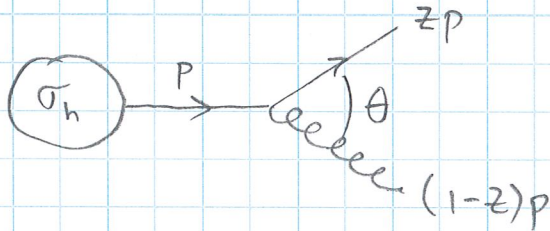


Hard gluon emission



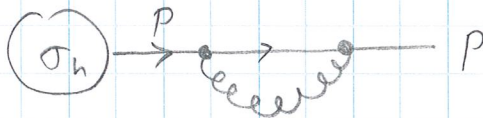
$$C_F = \frac{4}{3} \text{ for 3 colors.}$$

$$\sigma_{h+g} \approx \sigma_h \frac{\alpha_s C_F}{\pi} \frac{dE}{E} \frac{d\theta^2}{\theta^2} = \sigma_h \frac{\alpha_s C_F}{\pi} \frac{dz}{1-z} \frac{dk_t^2}{k_t^2}$$

$$E = (1-z)p$$

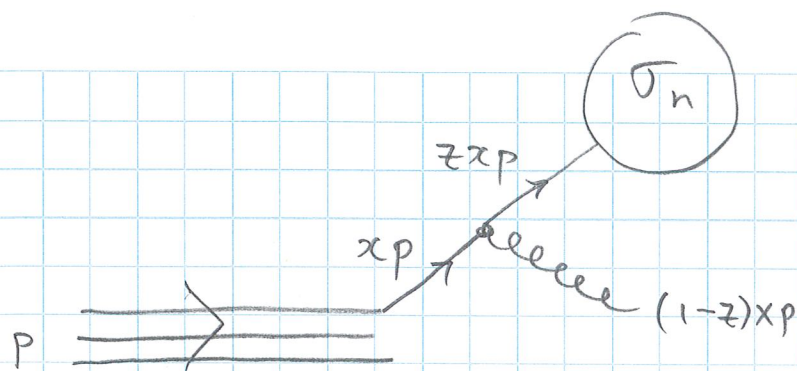
$$k_t = E \sin \theta = E \theta \quad (\theta \ll 1)$$

We must subtract the contribution to the cross section that is indistinguishable from virtual gluon emission



$$\sigma_{g+h} + \sigma_{v+h} = \underbrace{\frac{\alpha_s C_F}{\pi} \int_0^{Q^2} \frac{dk_t^2}{k_t^2}}_{\text{divergent}} \underbrace{\int \frac{dz}{1-z} (\sigma_h(zp) - \sigma_h(p))}_{\text{finite}}$$

We do not expect perturbative QCD to work below some scale, $\mu^2 \approx 1 \text{ GeV}^2$ so we cut off the first integral to make it finite.



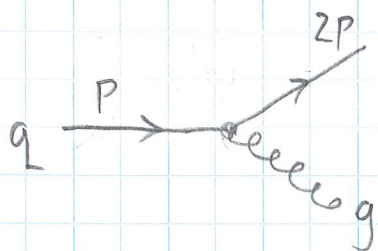
Leading order $\sigma_0 = \int dx \sigma_n(xp) q(x, \mu^2)$

First order

$$\sigma_1 \approx \underbrace{\frac{\alpha_s C_F}{\pi} \int_{\mu^2}^{Q^2} \frac{dk_t^2}{k_t^2}}_{\text{large but finite}} \underbrace{\int \frac{dx dz}{1-z} (\sigma_n(zxp) - \sigma_n(xp))}_{\text{Finite}} q(x, \mu^2)$$

μ^2 is called the factorization scale.

μ^2 is somewhat arbitrary. We need to predict how $q(x, \mu^2)$ will evolve as a function of this factorization scale.



$$q'(x) \sim \int \frac{dz}{1-z} \int_{\mu^2}^{Q^2} \frac{dk_t^2}{k_t^2} q(x/z)$$

$$\sim \log Q^2 \int \frac{dz}{1-z} q(x/z)$$

$$\frac{dq(x, \mu^2)}{d \log Q^2} \sim \int \frac{dz}{1-z} q(x/z, \mu^2)$$

This expresses how q will evolve with $\log Q^2$ due to gluon emission.

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But gluons can also split into quarks

$$\frac{d}{d \log Q^2} \begin{pmatrix} q \\ g \end{pmatrix} = \begin{pmatrix} P_{q \leftarrow q} & P_{q \leftarrow g} \\ P_{g \leftarrow q} & P_{g \leftarrow g} \end{pmatrix} \otimes \begin{pmatrix} q \\ g \end{pmatrix}$$

Splitting functions

$$P_{q \leftarrow q} = C_F \left(\frac{1+z^2}{1-z} \right) \quad C_F = \frac{4}{3}$$

$$P_{g \leftarrow q} = C_F \left(\frac{1+(1-z)^2}{z} \right)$$

$$P_{q \leftarrow g} = T_R (z^2 + (1-z)^2) \quad T_R = \frac{1}{2}$$

$$P_{g \leftarrow g} = 2C_A \left(\frac{z}{1-z} + \frac{1-z}{z} + z(1-z) \right) + \delta(1-z) \dots$$

$C_A = 3$

These allow us to evolve the parton densities to different Q^2 .

Once $f_{q/p}$ or $f_{g/p}$ is known or measured at one Q^2 it can be evolved to different Q^2 .

Evolution in $\log Q^2$: DGLAP

Evolution at small x : BFKL

Example of using pdf sets: LHAPDF

<http://arxiv.org/abs/1412.7420>

This is an interface to many different attempts to analyze DIS data and extrapolate pdf's to different Q^2, x .

Example: CT14 nnlo

Check the sum rules:

$$\int_0^1 \frac{1}{x} (xu(x) - x\bar{u}(x)) dx \sim 1.93$$

$$\int_0^1 \frac{1}{x} (xd(x) - x\bar{d}(x)) dx \sim 0.96$$

$$\sum_i \int_0^1 x q_i(x) dx = 0.446$$

$$\sum_i \int_0^1 x \bar{q}_i(x) dx = 0.100$$

$$\int_0^1 x g(x) dx = 0.440$$

$$\Sigma = 0.986$$