Purdue University Department of Physics

Physics 42200

Waves & Oscillations

Lecture 14 – French, Chapter 6

Spring 2013 Semester

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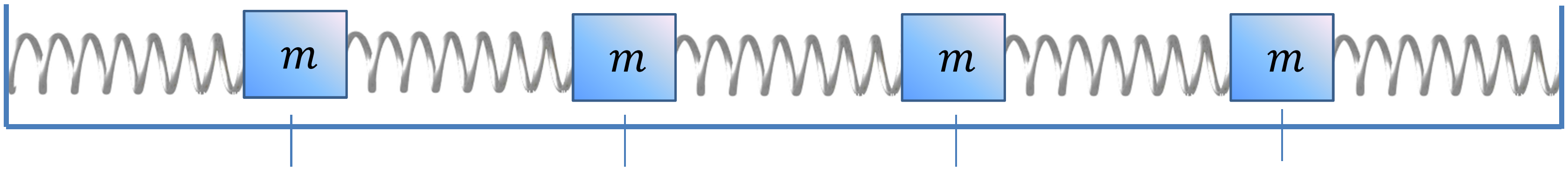
Vibrations of Continuous Systems



- Equations of motion for masses in the middle:

$$m\,\ddot{x}_i + 2kx_i \; - k(x_{i-1} + x_{i+1}) = 0$$

$$\ddot{x}_i + 2(\omega_0)^2 x_i - (\omega_0)^2 (x_{i-1} + x_{i+1}) = 0$$

- Proposed solution:

$$x_i(t) = A_i \cos \omega t$$

$$\frac{A_{i-1} + A_{i+1}}{A_i} = \frac{-\omega^2 + 2(\omega_0)^2}{(\omega_0)^2}$$

- We solved this to determine A_i and ω_j…

Vibrations of Continuous Systems

- Amplitude of mass n for normal mode k:

$$A_{n,k} = C \sin\left(\frac{nk\pi}{N+1}\right)$$

- Frequency of normal mode k:

$$\omega_k = 2\omega_0 \sin\left(\frac{k\pi}{2(N+1)}\right)$$

- Solution for normal modes:

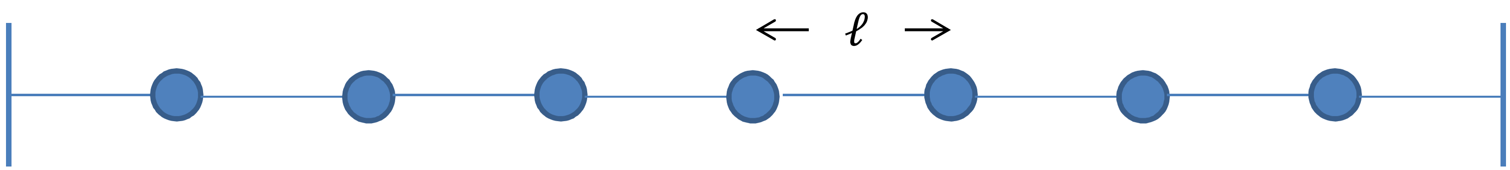
$$x_n(t) = A_{n,k} \cos \omega_k t$$

- General solution:

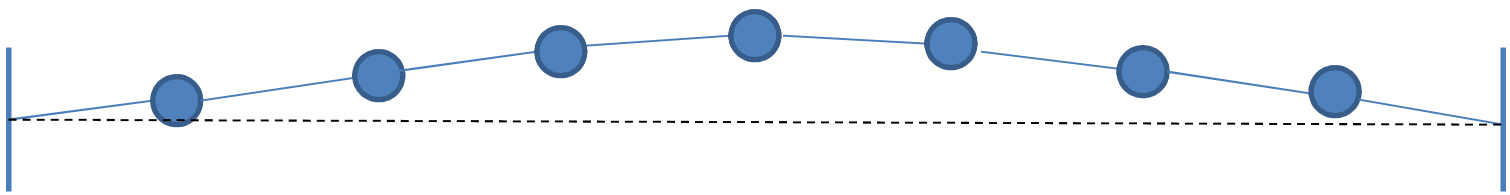
$$x_n(t) = \sum_{k=1}^{N} a_k \sin\left(\frac{nk\pi}{N+1}\right) \cos(\omega_k t - \delta_k)$$

Another Example

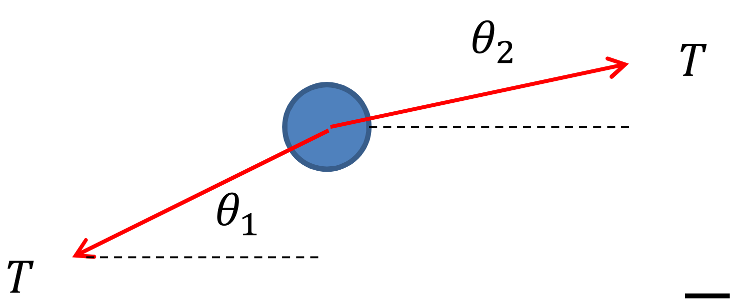
- Discrete masses on an elastic string with tension T:



- Consider transverse displacements:



- Vertical force on one mass:

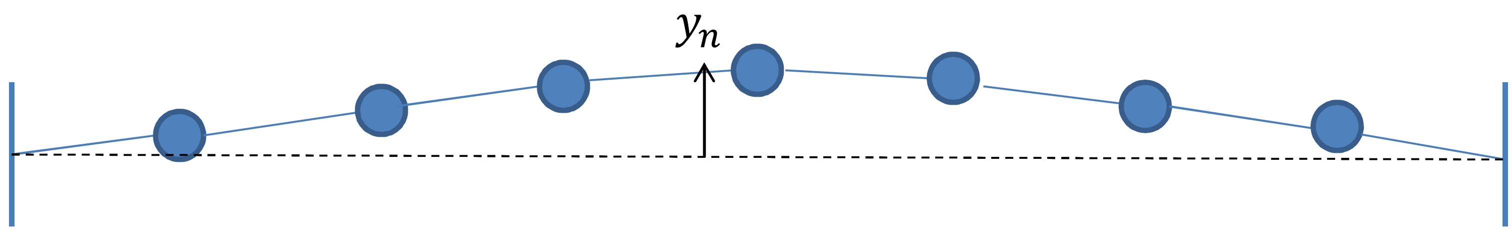


$$F_n = T\sin\theta_2 - T\sin\theta_1$$

$$= T(\theta_2 - \theta_1)$$

$$= \frac{T}{\ell}[(y_{n+1} - y_n) - (y_n - y_{n-1})]$$

Another Example



- Equation of motion for mass n:

$$m\,\ddot{y}_n = F_n = \frac{T}{\ell}\left[(y_{n+1} - y_n) - (y_n - y_{n-1})\right]$$

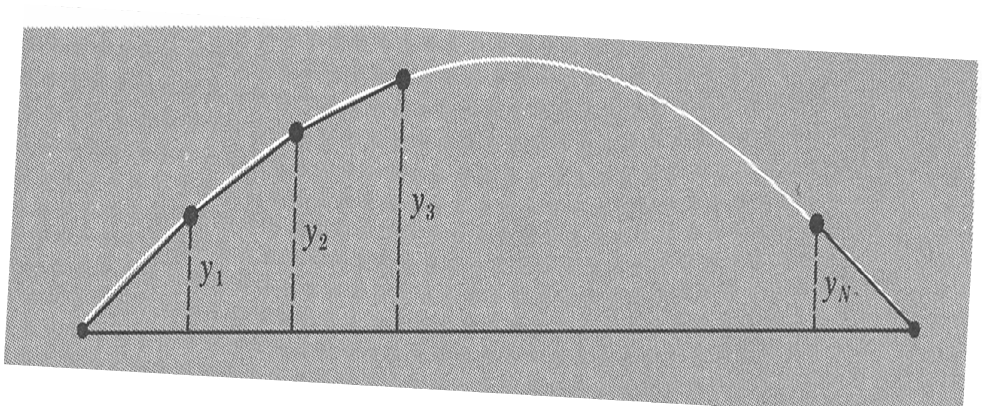
$$\ddot{y}_n + 2(\omega_0)^2 y_n - (\omega_0)^2 (y_{n+1} + y_{n-1}) = 0$$

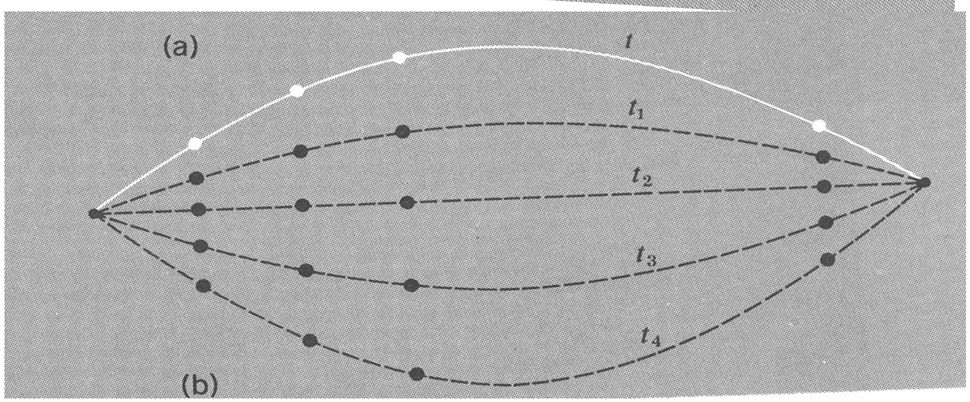
$$(\omega_0)^2 = \frac{T}{m\ell}$$

- Normal modes:

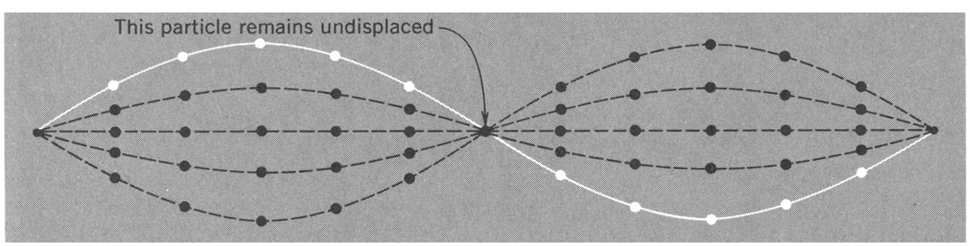
$$y_{n,k}(t) = A_{n,k}\cos(\omega_k t - \delta_k)$$

Masses on a String





First normal mode



Second normal mode

Continuous Systems

- What happens when the number of masses goes to infinity, while the linear mass density remains constant?

$$m\,\ddot{y}_n = \frac{T}{\ell}\left[(y_{n+1} - y_n) - (y_n - y_{n-1})\right]$$

$$\frac{m}{\ell} \rightarrow \mu$$

$$\frac{y_{n+1} - y_n}{\ell} \rightarrow \left(\frac{\partial y}{\partial x}\right)_{x+\Delta x} \qquad \frac{(y_n - y_{n-1})}{\ell} \rightarrow \left(\frac{\partial y}{\partial x}\right)_{x}$$

$$\mu\ell\,\frac{\partial^2 y}{\partial t^2} = T\left[\left(\frac{\partial y}{\partial x}\right)_{x+\Delta x} - \left(\frac{\partial y}{\partial x}\right)_{x}\right]$$

Continuous Systems

$$\mu \frac{\partial^2 y}{\partial t^2} = T \frac{\left(\frac{\partial y}{\partial x}\right)_{x+\Delta x} - \left(\frac{\partial y}{\partial x}\right)_{x}}{\ell}$$

$$\mu \frac{\partial^2 y}{\partial t^2} = T \frac{\partial^2 y}{\partial x^2}$$

$$\frac{\partial^2 y}{\partial x^2} = \frac{\mu}{T} \frac{\partial^2 y}{\partial t^2}$$

The Wave Equation: $$\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2}$$ $$v = \sqrt{T/\mu}$$

Solutions

- When we had N masses, the solutions were
$$y_{n,k}(t) = A_{n,k}\cos(\omega_k t - \delta_k)$$
  - n labels the mass along the string
  - With a continuous system, n is replaced by x.
- Proposed solution to the wave equation for the continuous string:
$$y(x,t) = f(x)\cos\omega t$$
- Derivatives:
$$\frac{\partial^2 y}{\partial t^2} = -\omega^2 f(x)\cos\omega t$$
$$\frac{\partial^2 y}{\partial x^2} = \frac{\partial^2 f}{\partial x^2}\cos\omega t$$

Solutions

- Substitute into the wave equation:

$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2}\frac{\partial^2 y}{\partial t^2}$$

$$\frac{\partial^2 f}{\partial x^2} = -\frac{\omega^2}{v^2} f(x)$$

$$\frac{\partial^2 f}{\partial x^2} + \frac{\omega^2}{v^2} f(x) = 0$$

- This is the same differential equation as for the harmonic oscillator.
- Solutions are $f(x) = A \sin(\omega x/v) + B \cos(\omega x/v)$

Solutions

$$f(x) = A\sin(\omega x/v) + B\cos(\omega x/v)$$

- Boundary conditions at the ends of the string:

$$f(0) = f(L) = 0$$

$$f(x) = A\sin(\omega x/v) \text{ where } \omega L/v = n\pi$$

- Solutions can be written:

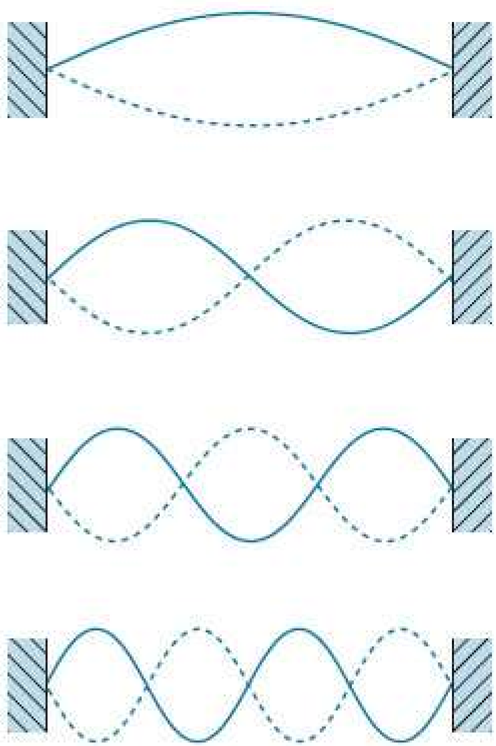
$$f_n(x) = A_n \sin\left(\frac{n\pi x}{L}\right)$$

- Complete solution describing the motion of the whole string:

$$y_n(x,t) = A_n \sin\left(\frac{n\pi x}{L}\right)\cos\omega_n t$$

Properties of the Solutions

$$y_n(x,t) = A_n \sin\left(\frac{n\pi x}{L}\right) \cos \omega_n t$$



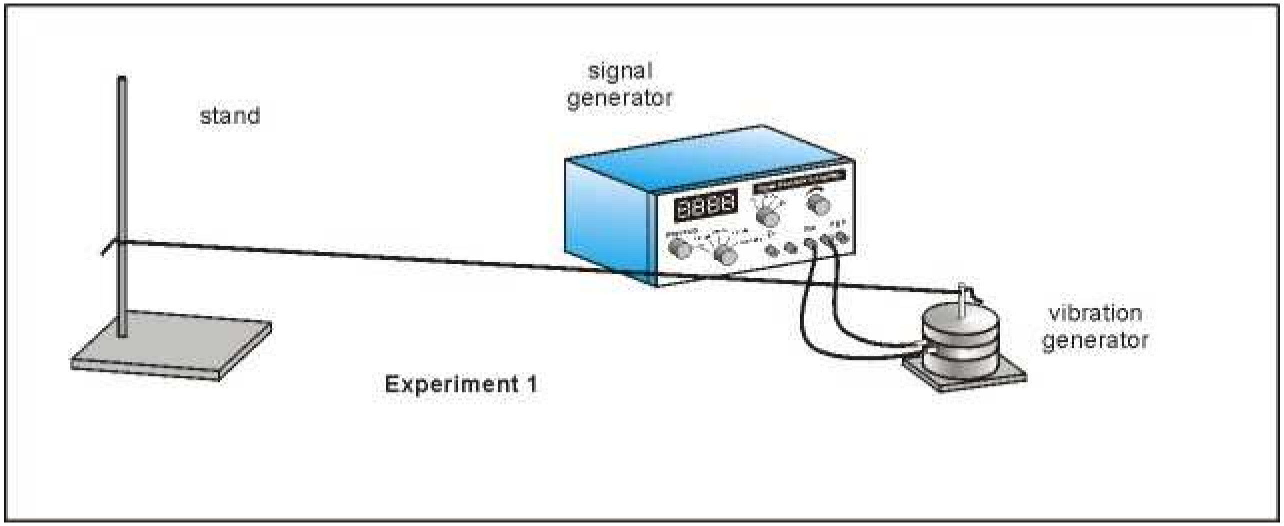
mode	wavelength	frequency
first	$2L$	$\frac{v}{2L}$
second	L	$\frac{v}{L}$
third	$\frac{2L}{3}$	$\frac{3v}{2L}$
fourth	$\frac{L}{2}$	$\frac{2v}{L}$

$$\lambda_n = \frac{2L}{n}$$

$$\omega_n = \frac{n\pi v}{L}$$

$$f_n = \frac{nv}{2L}$$

Forced Oscillations



- One end of the string is fixed, the other end is forced with the function $Y(t) = B\cos\omega t$.

$$y(0, t) = B\cos\omega t$$

$$y(L, t) = 0$$

- The wave equation still holds so we expect solutions to be of the form

$$y(x, t) = f(x)\cos\omega t$$

Forced Oscillations

- This time we can't constrain $f(x)$ to be zero at both ends.
- Now, let $f(x) = A\sin(kx + \alpha)$
  - The constant k is just ω/v.
  - We need to solve for A and α
- Boundary condition at $x = L$:

$$\sin\left(\frac{\omega L}{v} + \alpha\right) = 0 \quad \Rightarrow \quad \frac{\omega L}{v} + \alpha = p\pi$$

$$\alpha_p = p\pi - \frac{\omega L}{v}$$

- Condition at $x = 0$:

$$B = A_p \sin\alpha_p$$

Forced Oscillations

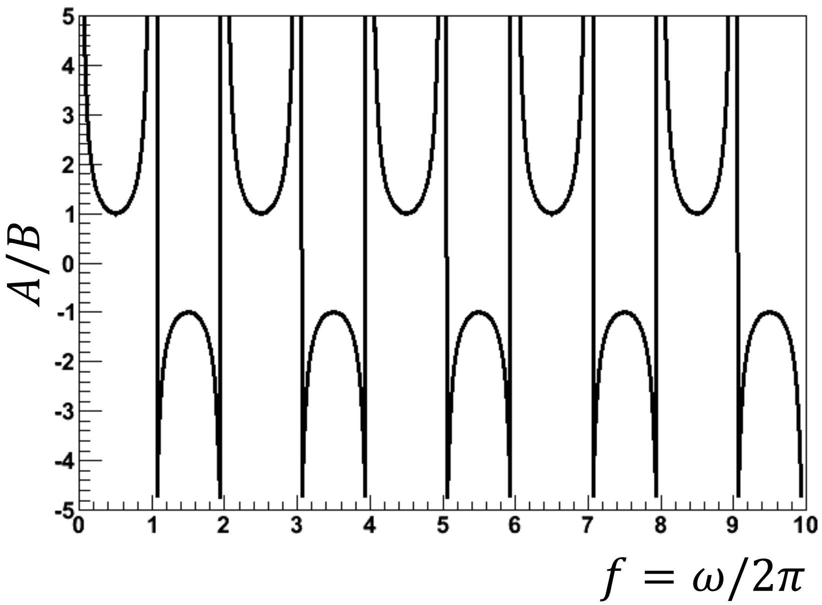
- Amplitude of oscillations:

$$A_p = \frac{B}{\sin(p\pi - \omega L/v)}$$

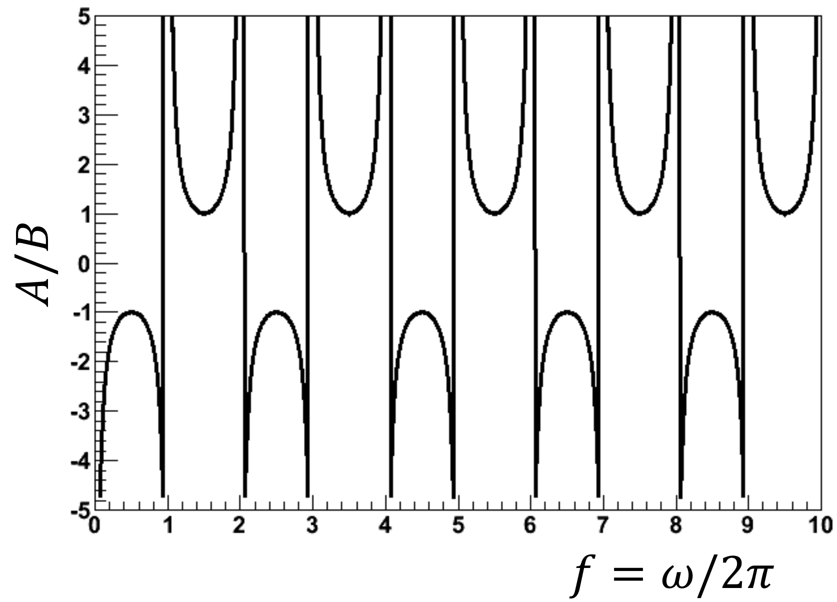
- What does this mean?
  - The driving force can excite many normal modes of oscillation
  - When $\omega = p\pi v/L$, the amplitude gets very large

Forced Oscillations

$p = 1$



$p = 2$



$$L = 5\,m$$

$$v = 10\,m/s$$