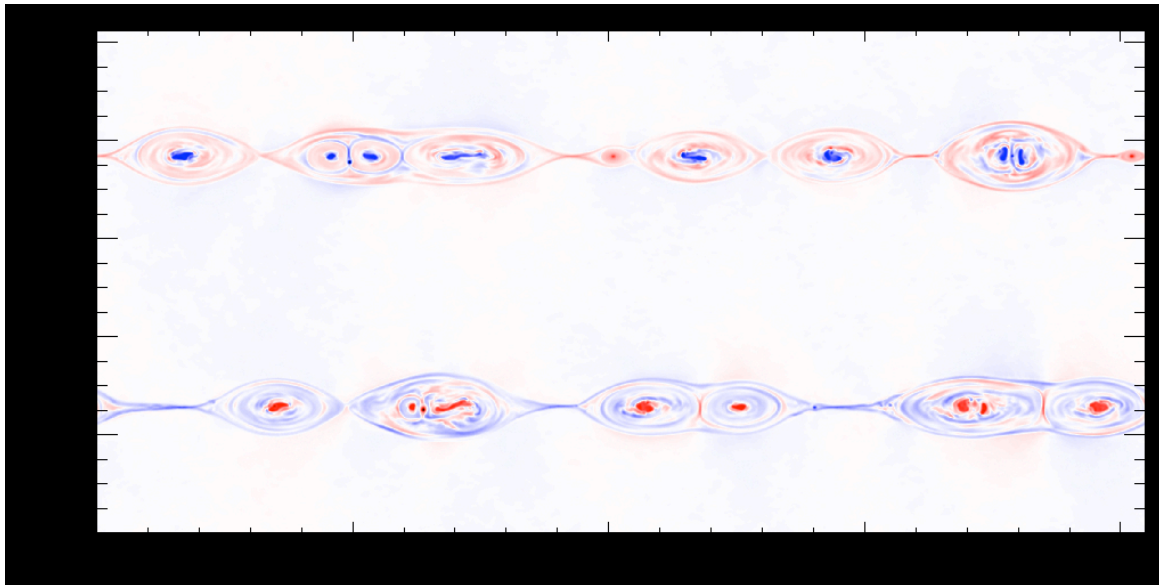
- Betatron acceleration increases perpendicular energy –  $\mu$  conservation

$$v_B = \frac{v_{\perp}^2}{2\Omega} \vec{b} \times \frac{\vec{\nabla} B}{B} \quad \mu = \frac{mv_{\perp}^2}{2B}$$

# Electron heating during reconnection

- Carry out PIC simulations of electron-proton system with a range guide fields
- Focus here on 2D --  $819.2d_i \times 409.6d_i$ 
  - Compare all of the heating mechanisms
  - Dahlin et al '14
- Do not see powerlaws in energetic particles
- Strong anisotropy – dominant acceleration parallel to the magnetic field

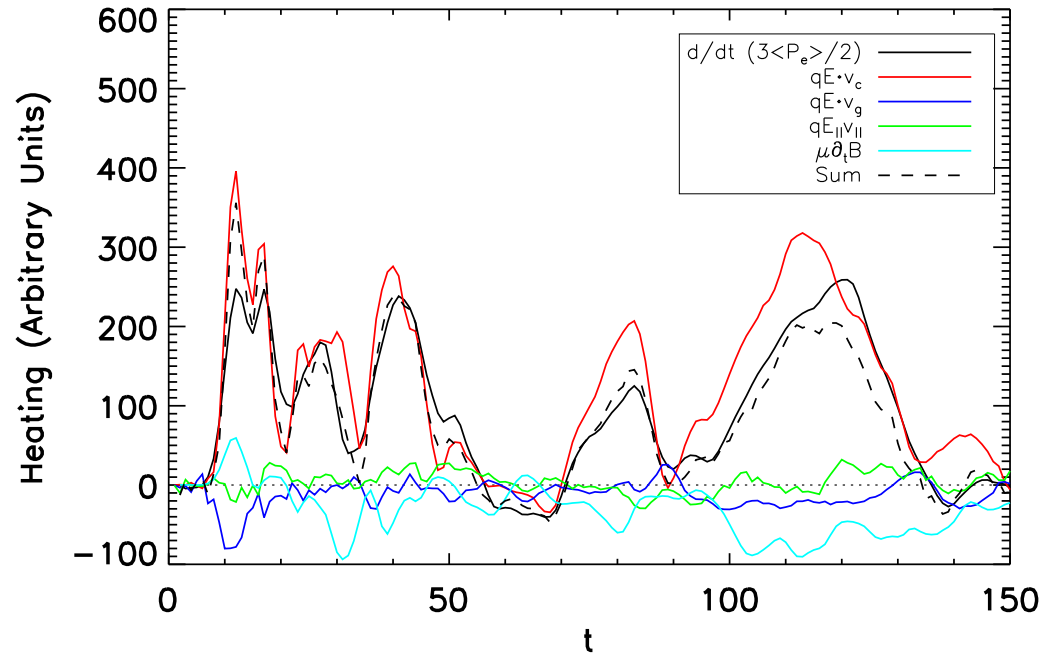
$$d_i = \frac{c}{\omega_{pi}}$$



# Electron heating mechanisms: weak guide field

- Slingshot term dominates (Fermi reflection)
- Parallel electric field term small – a surprise
- Grad B term is an energy sink
  - Electrons entering the exhaust where B is low lose energy because  $\mu$  is conserved.

$$B_g = 0.2B_r$$



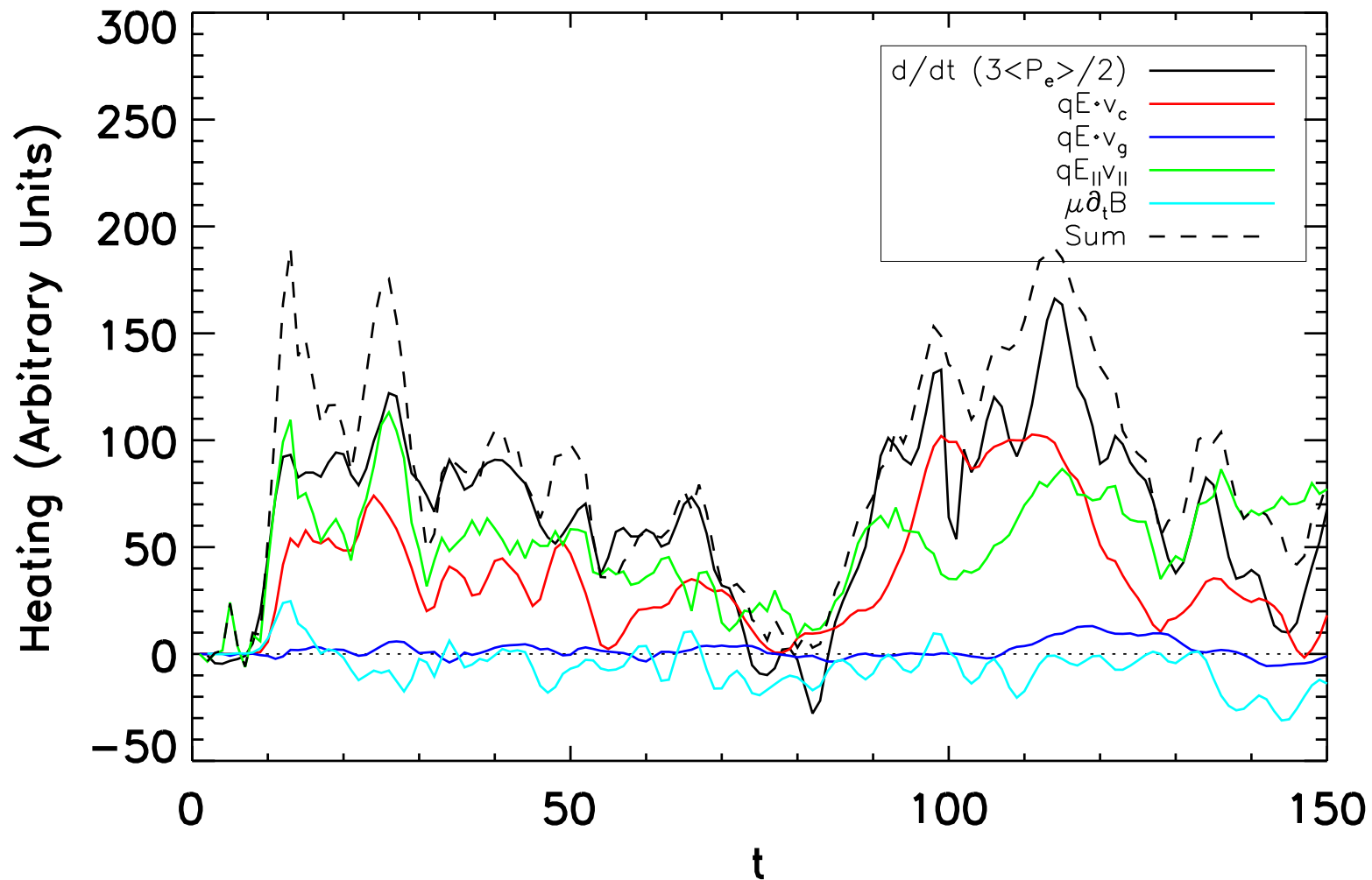


# Electron heating mechanisms: strong guide field

- Fermi and parallel electric field term dominate

– Longer current layers where  $E_{\parallel} \neq 0$  with a guide field

$$B_g = 1.0 B_r$$



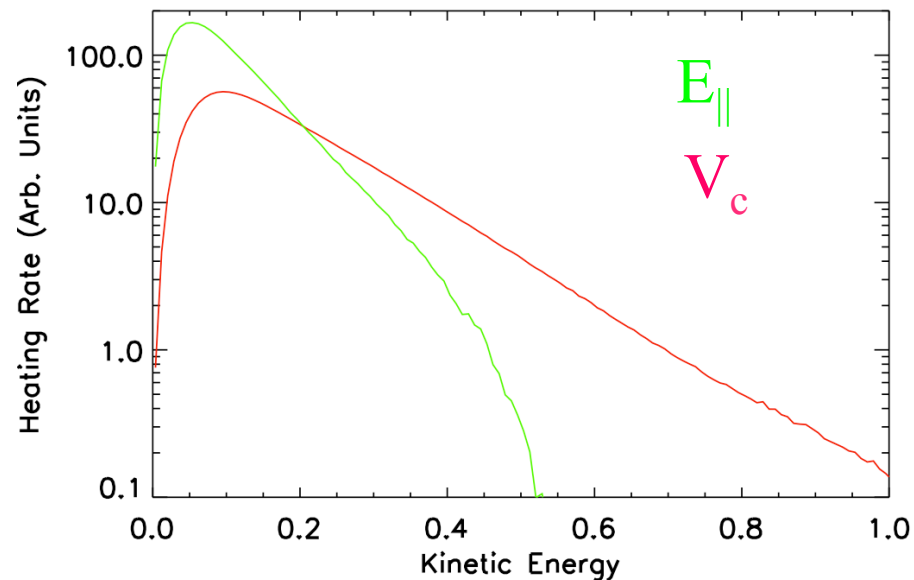
# Acceleration mechanism for highest energy electrons

- Fermi reflection dominates energy gain for highest energy electrons

– Where  $v_c \sim v_{\parallel}^2$

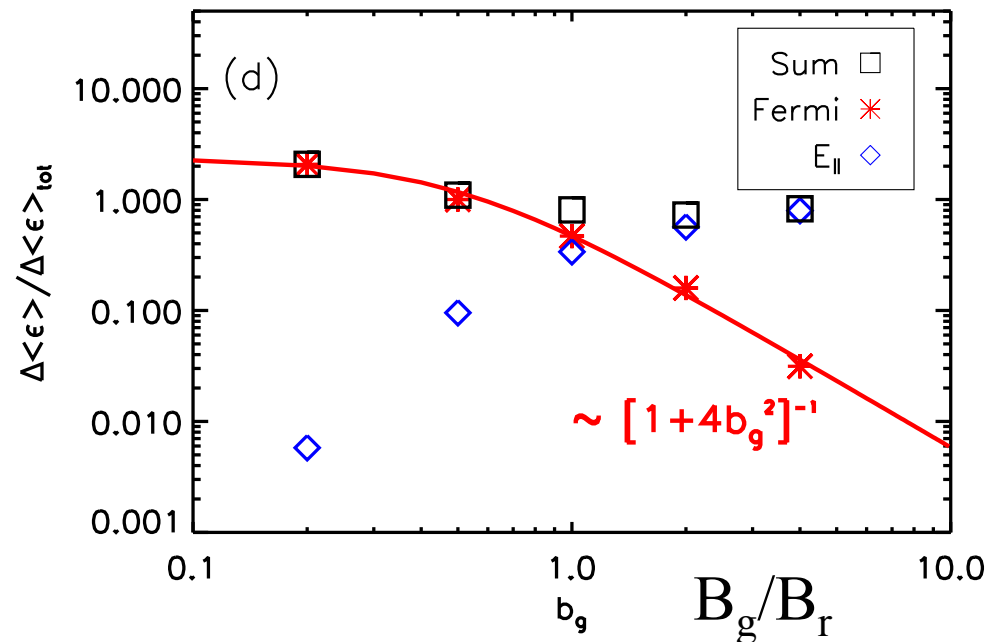
$$\frac{d\varepsilon}{dt} \sim qv_{\parallel}E_{\parallel} + q\vec{v}_c \cdot \vec{E}$$

- Recent simulations of pair and relativistic reconnection also see the dominance of Fermi reflection (Guo et al '14, Sironi and Spitkovsky '14)



# Electron heating: dependence on the guide field

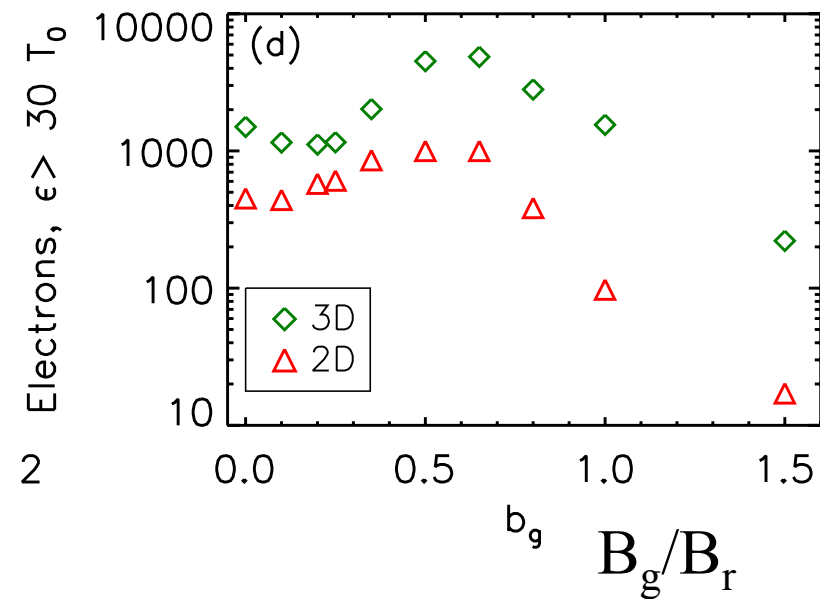
- Fermi reflection dominates for weak guide field
- $E_{\parallel}$  dominates for strong guide field
  - Consistent with gyro-kinetic ordering



Dahlin et al '16

# Production of energetic electrons: versus guide field

- Compare the production of energetic electrons versus the strength of guide field
  - Weak to modest guide field Fermi dominates
  - Large guide field  $E_{\parallel}$  dominates
- Virtually no energetic particles produced in strong guide field reconnection
- Parallel electric fields are not the driver of the most energetic electrons

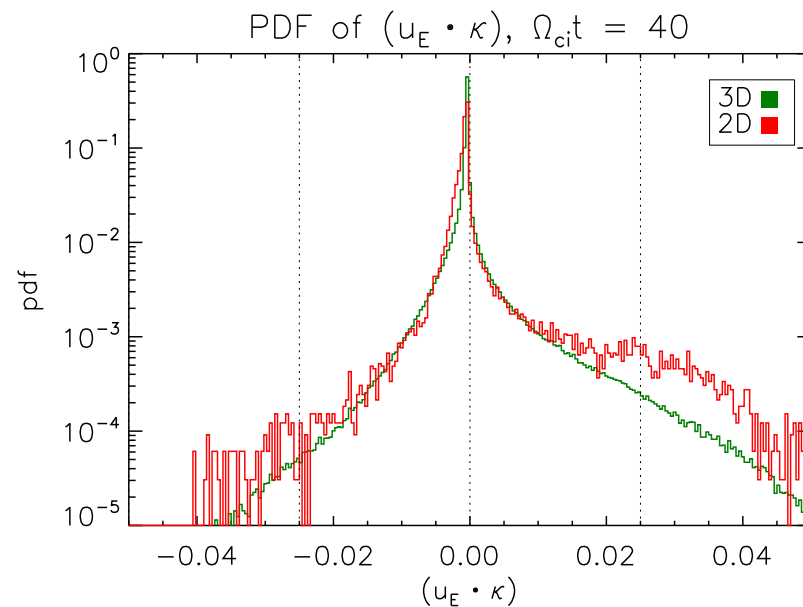
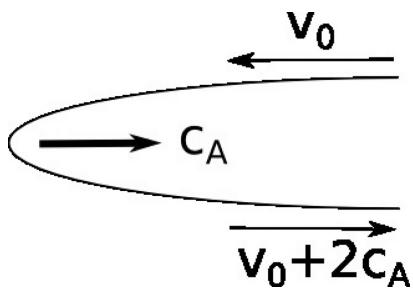


# A measure of particle acceleration efficiency

- A measure of the rate of energy release and particle acceleration during reconnection is the parameter

$$\vec{\kappa} \cdot \vec{V}_{ExB} = (\vec{b} \cdot \nabla \vec{b}) \cdot \frac{c\vec{E} \times \vec{B}}{B^2}$$

- Dominantly positive in a reconnecting system and negative in a dynamo systems
- The dominance of positive values in a reconnecting systems establishes that particle acceleration is a first order Fermi mechanism



# Spatial scales of energy dissipation in reconnection

- The canonical picture in MHD turbulence is that energy cascades to small spatial scales, where dissipation takes place
  - What about during reconnection in collisionless plasma?

- Parallel electric field  $\dot{\epsilon}_{\parallel} = J_{\parallel} E_{\parallel}$ 
  - Linked to the electron skin depth  $c/\omega_{pe}$

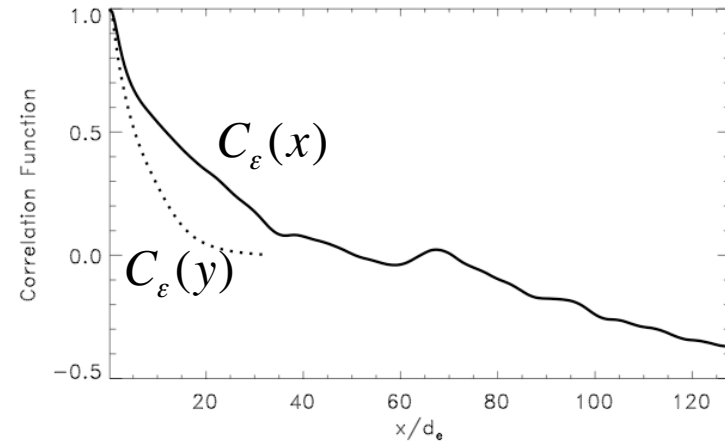
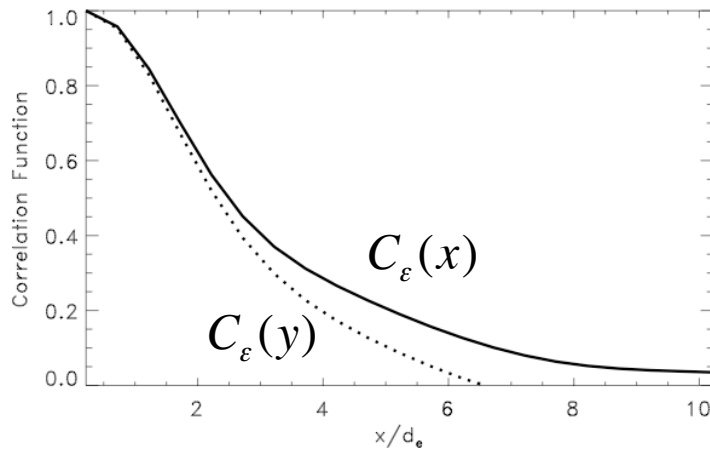
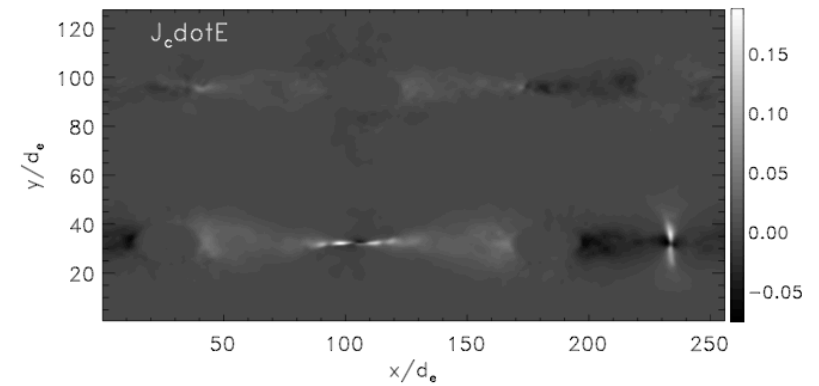
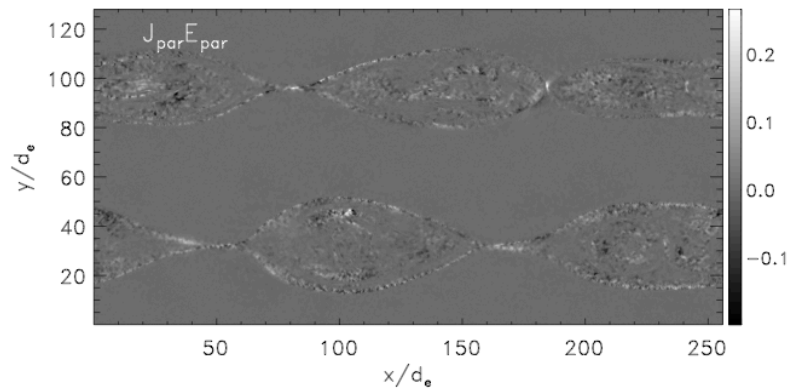
- Fermi reflection  $\dot{\epsilon}_c = \vec{J}_c \cdot \vec{E}$ 
  - Not linked to any kinetic scale

$$\vec{J}_c = -ne\vec{v}_c = -ne \frac{v_{\parallel}^2}{\Omega} \vec{b} \times (\vec{b} \cdot \vec{\nabla} \vec{b})$$

# Correlation scales of electron dissipation

- Calculate the correlation function of electron dissipation from reconnection simulations

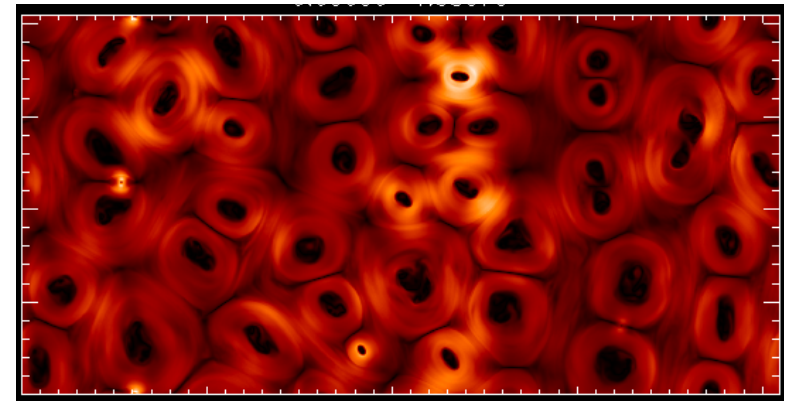
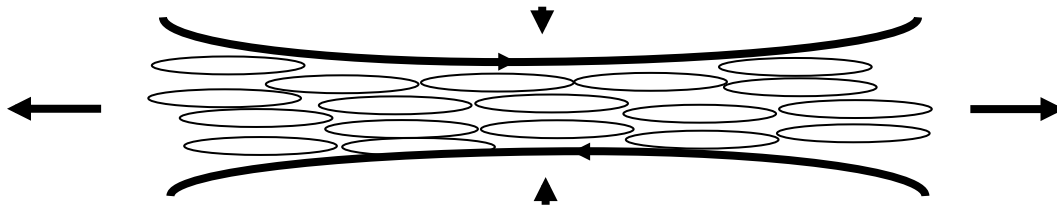
$$C_\varepsilon(\vec{r}) = \langle \dot{\varepsilon}(\vec{x} + \vec{r}) \dot{\varepsilon}(\vec{x}) \rangle$$

 $\dot{\varepsilon}_\parallel$ 
 $\dot{\varepsilon}_c$ 


# Particle acceleration in multi-island reconnection

- Single x-line reconnection can not explain the most energetic particles seen in solar flares
  - Energy gain limited to around 10keV
- Greater energy gain in contracting and merging magnetic islands

Tajima and Shibata '97  
Drake et al '06  
Oka et al '10



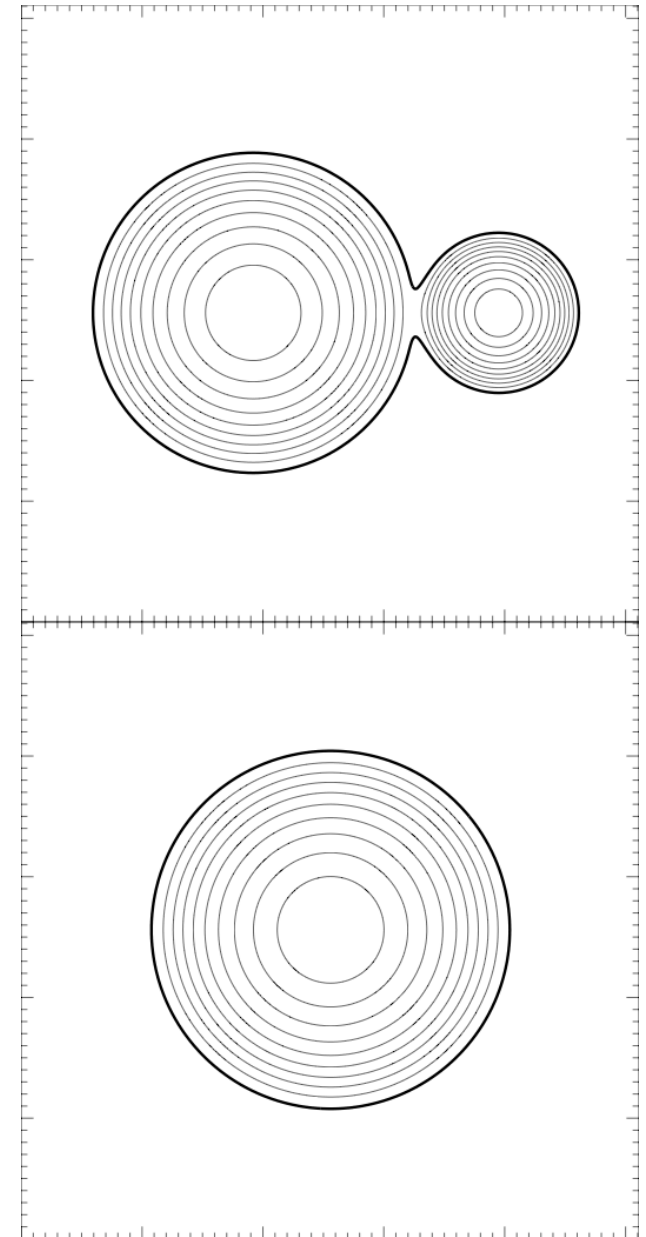
$C_{Ax}$



# Energy gain in a bath of merging islands

- Total area preserved during merger
- Magnetic flux of largest island is preserved
- Merging islands shorten field lines
- Parallel action is conserved  $p_{\parallel} L$ 
  - $L$  goes down during merger so  $P_{\parallel}$  goes up
  - Not valid for relativistic reconnection
- The merger of two equal sized islands doubles the parallel energy of particles within the islands

Drake et al '13, Montag et al '17



# MeV electrons in a coronal hard x-ray source

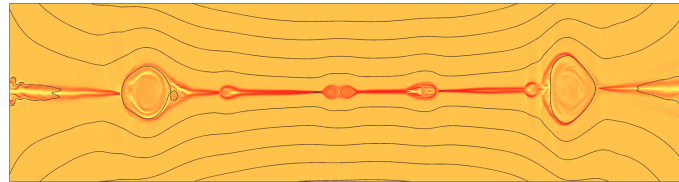
- How to get MeV electrons in the corona?
  - A two-step process – heating in single x-line reconnection following by island merging
- First step: single x-line reconnection splits released energy between electrons, ions and bulk flow
  - $\beta_e \sim 1/4$
  - For  $B \sim 50\text{G}$ , with  $n \sim 10^9\text{cm}^{-3}$ , obtain  $T_{\text{hot}} \sim 15\text{keV}$
- Second step: island mergers
  - Each merger doubles the electron energy – field line shortening
  - How many island mergers takes 10keV electrons to 1MeV?

$$15\text{keV} \times 2^N = 1\text{MeV} \Rightarrow N = 6$$

- Take typical island of size  $W \sim 10^3\text{km}$
- Two island merging time  $t_{\text{merge}} \sim (W/2)/0.1c_A \sim 1.5\text{s}$
- 1MeV electrons in  $t_{1\text{MeV}} \sim 6t_{\text{merge}} = 9\text{s}$

# Modeling reconnection-driven particle acceleration in macro-scale systems

- Eliminate all kinetic scales
  - Kinetic scales are needed to accurately describe parallel electric fields
  - They don't control the production of the most energetic particles
  - Particle production controlled by the dynamics of macro-islands
- Multi-island reconnection in MHD models produces fast reconnection (Bhattacharjee+ '09, Cassak+ '09, Huang+ '10)



- Test particle modeling in MHD fields (Onofri+ '06, Kowal+ '11, Guidoni+ '16)
  - Energy in energetic particles can run away
  - Need feedback of particles on MHD fields, e.g., firehose condition (Drake+ '13)

# A self-consistent MHD/guiding-center kinetic model

- A model with an MHD backbone and with macro-particles evolved in parallel using the guiding center equations
  - All kinetic scales ordered out
  - Energetic component evolved in the MHD fields
  - Energetic particle feedback on the MHD fluid through the pressure driven currents of the energetic component
  - Total energy of the MHD system plus the energetic component is conserved
  - Not appropriate for shock acceleration where particle scattering is essential
- Similar models
  - Park et al '92, gyrokinetic description – no Fermi contributions
  - Bai et al '15, full kinetic description for shocks

# Basic equations

- MHD momentum equation with MHD pressure  $P$  and the energetic particle current  $\mathbf{J}_h$

$$m_i n \frac{d}{dt} \vec{v} = \frac{1}{c} \vec{J} \times \vec{B} - \vec{\nabla} P - \frac{1}{c} \vec{J}_h \times \vec{B}$$

- Order out the Hall terms  $\vec{E} = -\frac{1}{c} \vec{v} \times \vec{B}$

- Particles – guiding center equations

$$\frac{d}{dt} p_{\parallel e} = \frac{c}{B} p_{\parallel e} \vec{b} \times \vec{k} \cdot \vec{E} - \frac{\mu_e}{\gamma_e} \vec{b} \cdot \vec{\nabla} B$$

$$\mu = \frac{p_{\perp e}^2}{2m_e B} = \text{const.}$$

# Basic Equations (cont.)

- Particle Moments

$$\vec{P}_h = P_{\parallel h} \vec{b}\vec{b} + P_{\perp h} (\vec{I} - \vec{b}\vec{b}) \quad \vec{\nabla} \cdot \vec{P}_h = \frac{1}{c} \vec{J}_h \times \vec{B}$$

$$\vec{J}_h = \frac{c}{B} \left( P_{\parallel h} \vec{b} \times \vec{\kappa} + P_{\perp h} \vec{b} \times \vec{\nabla} \ln(B) \right) - c \vec{\nabla} \times \left( \frac{P_{\perp h}}{B} \vec{b} \right)$$

curvature drift    grad B drift    magnetization current

- Energy Conservation

$$\frac{d}{dt} W_{MHD} = - \int d\vec{x} \vec{J}_h \cdot \vec{E} = - \frac{d}{dt} W_h$$

$$\frac{d}{dt} W_h = \int d\vec{x} \left( P_{\parallel h} \vec{V}_E \cdot \vec{\kappa} + P_{\perp h} \vec{V}_E \cdot \vec{\nabla} \ln B + \frac{P_{\perp h}}{B} \frac{\partial}{\partial t} B \right)$$

$$\vec{V}_E = \frac{c}{B^2} \vec{E} \times \vec{B}$$

# Main Points

- Solar observations suggest that magnetic energy conversion into energetic electrons is extraordinarily efficient
- Fermi reflection the dominant driver of energetic electron production during reconnection
  - $E_{\parallel}$  is not the main driver of the energetic component
  - Fermi reflection is dominated by the dynamics of island formation and merger
- Energetic particle production can be modeled without including the kinetic boundary layers necessary for describing  $E_{\parallel}$
- A new hybrid model with an MHD backbone and but with macro-particles evolved in parallel using guiding center equations is being developed
  - The particles self-consistently feed back on the MHD fluid through their perpendicular currents
  - The total fluid plus energetic particle energy is conserved

## Main Points (cont.)

- The particle-in-cell code *p3d* is being modified to advance the new hybrid equations
  - *p3d* is already set up to evolve the MHD equations
  - The coupling of the MHD equations to the particle stepping algorithm is now undergoing development
- Stay tuned for test results of the new model