

Assignment # 3

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1. Consider the current operator

$$j^\mu = \bar{\psi} \gamma^\mu \psi$$

(a) Show that $\partial_\mu j^\mu = 0$

$$\partial_\mu j^\mu = (\partial_\mu \bar{\psi}) \gamma^\mu \psi + \bar{\psi} \gamma^\mu \partial_\mu \psi$$

$$\text{But } (i \gamma^\mu \partial_\mu - m) \psi = 0$$

$$\text{and } \bar{\psi} (i \overleftarrow{\partial}_\mu \gamma^\mu + m) = 0$$

$$\text{So, } (\partial_\mu \bar{\psi}) \gamma^\mu = i m \bar{\psi}$$

$$\gamma^\mu (\partial_\mu \psi) = -i m \psi$$

$$\text{So } \partial_\mu j^\mu = i m \bar{\psi} \psi - i m \bar{\psi} \psi = 0$$

(b) Express $Q = \int d^3x j^0(x)$ in terms of creation and annihilation operators in normal order.

$$j^0 = \bar{\psi} \gamma^0 \psi = \psi^\dagger \psi$$

$$\text{Using } \psi(x) = \int \frac{d^3k}{(2\pi)^3} \frac{m}{\omega_k} \sum_{\lambda=1,2} \left[u^{(\lambda)}(k) b^{(\lambda)}(k) e^{-ik \cdot x} + v^{(\lambda)}(k) d^{(\lambda)\dagger}(k) e^{ik \cdot x} \right]$$

$$\psi^\dagger(x) = \int \frac{d^3k}{(2\pi)^3} \frac{m}{\omega_k} \sum_{\lambda=1,2} \left[u^{(\lambda)\dagger}(k) b^{(\lambda)\dagger}(k) e^{ik \cdot x} + v^{(\lambda)\dagger}(k) d^{(\lambda)}(k) e^{-ik \cdot x} \right]$$

$$j^0(x) = \int \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} \frac{m}{\omega_k} \frac{m}{\omega_{k'}} \sum_{\lambda, \lambda'} \left[u^{(\lambda)\dagger}(k) u^{(\lambda')}(k') b^{(\lambda)\dagger}(k) b^{(\lambda')}(k') e^{i(x \cdot (k-k'))} \right. \\ \left. + v^{(\lambda)\dagger}(k) v^{(\lambda')}(k') d^{(\lambda)}(k) d^{(\lambda')\dagger}(k') e^{i(x \cdot (k'-k))} \right] \\ + (\text{stuff}) [b^{(\lambda)}(k) d^{(\lambda')}(k) + \text{etc}]$$

Hence, $Q = \int j^0(x) d^3x$
 $= \int d^3x \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} \frac{m}{\omega_k} \frac{m}{\omega_{k'}}$
 $\times \sum_{\lambda, \lambda'} \left[u^{(\lambda)\dagger}(k) u^{(\lambda')}(k') b^{(\lambda)\dagger}(k) b^{(\lambda')}(k') e^{i\vec{x} \cdot (\vec{k} - \vec{k}')} \right.$
 $\left. + v^{(\lambda)\dagger}(k) v^{(\lambda')}(k') d^{(\lambda)}(k) d^{(\lambda')\dagger}(k') e^{i\vec{x} \cdot (\vec{k}' - \vec{k})} \right]$

But $\int \frac{d^3x}{(2\pi)^3} e^{i\vec{x} \cdot (\vec{k} - \vec{k}')} = \delta^3(\vec{k} - \vec{k}')$ so

$Q = \int \frac{d^3k}{(2\pi)^3} \frac{d^3k'}{(2\pi)^3} \frac{m}{\omega_k} \frac{m}{\omega_{k'}} \sum_{\lambda, \lambda'} \left[u^{(\lambda)\dagger}(k) u^{(\lambda')}(k') b^{(\lambda)\dagger}(k) b^{(\lambda')}(k') \right.$
 $\left. + v^{(\lambda)\dagger}(k) v^{(\lambda')}(k') d^{(\lambda)}(k) d^{(\lambda')\dagger}(k') \right] \delta^3(\vec{k} - \vec{k}')$

Integrating over d^3k' gives:

$Q = \int \frac{d^3k}{(2\pi)^3} \left(\frac{m}{\omega_k}\right)^2 \sum_{\lambda, \lambda'} \left[u^{(\lambda)\dagger}(k) u^{(\lambda')}(k) b^{(\lambda)\dagger}(k) b^{(\lambda')}(k) \right.$
 $\left. + v^{(\lambda)\dagger}(k) v^{(\lambda')}(k) d^{(\lambda)}(k) d^{(\lambda')\dagger}(k) \right]$

But $u^{(\lambda)\dagger}(k) u^{(\lambda')}(k) = \frac{\omega_k}{m} \delta^{\lambda\lambda'}$

and $v^{(\lambda)\dagger}(k) v^{(\lambda')}(k) = \frac{\omega_k}{m} \delta^{\lambda\lambda'}$

So $Q = \int \frac{d^3k}{(2\pi)^3} \frac{m}{\omega_k} \sum_{\lambda} \left[b^{(\lambda)\dagger}(k) b^{(\lambda)}(k) + d^{(\lambda)}(k) d^{(\lambda)\dagger}(k) \right]$

Next, since $\{d^{(\lambda)}(k), d^{(\lambda)\dagger}(k)\} = 1$ this

is $Q = \int \frac{d^3k}{(2\pi)^3} \frac{m}{\omega_k} \sum_{\lambda} \left[b^{(\lambda)\dagger}(k) b^{(\lambda)}(k) - d^{(\lambda)\dagger}(k) d^{(\lambda)}(k) \right]$
 $+ Q_0$

(c)

Consider the states $b^{(\lambda)\dagger}(\mathbf{k})|0\rangle$ and $d^{(\lambda)\dagger}(\mathbf{k})|0\rangle$.

These are eigenstates of Q :

$$Q b^{(\lambda)\dagger}(\mathbf{k})|0\rangle = \int \frac{d^3k'}{(2\pi)^3} \frac{m}{\omega_{k'}} \sum_{\lambda'} b^{(\lambda')\dagger}(\mathbf{k}') b^{(\lambda')\dagger}(\mathbf{k}') b^{(\lambda)\dagger}(\mathbf{k})|0\rangle$$

$$\text{But } \{ b^{(\lambda')\dagger}(\mathbf{k}'), b^{(\lambda)\dagger}(\mathbf{k}) \} = (2\pi)^3 \frac{\omega_{\mathbf{k}}}{m} \delta^3(\vec{\mathbf{k}} - \vec{\mathbf{k}}') \delta^{\lambda\lambda'}$$

$$\text{So } Q b^{(\lambda)\dagger}(\mathbf{k})|0\rangle =$$

$$= \int \frac{d^3k}{(2\pi)^3} \frac{m}{\omega_{\mathbf{k}}} \sum_{\lambda'} \cdot (2\pi)^3 \frac{\omega_{\mathbf{k}}}{m} \delta^3(\vec{\mathbf{k}} - \vec{\mathbf{k}}') \delta^{\lambda\lambda'} b^{(\lambda')\dagger}(\mathbf{k}')|0\rangle$$

$$= b^{(\lambda)\dagger}(\mathbf{k})|0\rangle \quad (\text{Eigenvalue is } +1)$$

$$Q d^{(\lambda)\dagger}(\mathbf{k})|0\rangle = \int \frac{d^3k'}{(2\pi)^3} \frac{m}{\omega_{k'}} \sum_{\lambda'} (-d^{(\lambda')\dagger}(\mathbf{k}') d^{(\lambda')\dagger}(\mathbf{k}')) d^{(\lambda)\dagger}(\mathbf{k})|0\rangle$$

$$\text{But } \{ d^{(\lambda')\dagger}(\mathbf{k}'), d^{(\lambda)\dagger}(\mathbf{k}) \} = (2\pi)^3 \frac{\omega_{\mathbf{k}}}{m} \delta^3(\vec{\mathbf{k}} - \vec{\mathbf{k}}') \delta^{\lambda\lambda'}$$

$$\text{So } Q d^{(\lambda)\dagger}(\mathbf{k})|0\rangle =$$

$$= - \int \frac{d^3k}{(2\pi)^3} \frac{m}{\omega_{\mathbf{k}}} \sum_{\lambda'} (2\pi)^3 \frac{\omega_{\mathbf{k}}}{m} \delta^3(\vec{\mathbf{k}} - \vec{\mathbf{k}}') \delta^{\lambda\lambda'} d^{(\lambda')\dagger}(\mathbf{k}')|0\rangle$$

$$= - d^{(\lambda)\dagger}(\mathbf{k})|0\rangle \quad (\text{Eigenvalue is } -1)$$

(4)

2. In the chiral representation,

$$\frac{1}{2} \vec{\Sigma} \cdot \hat{k} = \frac{1}{2|\vec{k}|} \begin{pmatrix} \vec{\sigma} \cdot \vec{k} & 0 \\ 0 & \vec{\sigma} \cdot \vec{k} \end{pmatrix}$$

(a) In the limit $E \gg m$,

$$u(k) = \frac{1}{\sqrt{2m(E+m)}} \begin{pmatrix} E+m + \vec{\sigma} \cdot \vec{k} & 0 \\ 0 & E+m - \vec{\sigma} \cdot \vec{k} \end{pmatrix} u(0)$$

$$\approx \frac{\sqrt{E}}{\sqrt{2m}} \begin{pmatrix} 1 + \vec{\sigma} \cdot \hat{k} & 0 \\ 0 & 1 - \vec{\sigma} \cdot \hat{k} \end{pmatrix} u(0)$$

$$\text{Thus, } \frac{1}{2} \vec{\Sigma} \cdot \hat{k} u(k) = \frac{1}{2} \frac{\sqrt{E}}{\sqrt{2m}} \begin{pmatrix} \vec{\sigma} \cdot \hat{k} (1 + \vec{\sigma} \cdot \hat{k}) & 0 \\ 0 & \vec{\sigma} \cdot \hat{k} (1 - \vec{\sigma} \cdot \hat{k}) \end{pmatrix} u(0)$$

$$= \frac{1}{2} \frac{\sqrt{E}}{\sqrt{2m}} \begin{pmatrix} 1 + \vec{\sigma} \cdot \hat{k} & 0 \\ 0 & -1 + \vec{\sigma} \cdot \hat{k} \end{pmatrix} u(0)$$

$$= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} u(k)$$

$$= \frac{1}{2} \gamma^5 u(k)$$

$$(b) P_{\pm} = \frac{1}{2} (1 \pm \gamma^5)$$

$$P_+ = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

$$P_- = \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} - \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix}$$

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$$\text{Thus, } P_+ u(k) = \frac{1}{\sqrt{2m(E+m)}} \begin{pmatrix} E+m+\vec{\sigma}\cdot\vec{k} & 0 \\ 0 & 0 \end{pmatrix} u(0)$$

which transforms like $e^{\vec{\sigma}\cdot\vec{p}/2}$

$$\text{and } P_- u(k) = \frac{1}{\sqrt{2m(E+m)}} \begin{pmatrix} 0 & 0 \\ 0 & E+m-\vec{\sigma}\cdot\vec{k} \end{pmatrix} u(0)$$

which transforms like $e^{-\vec{\sigma}\cdot\vec{p}/2}$.

$$\begin{aligned} \text{(c) } \frac{1}{2} \vec{\Sigma} \cdot \hat{k} P_+ u(k) &= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} u(k) \\ &= \frac{+1}{2} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} u(k) \\ &= \frac{+1}{2} P_+ u(k) . \end{aligned}$$

$$\begin{aligned} \frac{1}{2} \vec{\Sigma} \cdot \hat{k} P_- u(k) &= \frac{1}{2} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} u(k) \\ &= \frac{-1}{2} \begin{pmatrix} 0 & 0 \\ 0 & 1 \end{pmatrix} u(k) \\ &= \frac{-1}{2} P_- u(k) . \end{aligned}$$