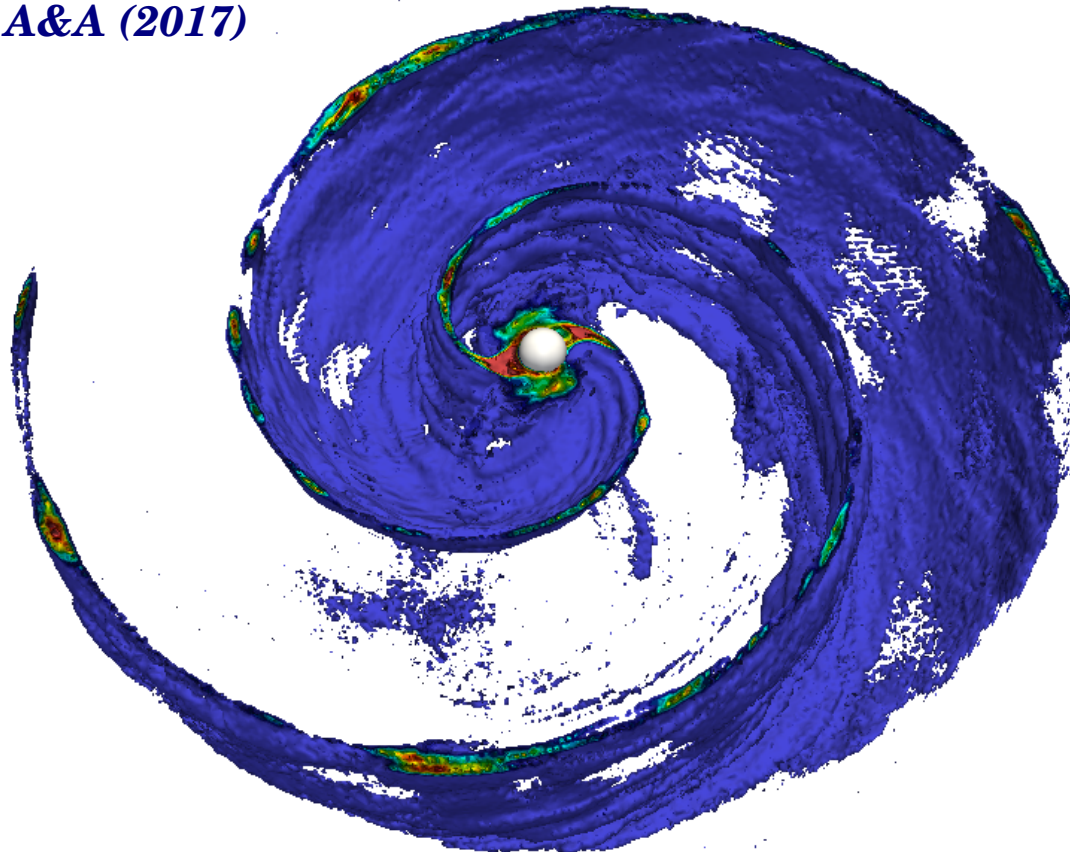


Magnetic dissipation in pulsar winds

Benoît Cerutti, *CNRS & Univ. Grenoble Alpes, France.*

In collaboration with **Sasha Philippov**, *UC Berkeley, USA.*

Cerutti & Philippov, A&A (2017)



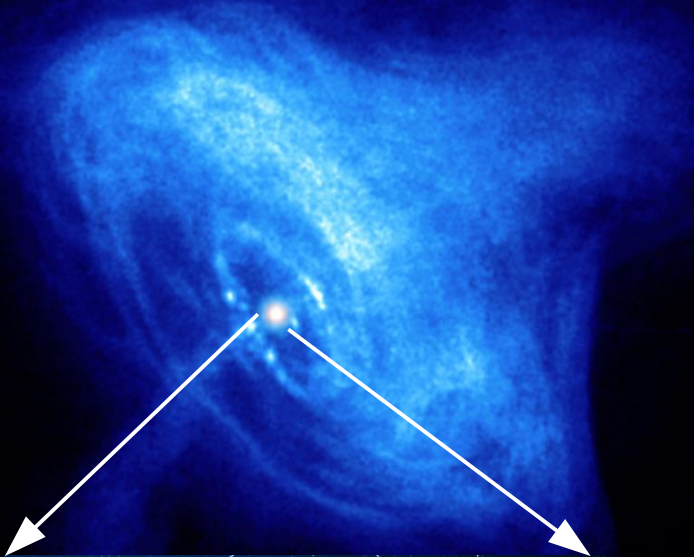
The “sigma” problem

e.g., Rees & Gunn 1974
Kennel & Coroniti 1984a,b
Begelman & Li 1992
Coroniti 1990
Komissarov & Lyubarsky 2004
Lyutikov 2010
Porth+2014 ...

$$\sigma = \frac{B^2}{8 \pi \Gamma n m c^2}$$

Crab

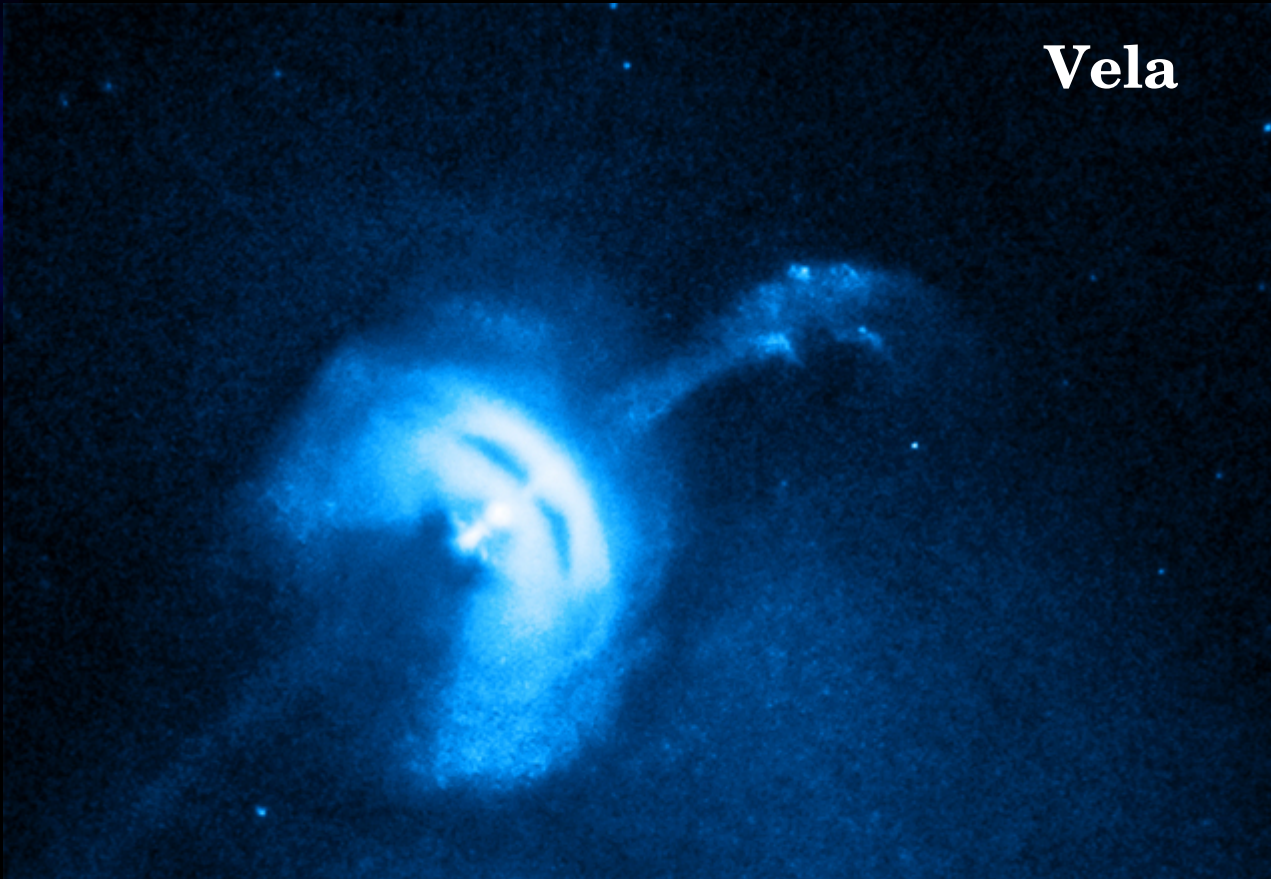
Nebula: $\sigma \ll 1$



Magnetosphere: $\sigma \gg 1$

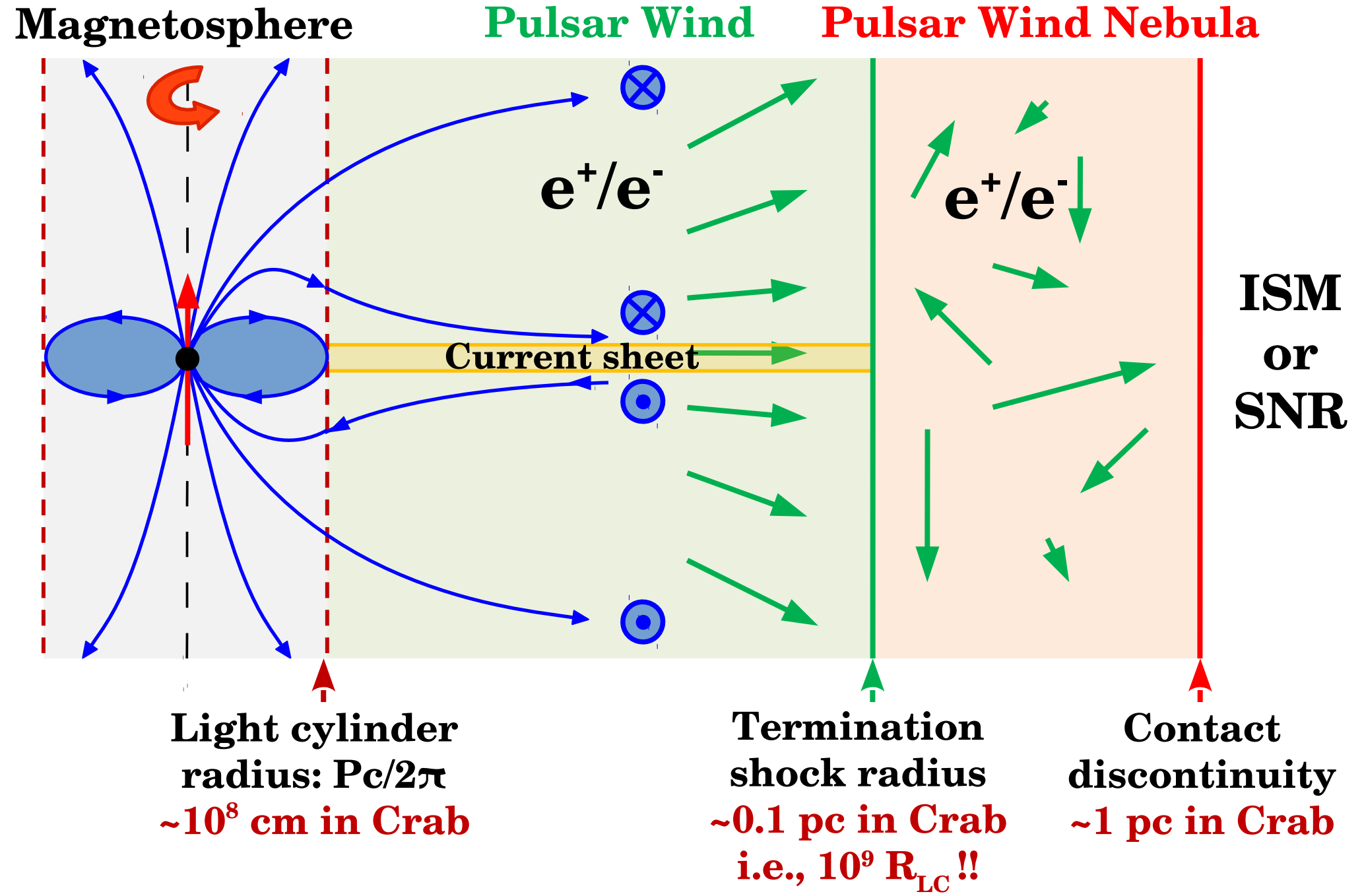


Vela



How and where is the magnetic energy dissipated ?
Generic problem in magnetized astrophysical outflows

The classical picture of pulsar wind nebulae

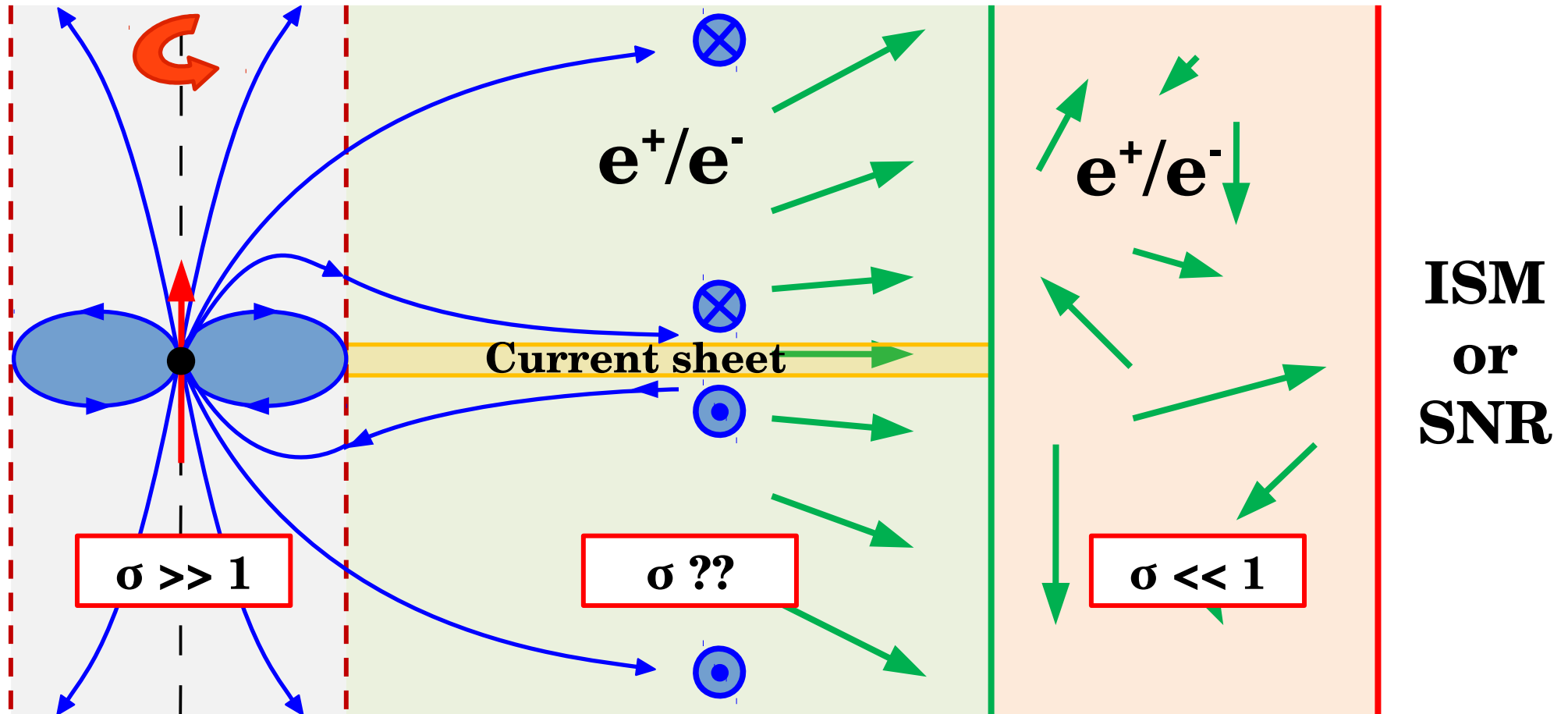


The σ -problem

Magnetosphere

Pulsar Wind

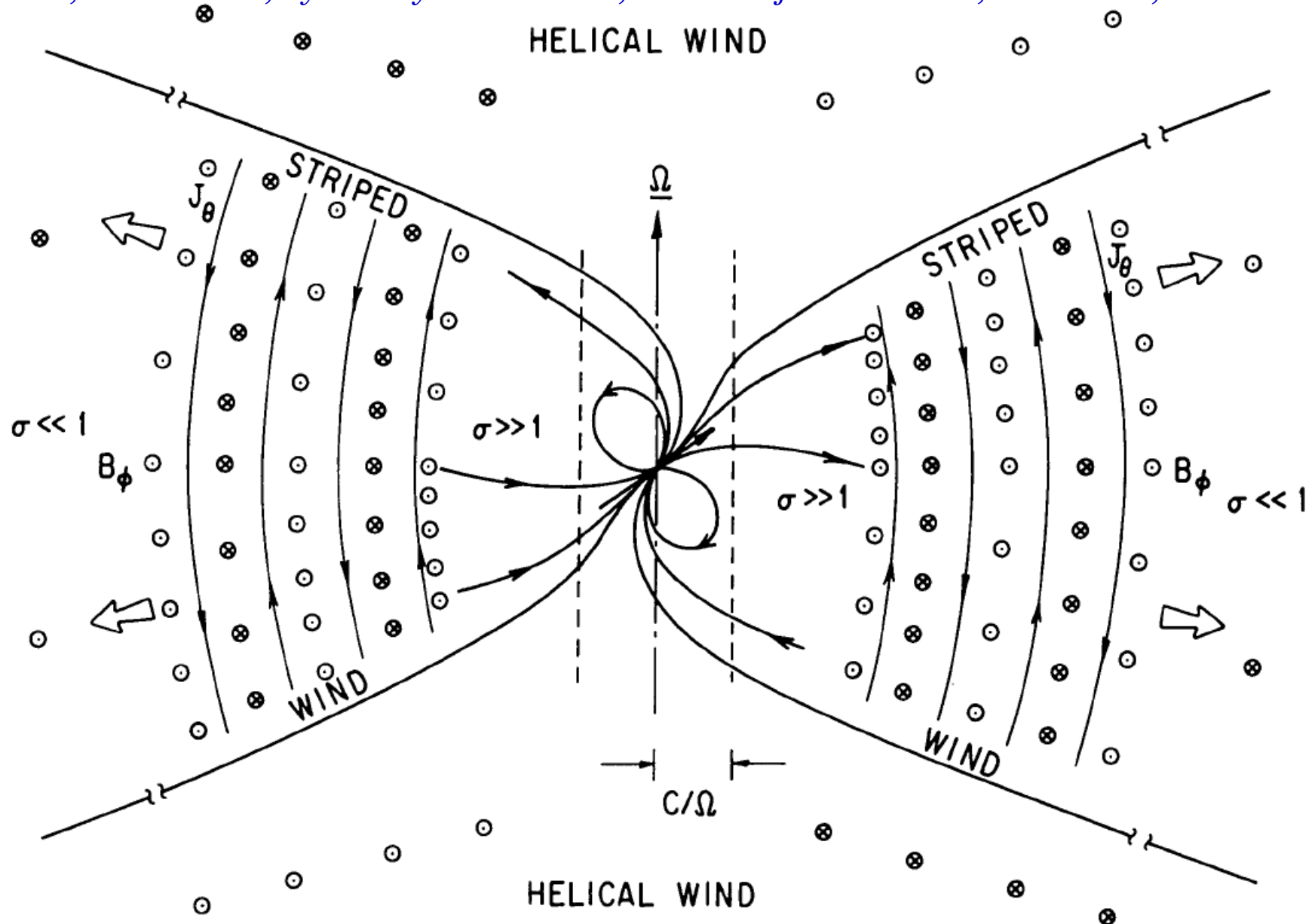
Pulsar Wind Nebula



Transition $\sigma \gg 1$ to $\sigma \ll 1$ unknown: “sigma” problem
=> Dissipation somewhere in between needed!

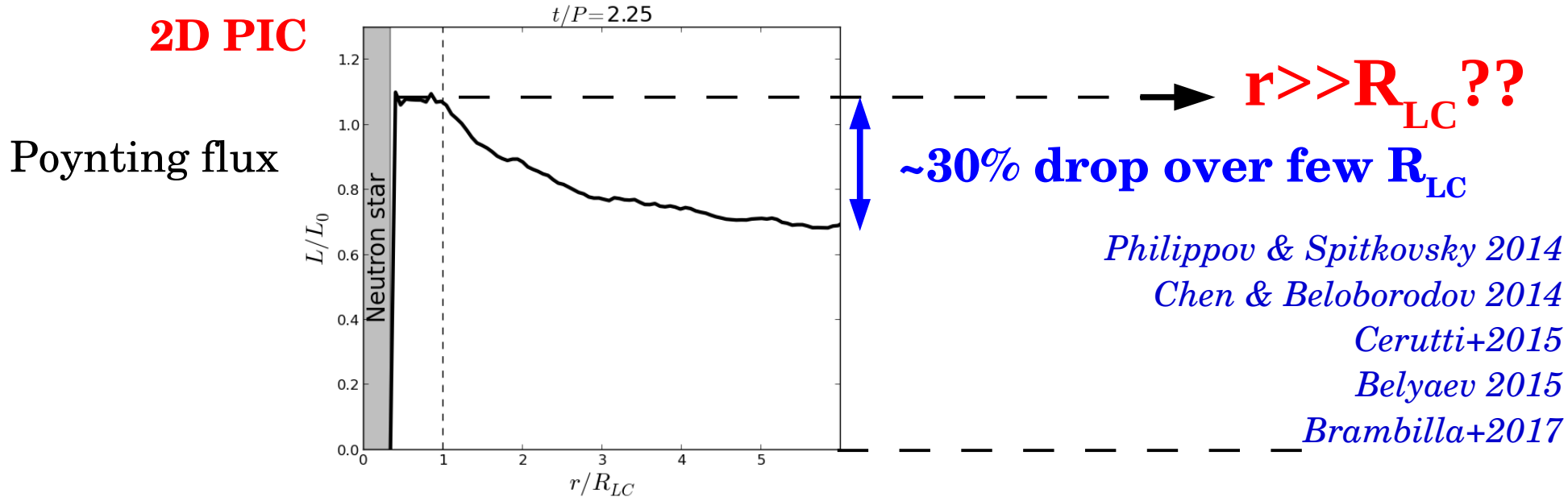
The equatorial current sheet : The ideal culprit !

Coroniti 1990 ; Michel 1994 ; Lyubarsky & Kirk 2001 ; Kirk & Skjæraasen 2003, Zrake 2016, Zrake & Arons 2016



How **far** and how **fast** does magnetic reconnection proceed in the wind?

Close in or at the termination shock ?

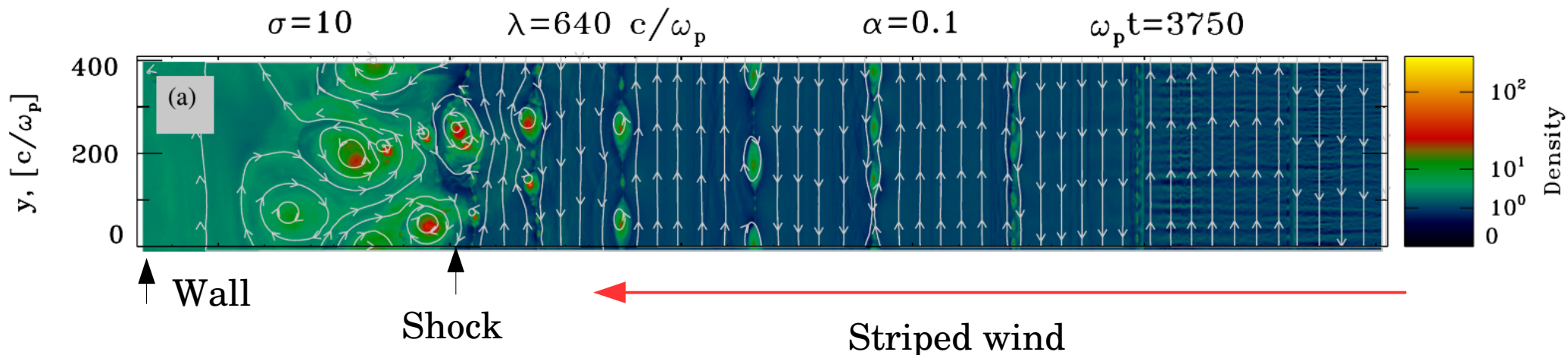


Scenario 1 : Complete dissipation far before the shock

Coroniti 1990 ; Michel 1994 ; Lyubarsky & Kirk 2001 ; Kirk & Skjæraasen 2003

Scenario 2 : Shock-driven reconnection at the termination shock

Lyubarsky 2003 ; Pétri & Lyubarsky 2007 ; Sironi & Spitkovsky 2011

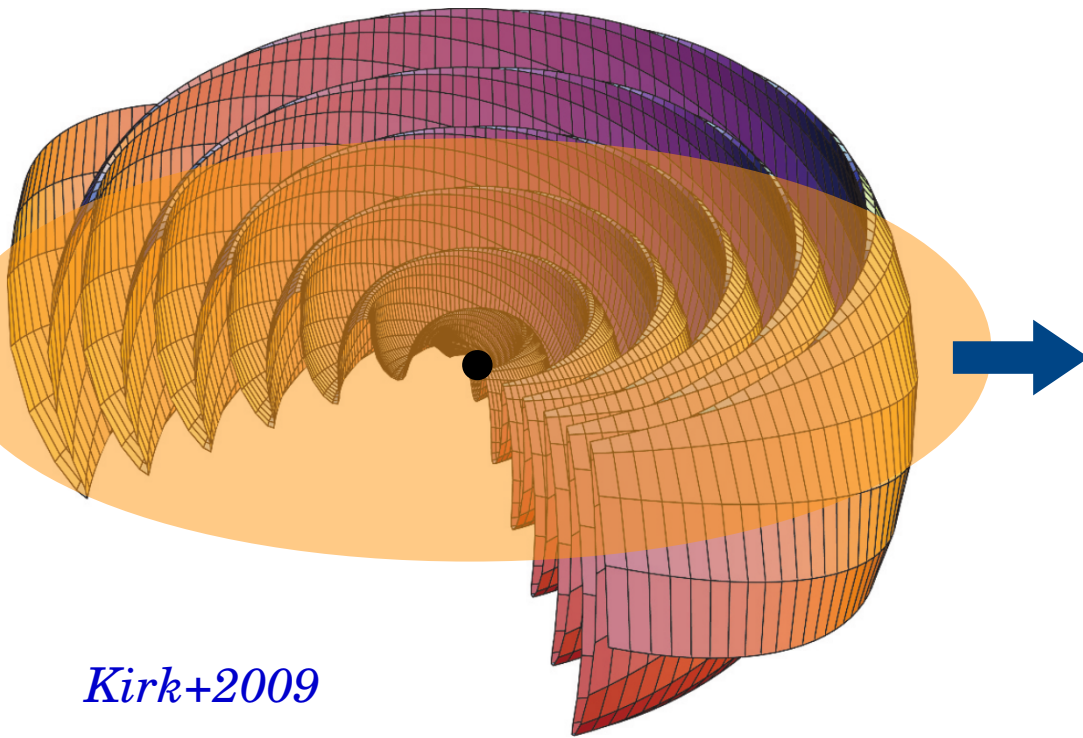


The numerical challenge

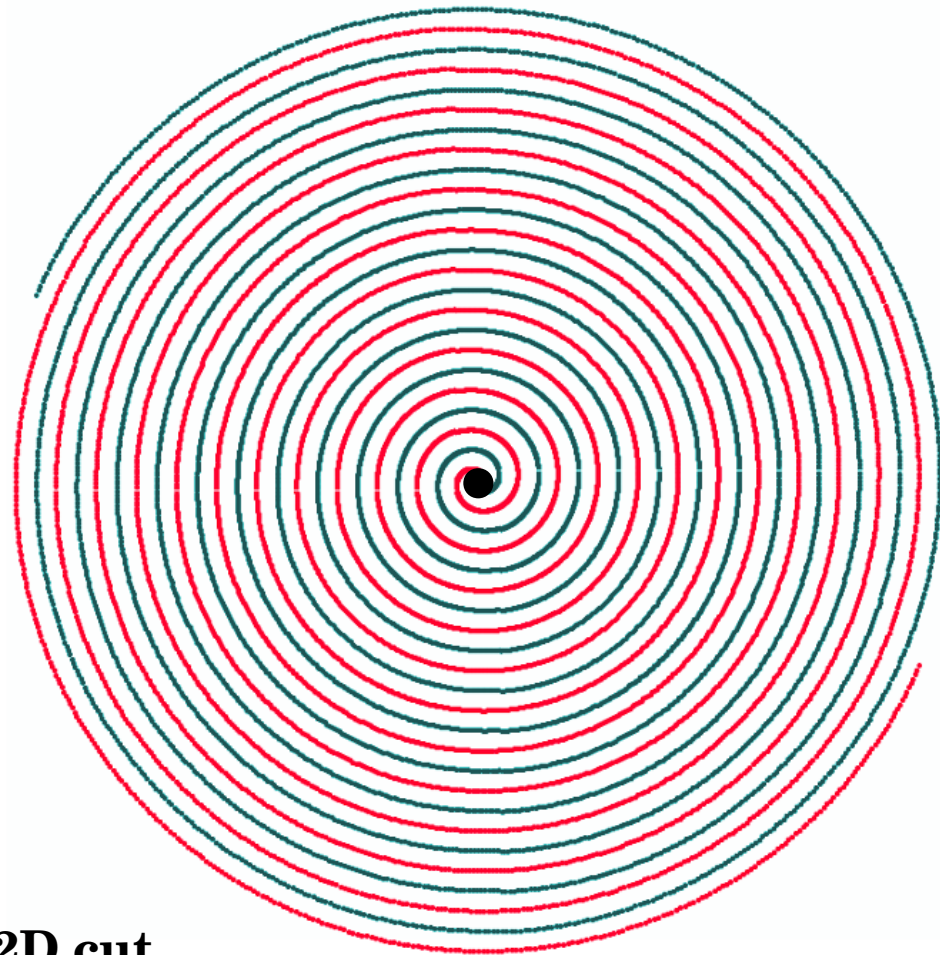
PIC simulations are needed to capture dissipation.
We need to probe large radii, very large box $r \gg R_{LC}$!!

3D simulations too expansive!

=> 2D simulations of the oblique rotator!?



Kirk+2009



Equatorial 2D cut

The numerical setup

Split monopole

[Michel 1973 ; Bogovalov 1999]

4096×4096 cells

$r\phi$ -plane spherical

- logarithmic in r

- uniform in ϕ

$$r_{\max} = 150 R_{\text{LC}}$$

$$R_{\text{LC}}/R_* = 3, 6, 9$$

$$\sigma_* = 250, 1000, 2500$$

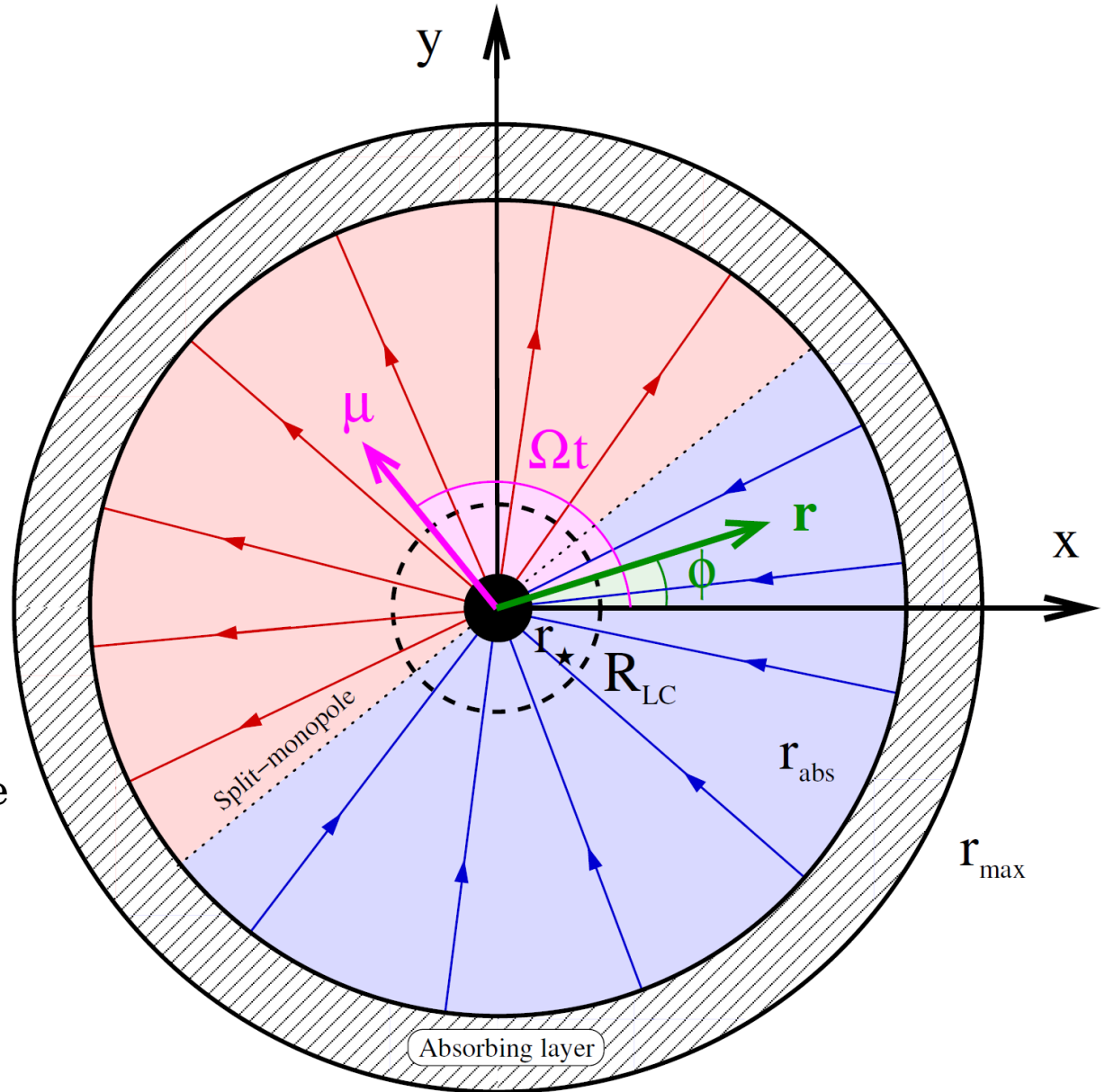
$$\kappa_* = 5, 10, 20$$

Neutral e^{\pm} plasma injected from the surface at the $\mathbf{E} \times \mathbf{B}$ drift velocity :

$$V_r = \frac{c}{1 + R_{\text{LC}}^2/r^2},$$

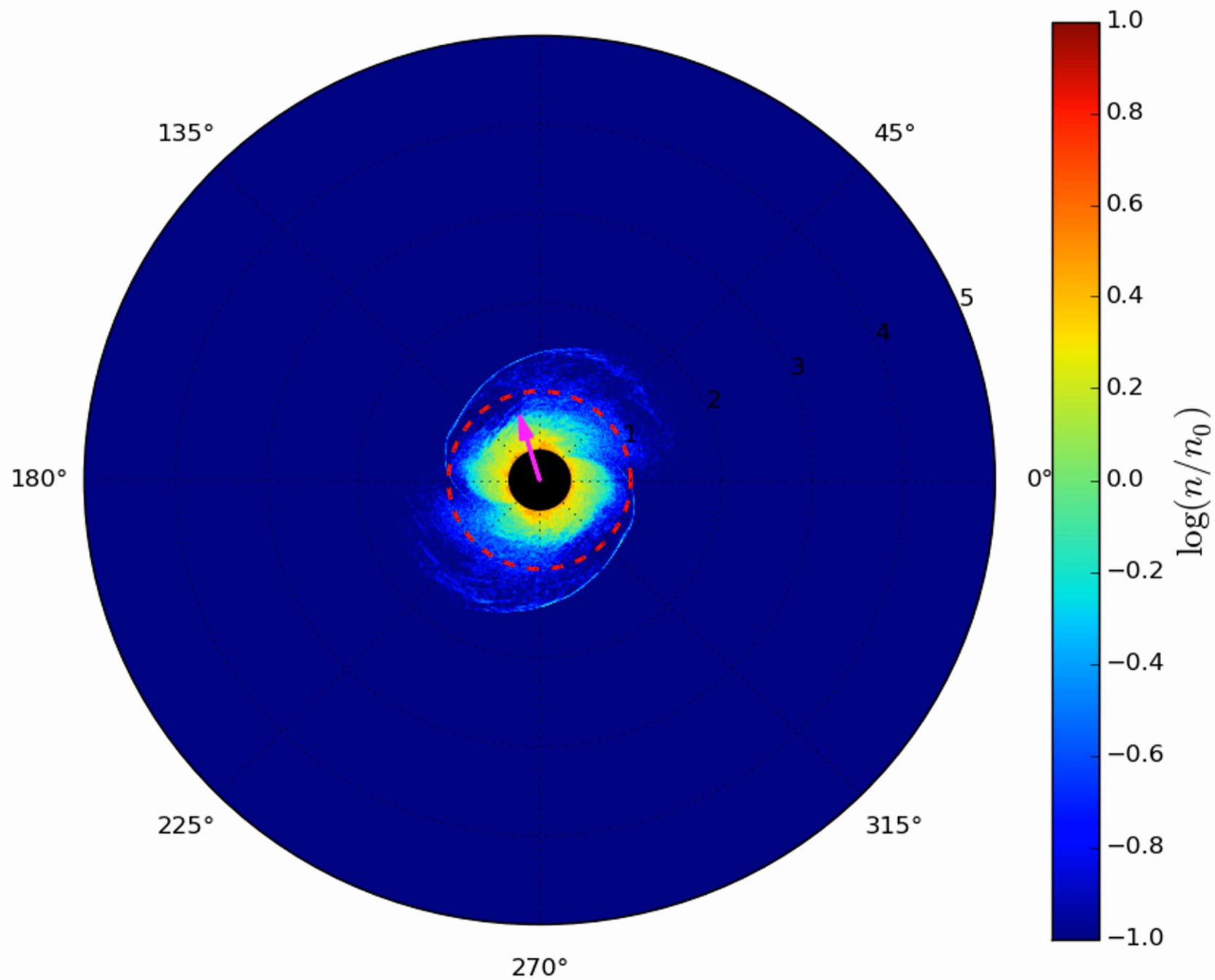
$$V_\theta = 0,$$

$$V_\phi = \frac{r\Omega}{1 + r^2/R_{\text{LC}}^2}.$$

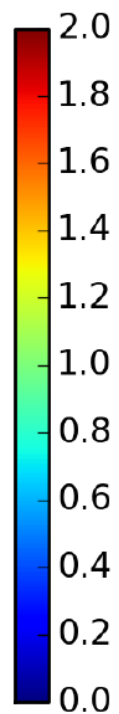
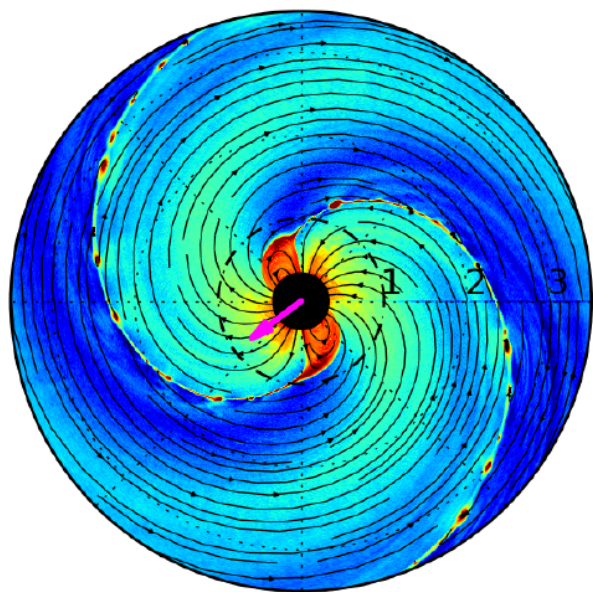


NB: A “cylindrical” ($R\phi$ -plane) pulsar wind does not work ($\mathbf{E}_z \neq \mathbf{B}_\phi$)

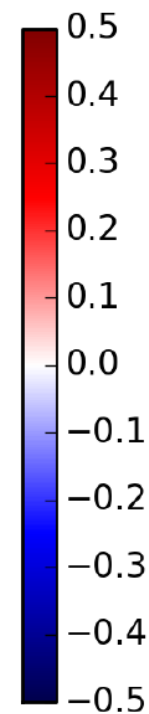
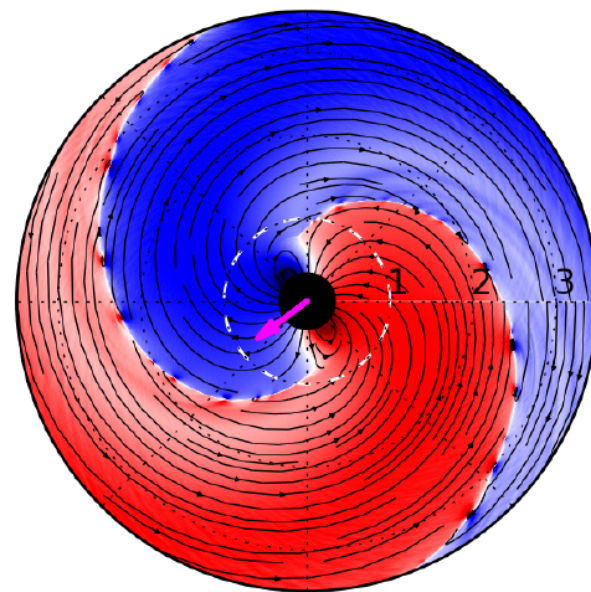
Time=0.2972
90°



Density

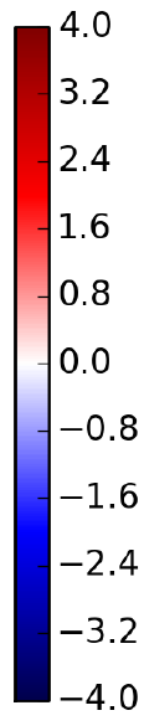
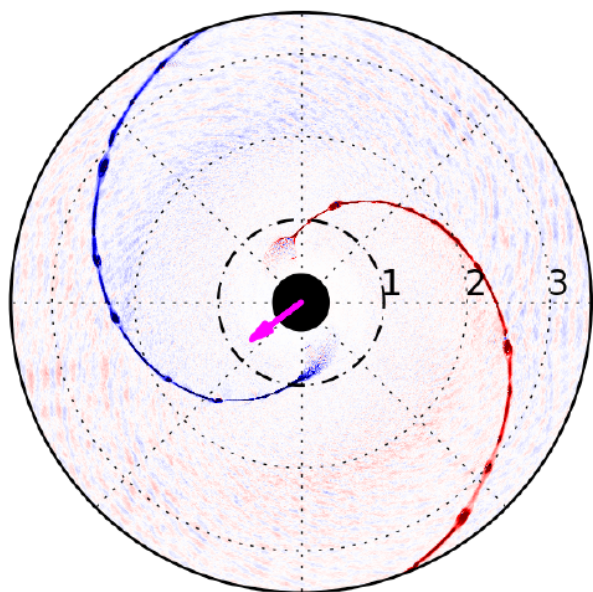


B_ϕ

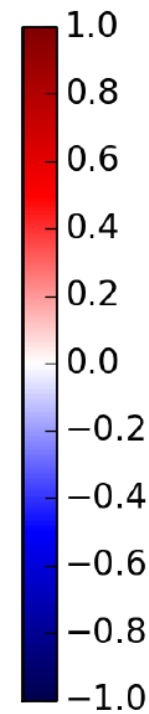
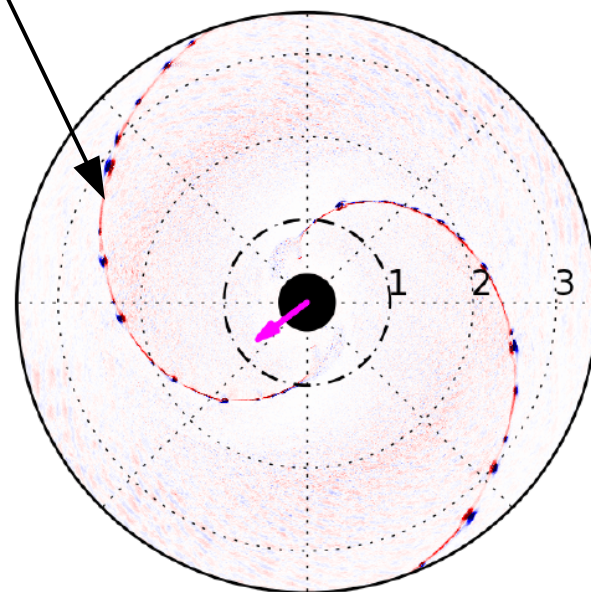


Dissipation within the sheets (between islands)

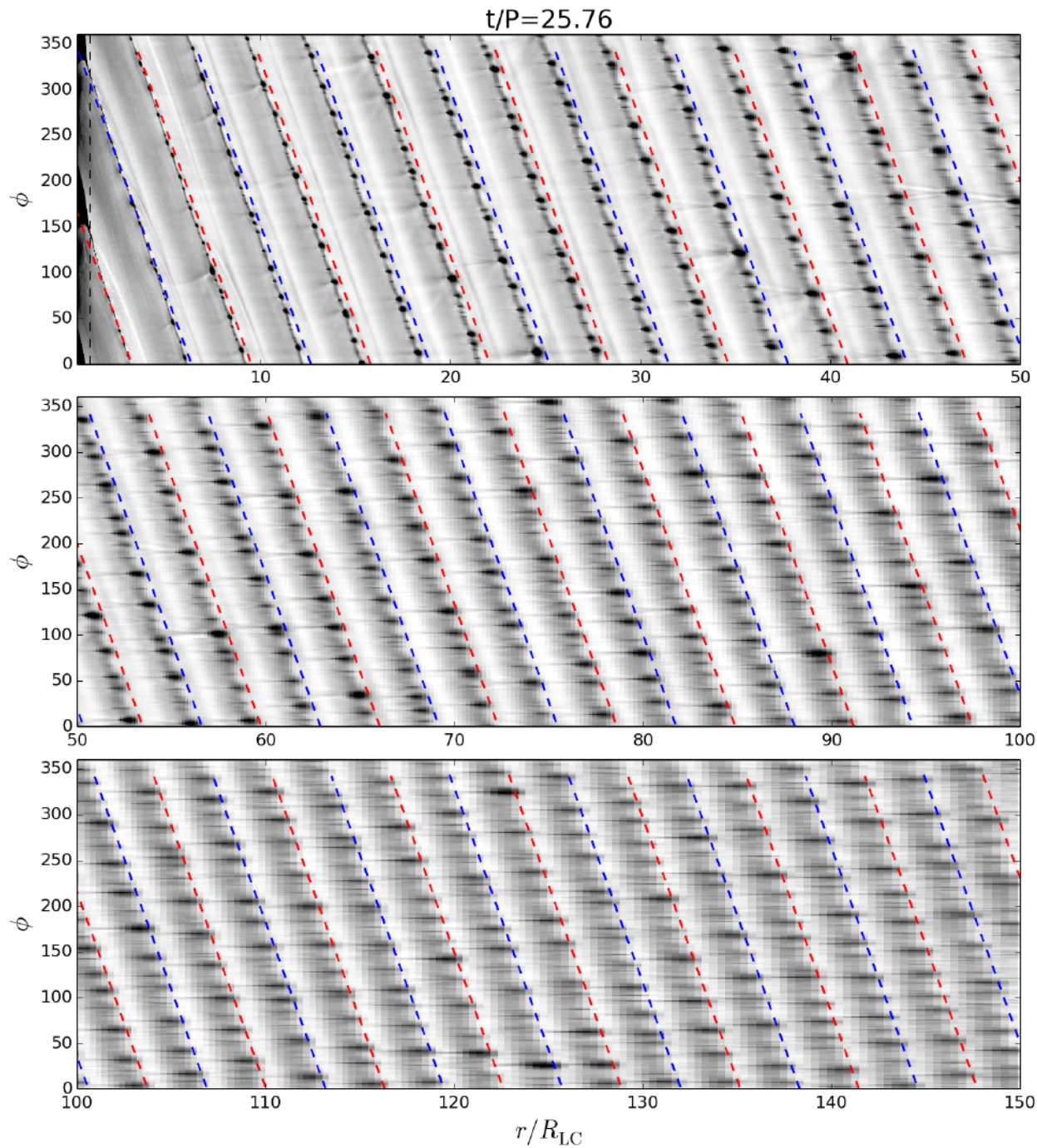
J_θ



$E \cdot J$



The whole box



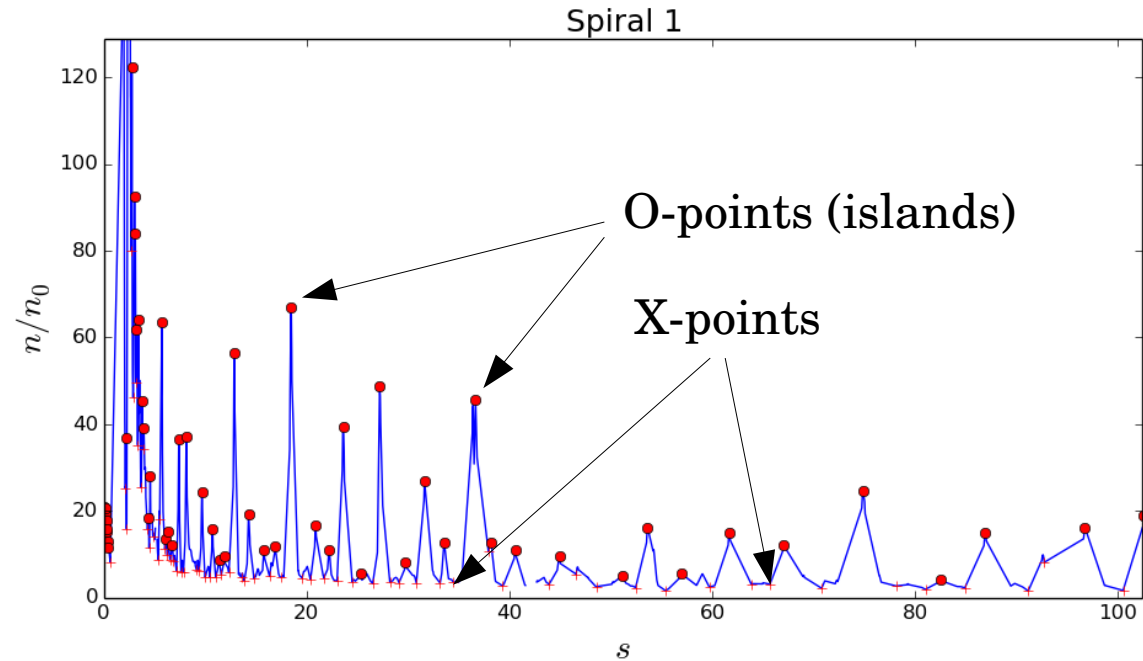
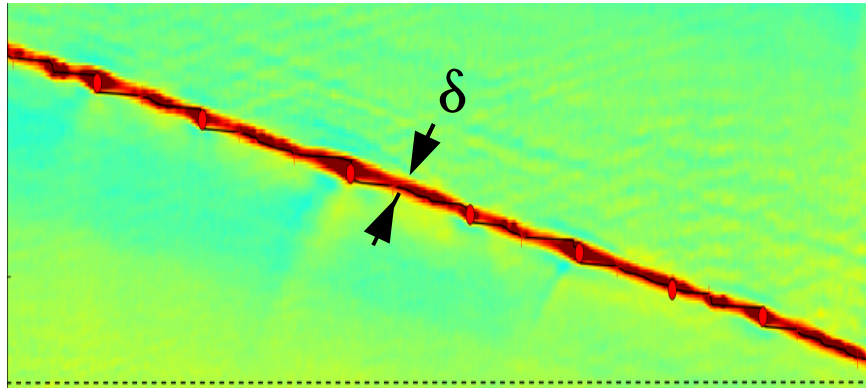
Most dynamical region
Formation of plasmoids
and mergers

Structure frozen by the
expansion of the wind
islands ~ constant

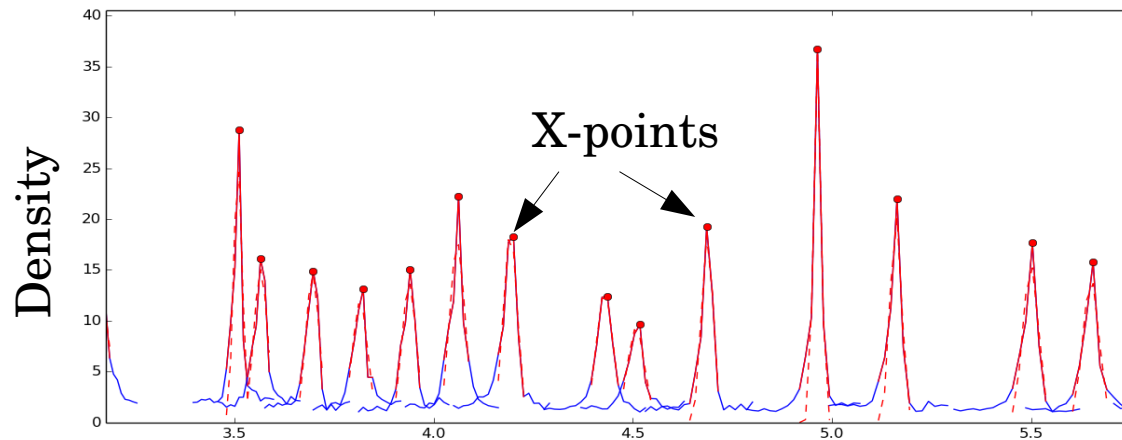


Layer thickness analysis

1. Look for plasma **density minima** along the sheet to **identify X-points**



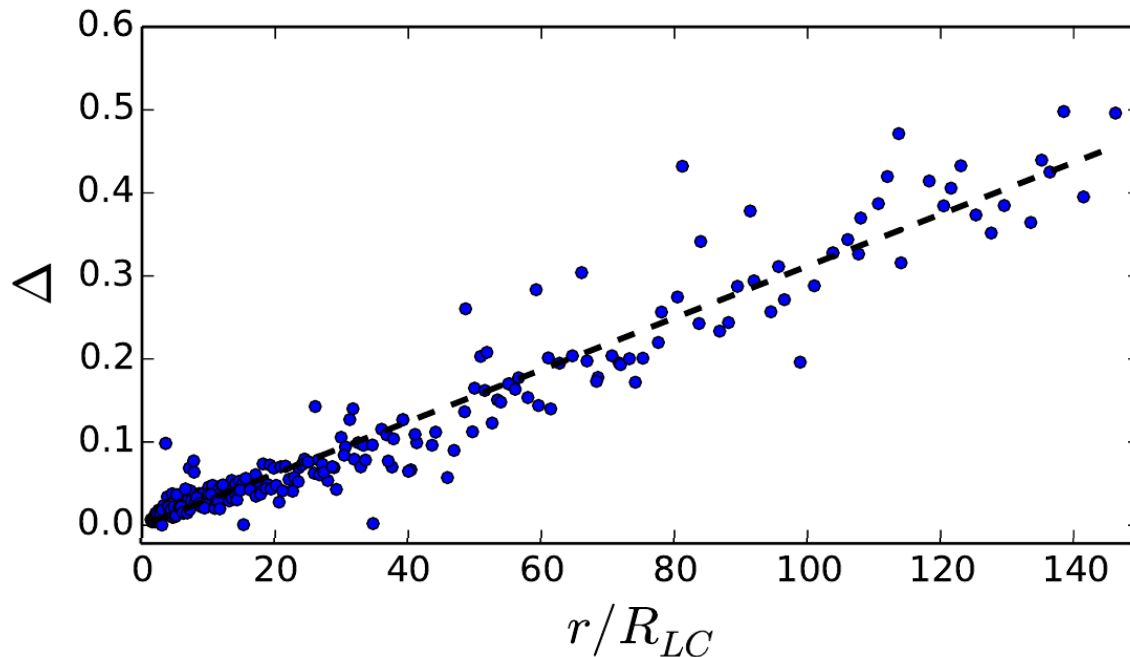
2. Fit the **transverse density profile** with a gaussian



3. Measure the local sheet width with the Full Width at Half Maximum, $\delta(\mathbf{r})=\text{FWHM}$

Expansion of the sheet thickness

Linear expansion of the sheet :



$$\Delta = \delta / \pi R_{LC}$$

$$\Delta(r) \approx \Delta_{LC} \left(\frac{r}{R_{LC}} \right)$$

$$\Delta(r) \approx \frac{1}{\pi \Gamma_{LC} \kappa_{LC}} \left(\frac{r}{R_{LC}} \right) = \frac{1}{\pi \Gamma_{LC} \kappa}$$

Complete dissipation $\Delta=1 \Rightarrow r_{\text{diss}}/R_{LC} = \pi \Gamma_{LC} \kappa_{LC} \sim 10^3 - 10^6 \ll R_{\text{shock}}/R_{LC}$

See also Lyubarsky & Kirk 2001

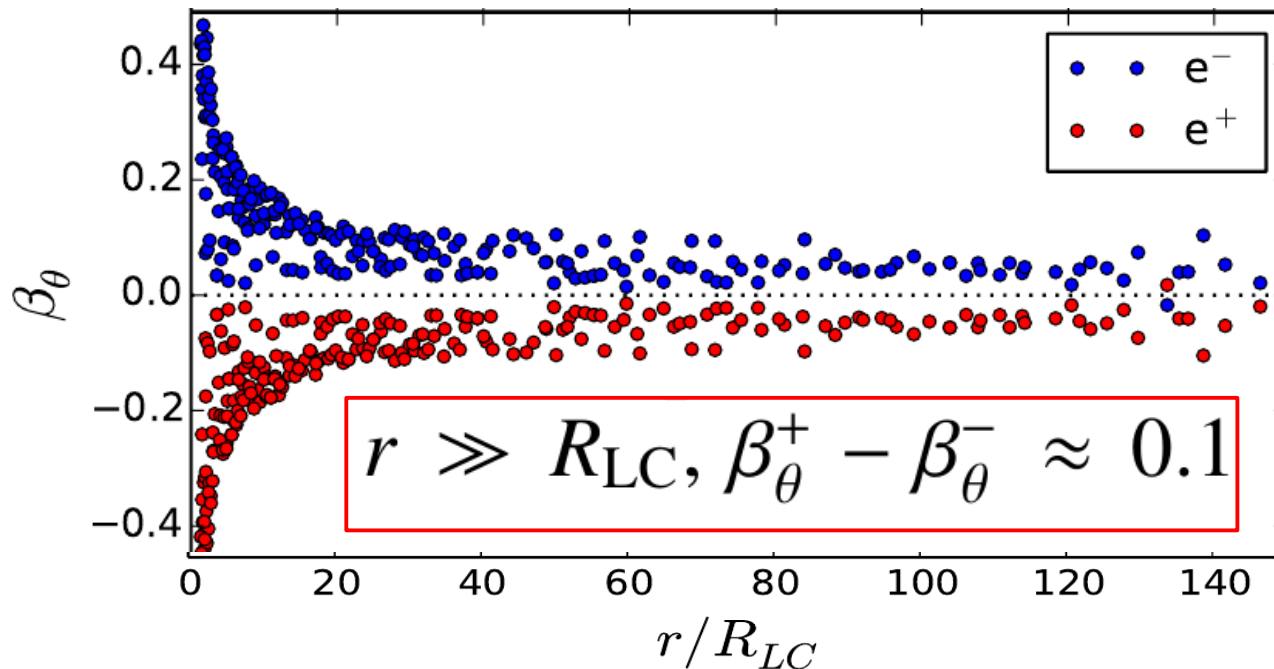
Expansion of the sheet thickness

Results consistent with *Coroniti (1990)* and *Michel (1994)*

Ampère's law accross the sheet yields ($r \gg R_{LC}$) :

$$2B_{\phi} = \frac{4\pi}{c} J_{\theta} \delta = 4\pi n e (\beta_{\theta}^{+} - \beta_{\theta}^{-}) \delta.$$

$\sim 1/r$ $\sim 1/r^2$ $\sim r$

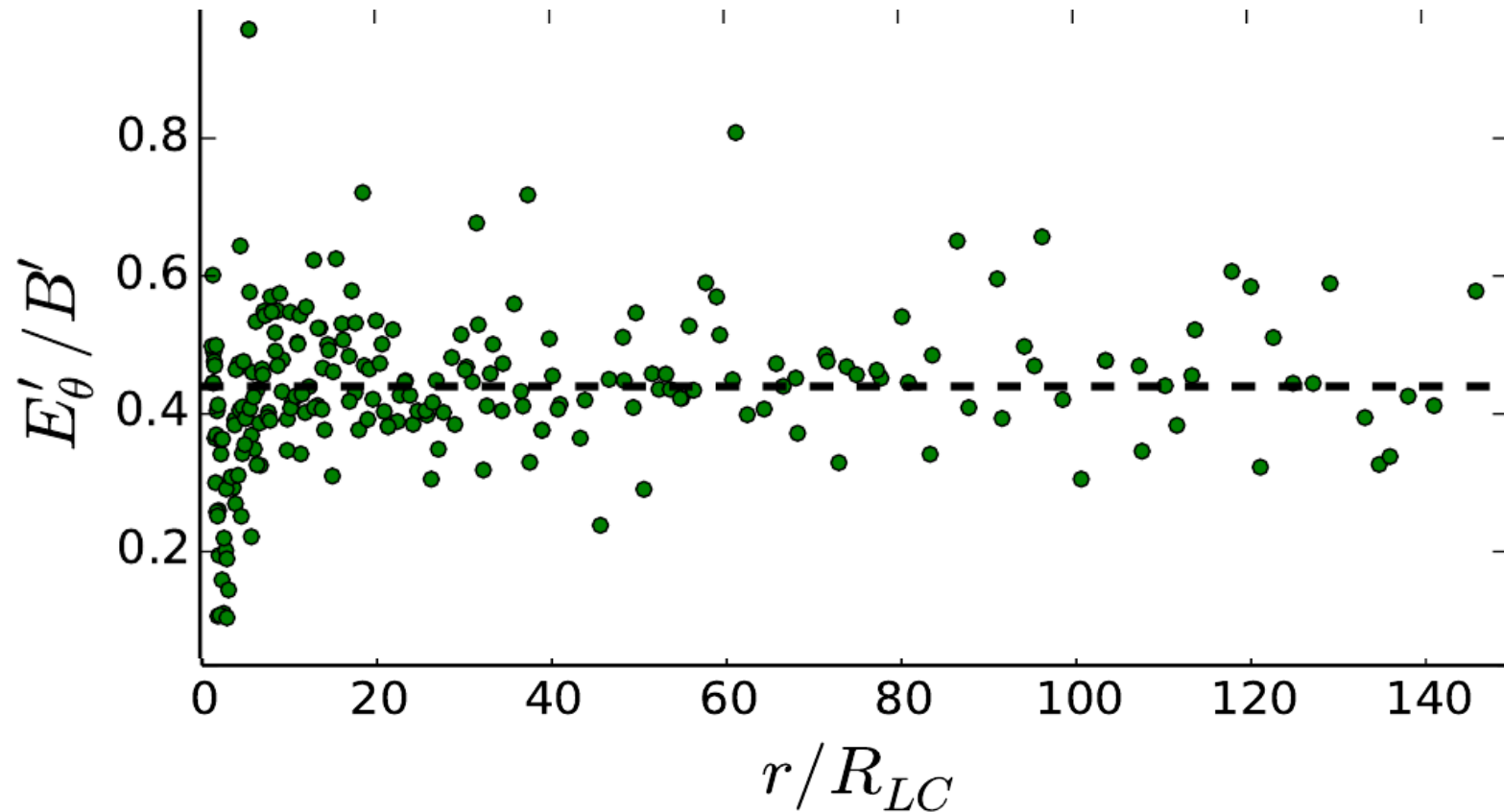


Important consequence : **Current starvation does not happen** if there are enough charges to begin with, in agreement with *Arons 2012*, but in contradiction with *Usov 1975*.

Reconnection rate

Appropriately described into the co-moving frame

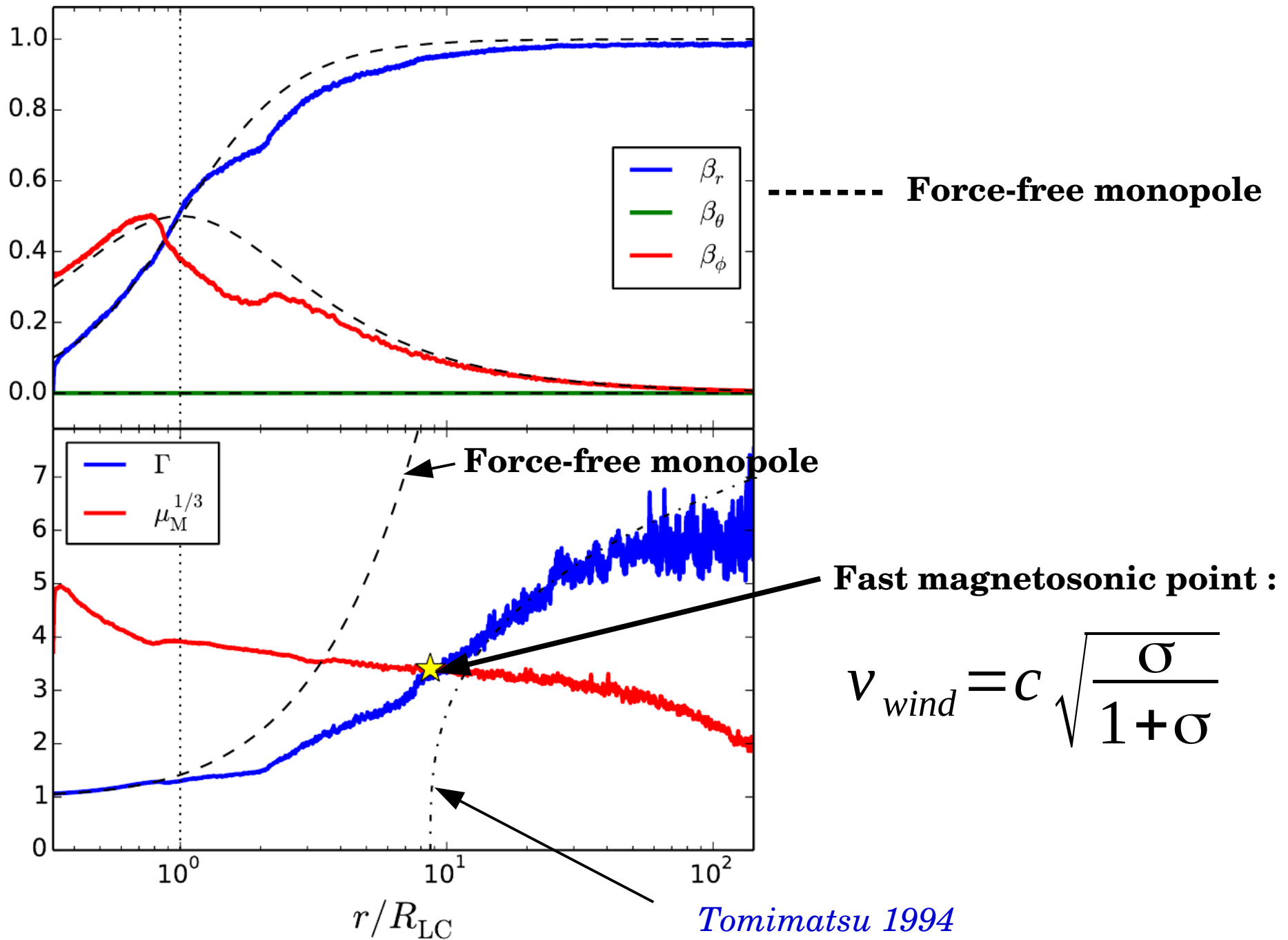
[Lyubarskii 1996 ; Uzdensky & Spitkovsky 2014]



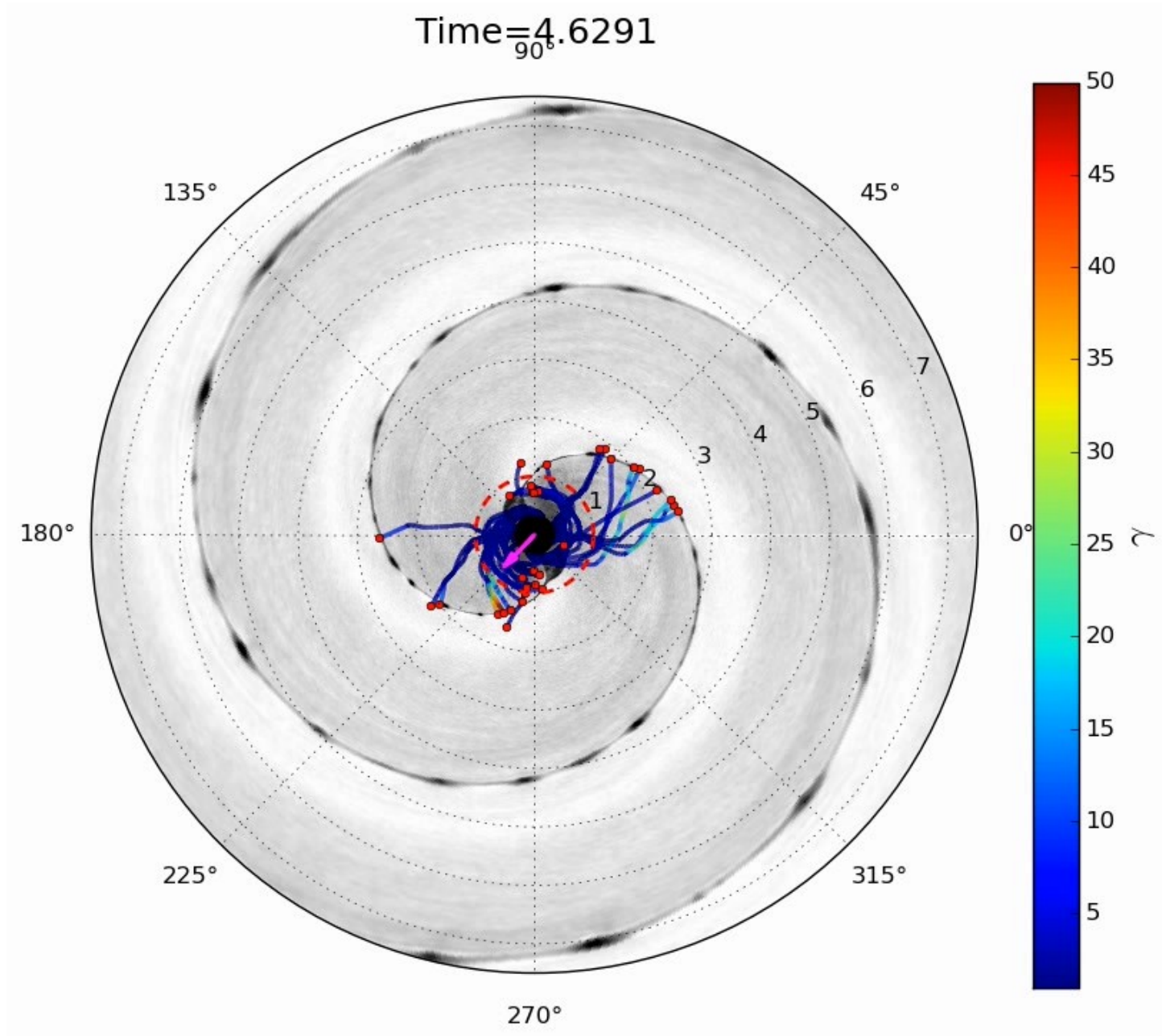
Hard to measure (the wind accelerates) but approximatively constant ~ 0.45
High rate compared with local simulations of reconnection ($\sim 0.1-0.2$).

Reconnection driven by large scale plasma motion ?

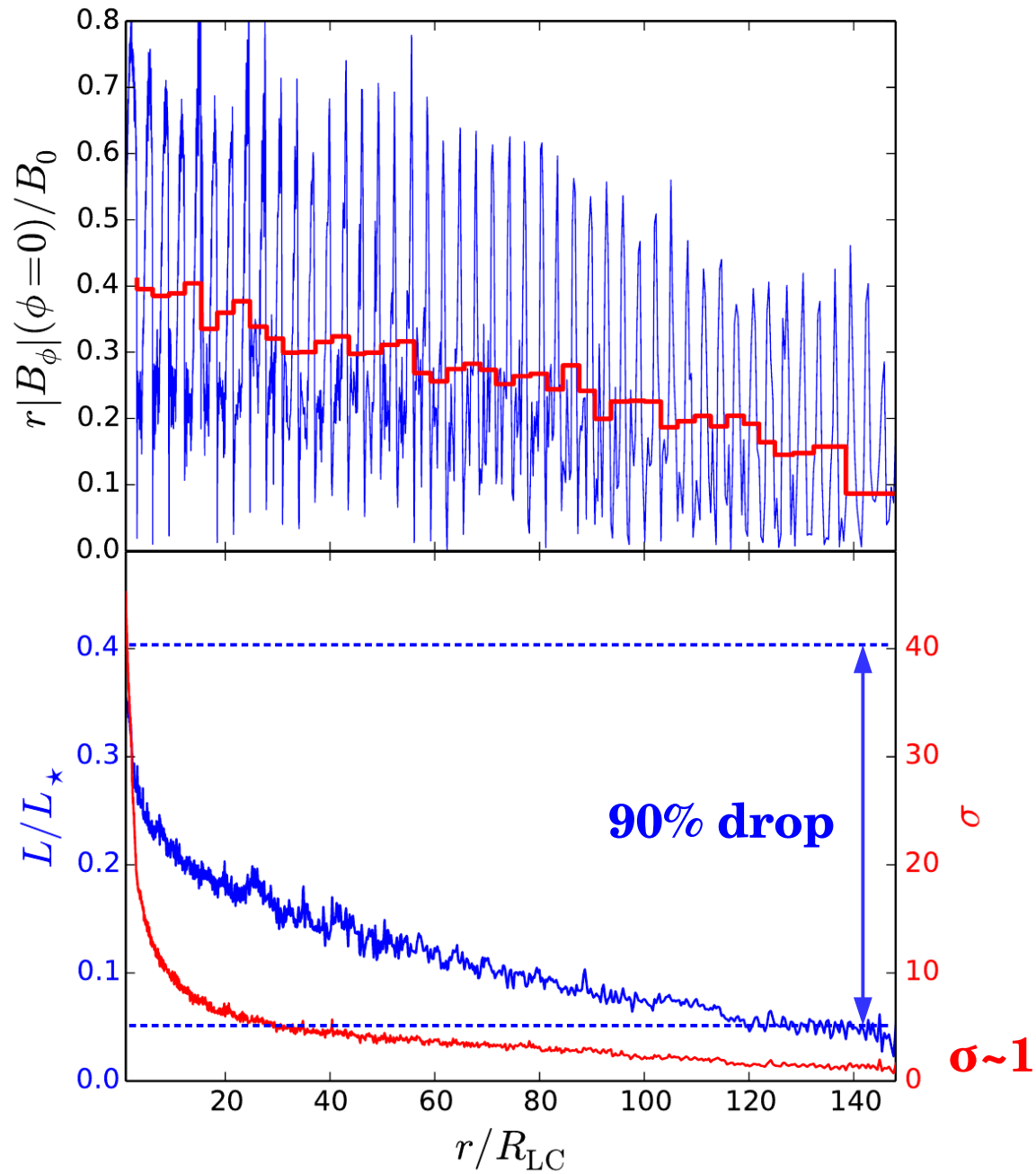
Wind kinematics



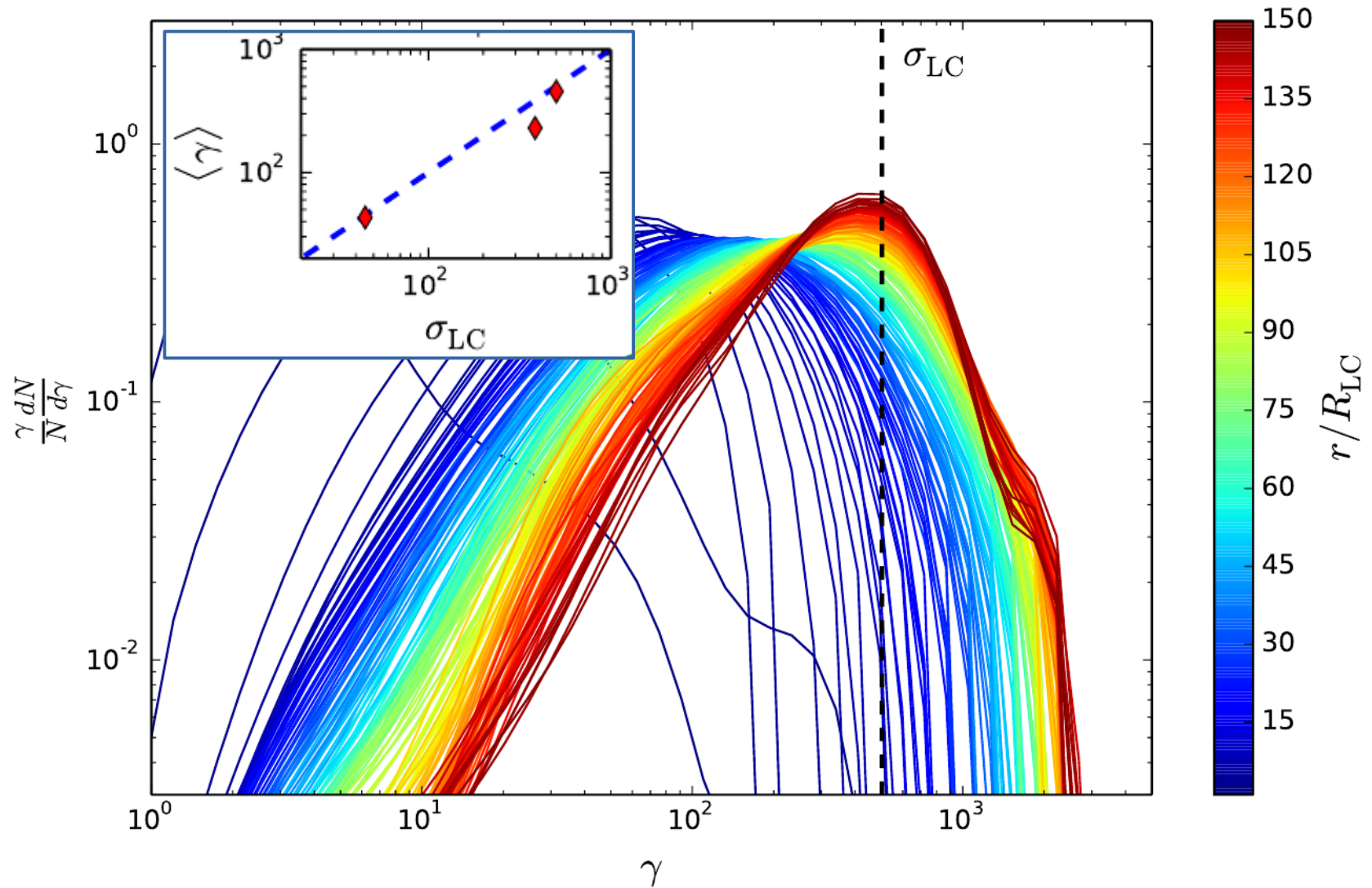
Particle acceleration



Magnetic dissipation in the wind



Particle spectral evolution in the wind



“Narrow” particle energy distribution set by :

$$\sigma_{LC} \sim \Phi_{pc} / \Gamma_{LC} \kappa_{LC}$$

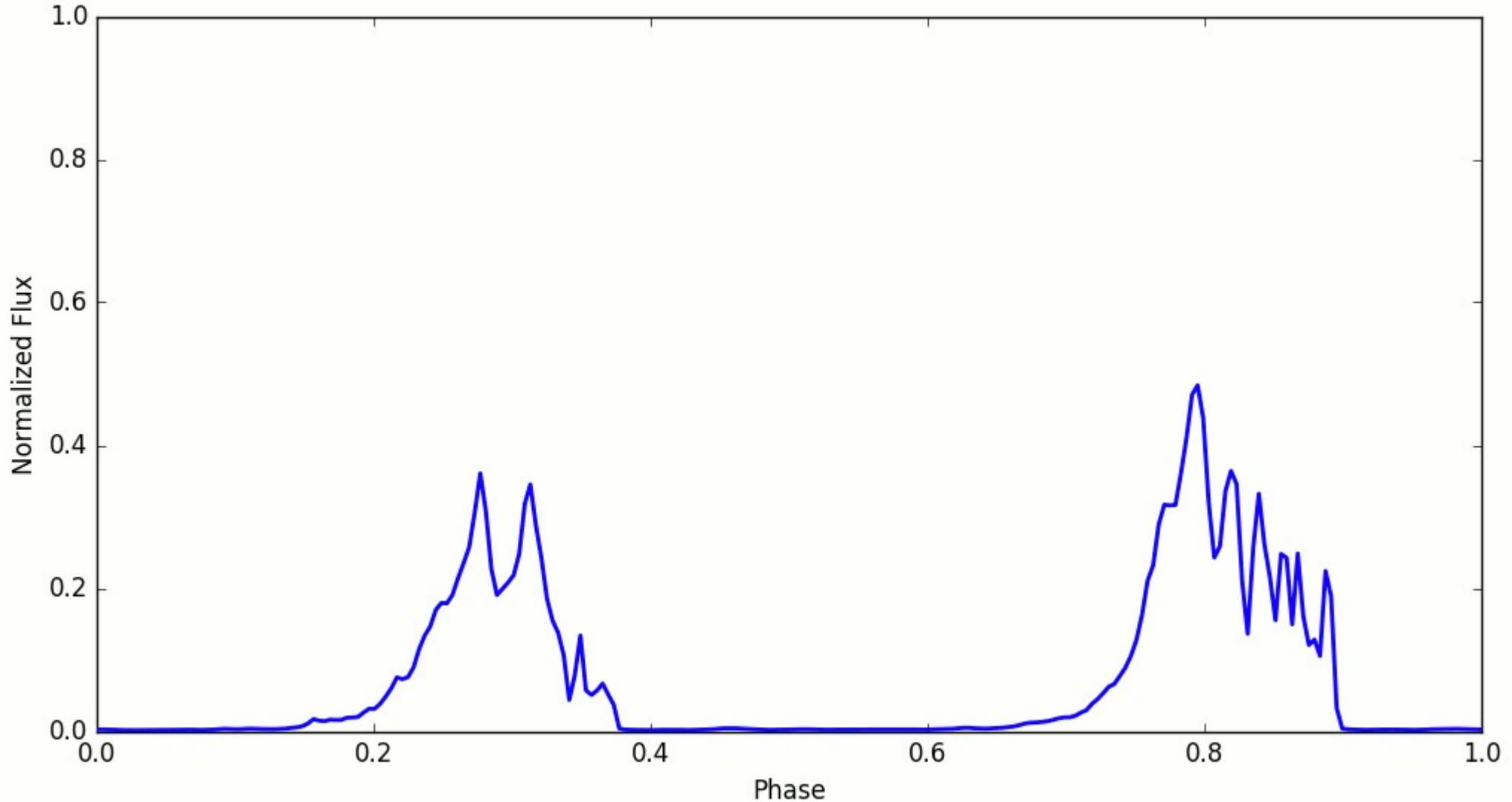
Macroscopic

Microscopic

Radiative signatures

Synchrotron pulse profile : One peak per line-of-sight crossing of the sheet.

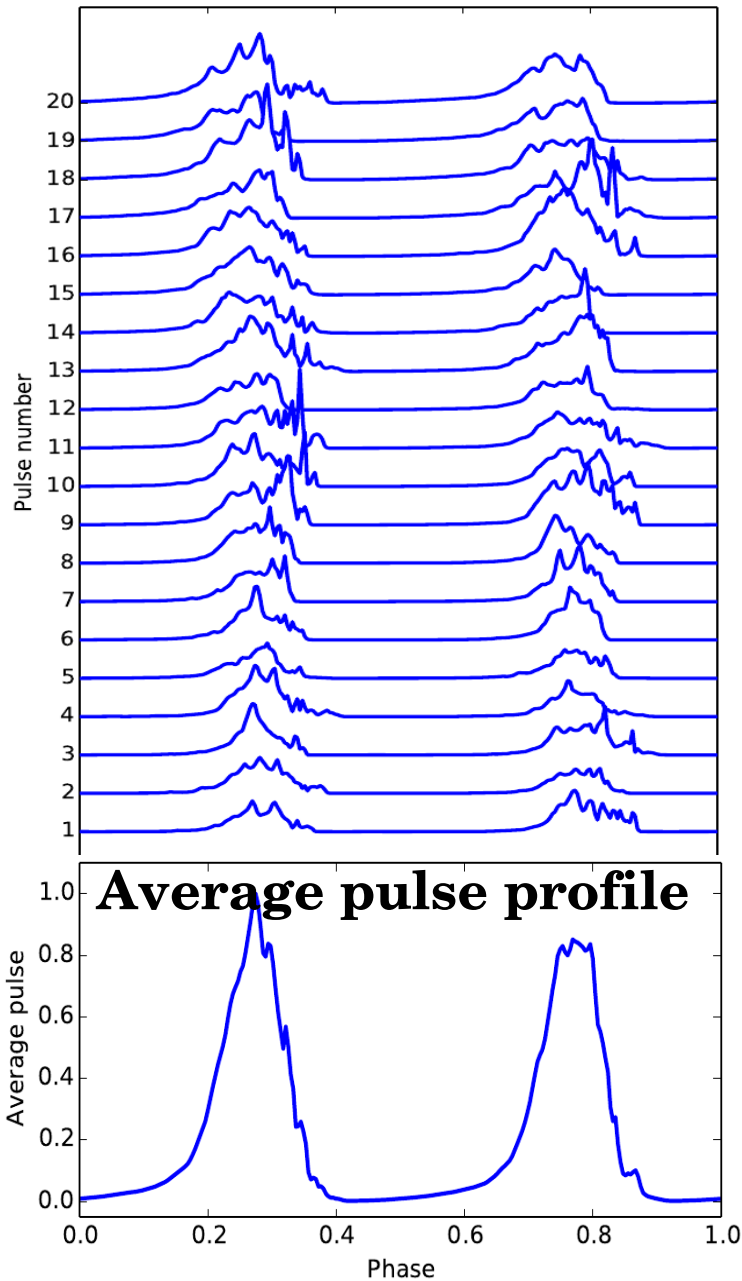
[Cerutti+16, Philippov & Spitkovsky 2017, Kalapotharakos+2017]



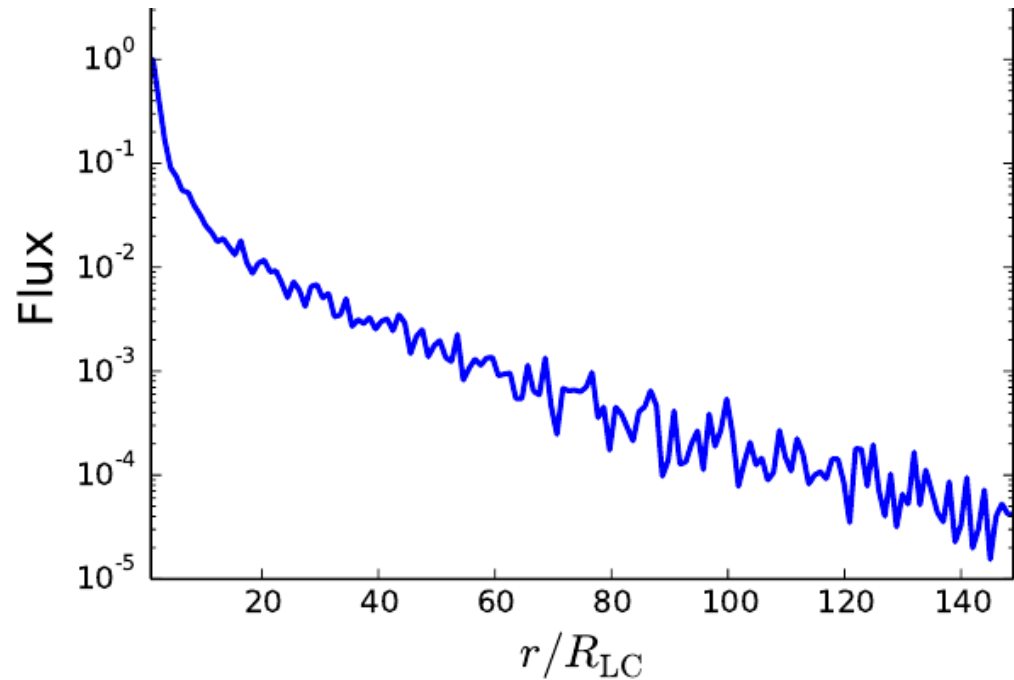
Significant **pulse-to-pulse variability**

Radiative signatures

Synchrotron pulse profiles



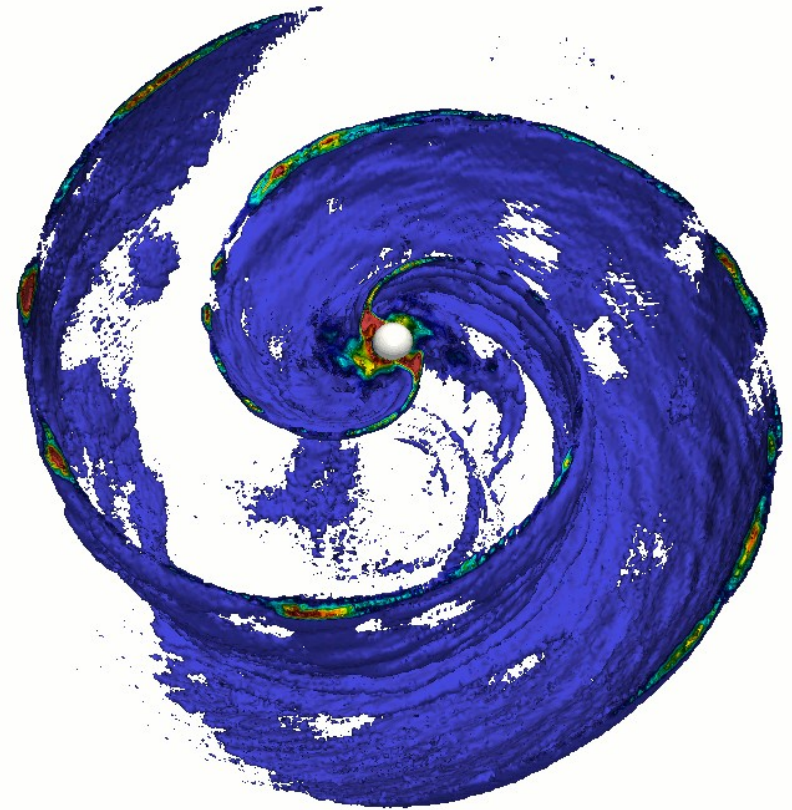
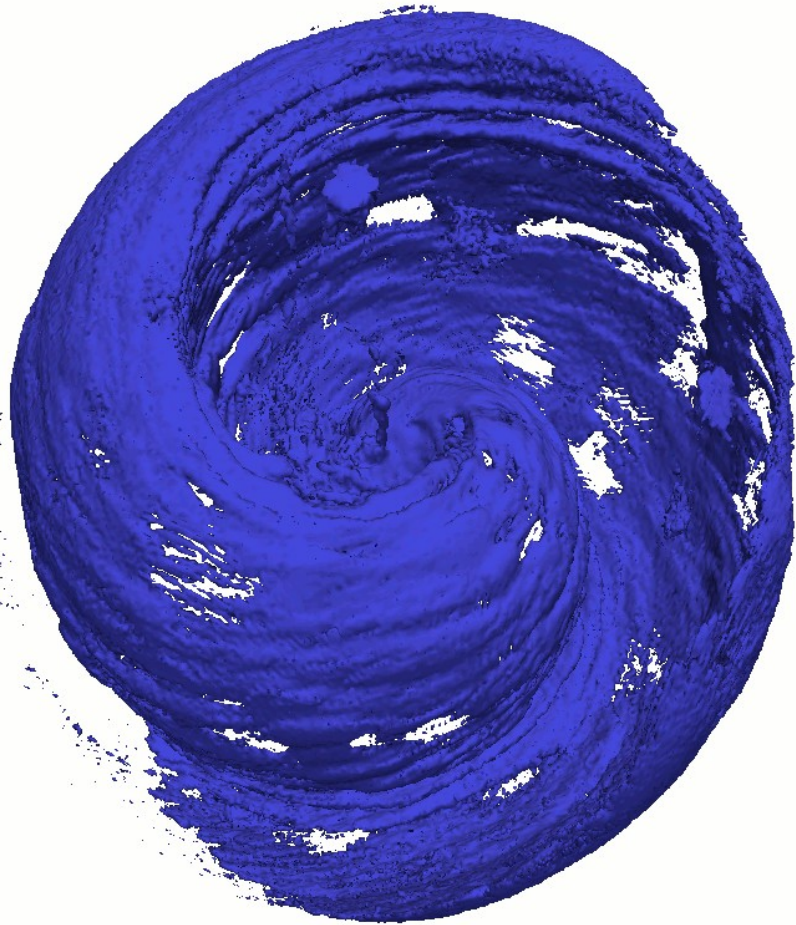
Significant **pulse-to-pulse variability** due to the passage of plasmoids along the line of sight.



Synchrotron emission mostly concentrated **near the light-cylinder**.

=> Gamma-rays probe the **most active regions** of the sheet (formation of islands & merging episodes)

Coming up : Large 3D simulations



Conclusions

- Relativistic reconnection **proceeds in the wind**
- **Complete dissipation** most likely far before the termination shock radius, $R_{\text{diss}}/R_{\text{LC}} \sim \kappa_{\text{LC}} \sim 10^2 - 10^5 \ll R_{\text{shock}}$
- Particle distribution “thermalize” into a **narrow distribution** centered around Lorentz factor given by $\sigma_{\text{LC}} \sim \Phi_{\text{pc}}/\kappa_{\text{LC}}$
- **=> Need better constraints for pair creation @ LC!**
- **Current starvation does not happen** as long as there are enough charges to begin with (at LC)
- **Pulsars in binary** systems (transitional ms & γ -ray binaries) good targets to probe magnetic dissipation within the wind as $R_{\text{shock}} \sim R_{\text{diss}}$