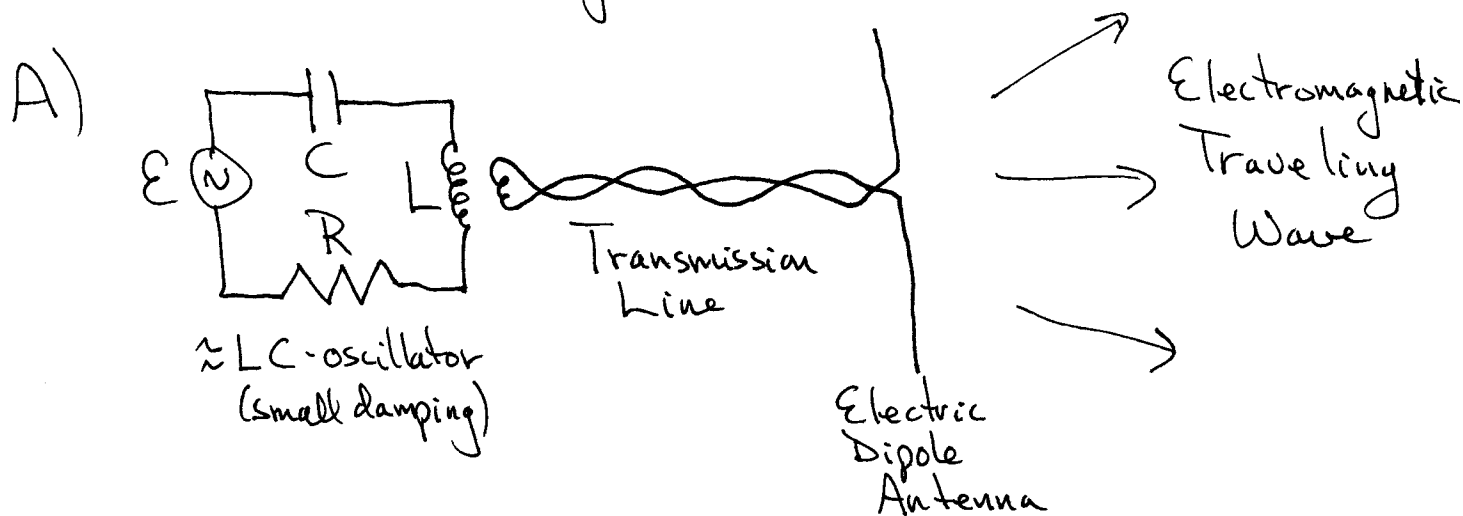
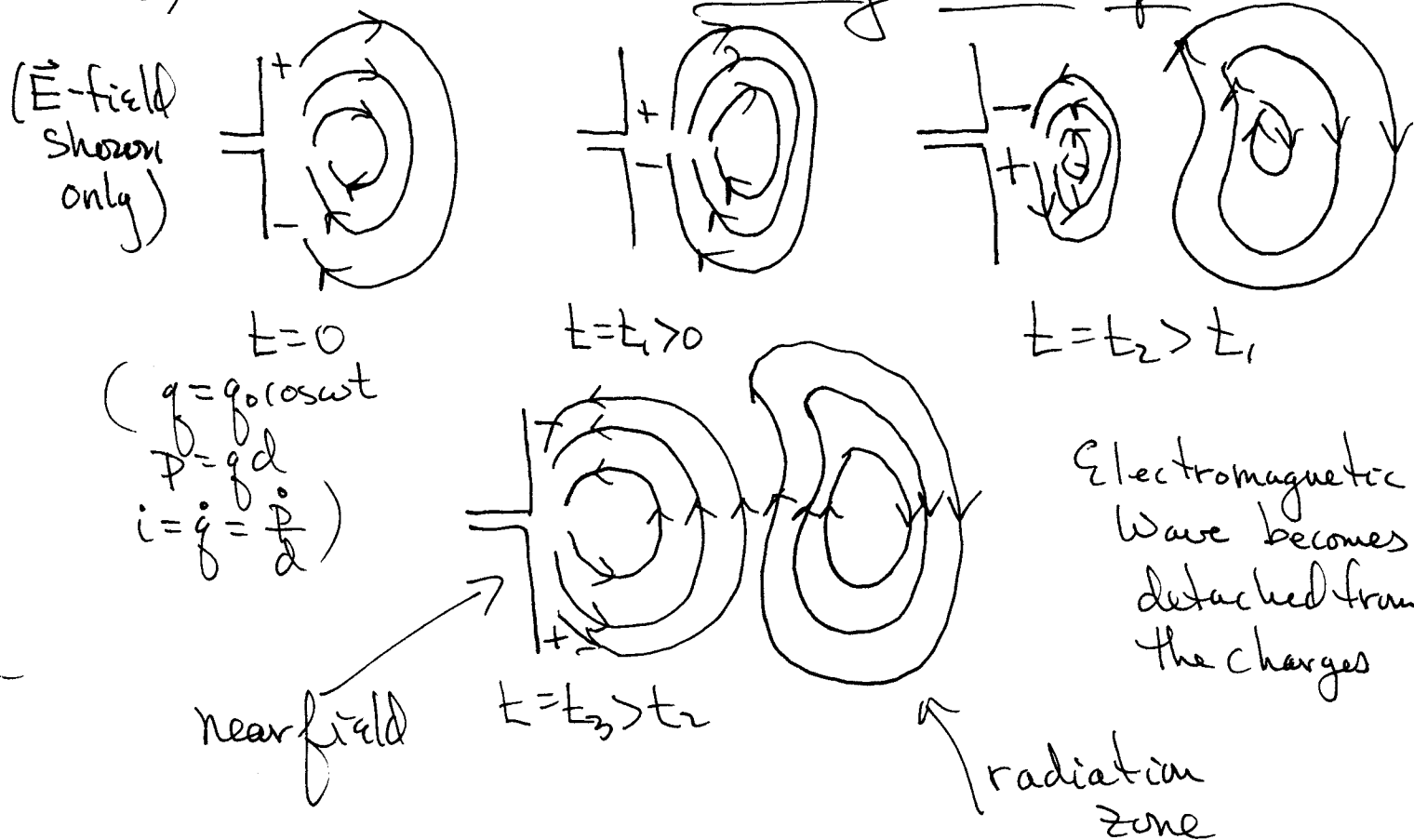


IX) Electromagnetic Radiation: Charges at rest or moving at constant velocity do not radiate.

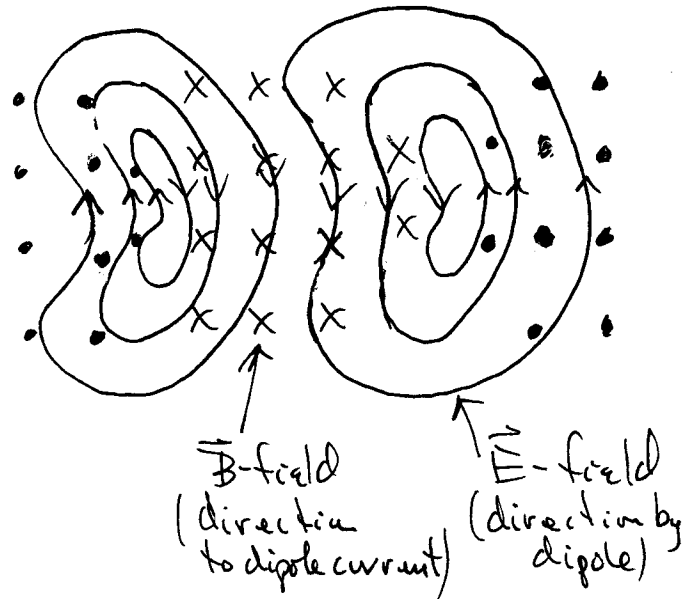
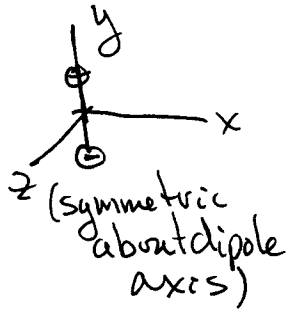
Accelerated charges radiate



B) Antenna acts like oscillating electric dipole moment



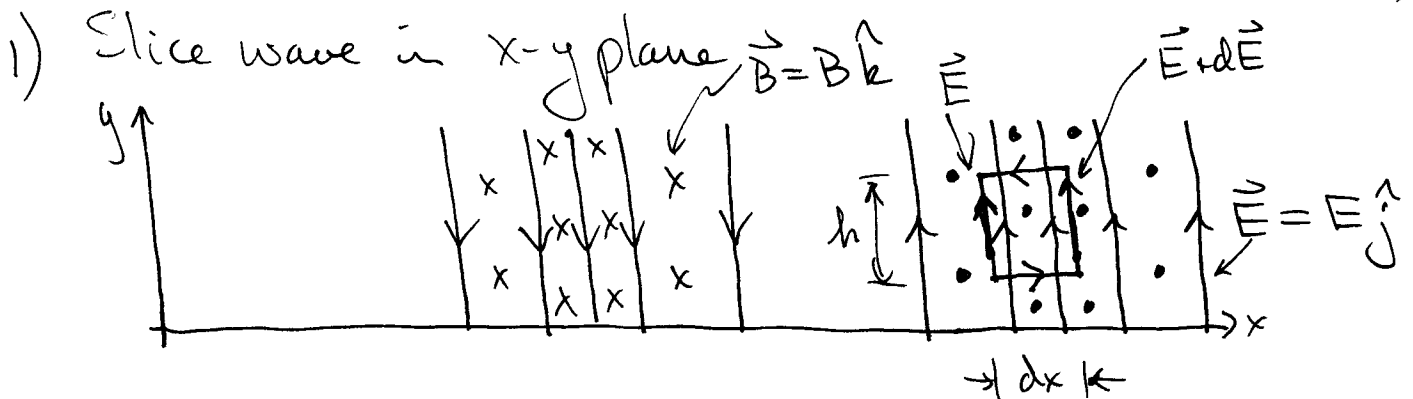
IX) C) Closer look at wave:



D) Maxwell's Equations applied to wave in Rad. zone

$$\oint_C \vec{E} \cdot d\vec{s} = - \int_S \frac{\partial \vec{B}}{\partial t} \cdot d\vec{A} \quad : \text{Faraday's law}$$

$$\oint_C \vec{B} \cdot d\vec{s} = \epsilon_0 \mu_0 \int_S \frac{\partial \vec{E}}{\partial t} \cdot d\vec{A} \quad : \text{Ampere's law (far from dipole in Radiation Zone)}$$



Faraday's Law

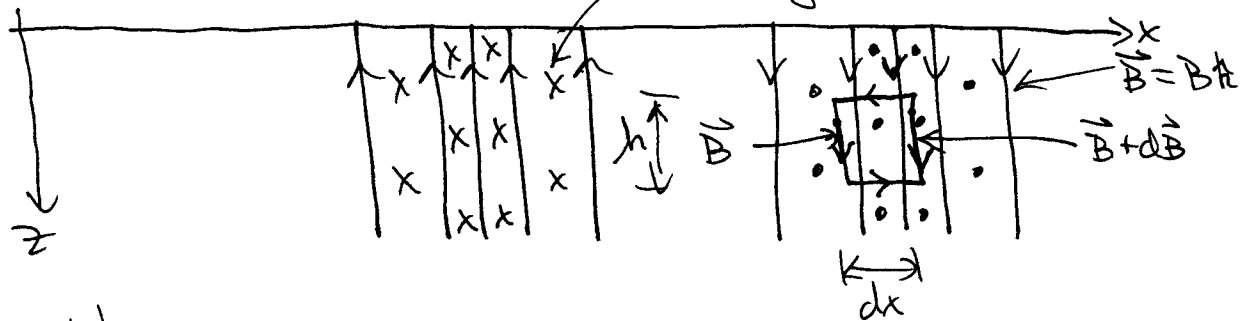
$$\begin{aligned} \oint_C \vec{E} \cdot d\vec{s} &= (E + dE)h - (E)h = dEh \\ &= - \frac{\partial B}{\partial t} (h dx) \end{aligned}$$

IX. D) $\Rightarrow$

$$\frac{\partial E}{\partial x} = - \frac{\partial B}{\partial t}$$

-124-

2) Slice wave in x - z plane $\vec{E} = E \hat{j}$



Ampere's Law:

$$\oint_C \vec{B} \cdot d\vec{s} = -(B + dB)h + Bh = -h dB$$

$$= \epsilon_0 \mu_0 \frac{\partial E}{\partial t} (h dx)$$

$\Rightarrow$

$$\frac{\partial B}{\partial x} = - \epsilon_0 \mu_0 \frac{\partial E}{\partial t}$$

E) Wave Equation: uncouple the above by differentiating again $\Rightarrow$

$$\left[\epsilon_0 \mu_0 \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} \right] E(x, t) = 0$$

$$\left[\epsilon_0 \mu_0 \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} \right] B(x, t) = 0$$

IX. E) General Solution to wave equation

$$\Rightarrow \begin{cases} E(x,t) = E(x-vt) \\ B(x,t) = B(x-vt) \\ v = \frac{1}{\sqrt{\epsilon_0 \mu_0}} \equiv c \end{cases}$$

All electromagnetic waves travel at speed c !

and $\left[\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \frac{\partial^2}{\partial x^2} \right] \begin{matrix} \vec{E} \\ \text{or} \\ \vec{B} \end{matrix} = 0$

F) Plane Waves: Simplest Solution

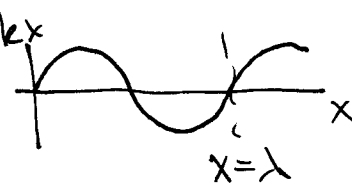
$$\begin{aligned} E(x,t) &= E_m \sin(kx - \omega t) \\ B(x,t) &= B_m \sin(kx - \omega t) \end{aligned}$$

Wave Equation $\Rightarrow \frac{\omega}{k} = c$

k = wave number

$kx = 2\pi$ gives zeroes of $\sin kx$

λ is called wave length



$$k\lambda = 2\pi$$

$$\Rightarrow \boxed{k = \frac{2\pi}{\lambda}}$$

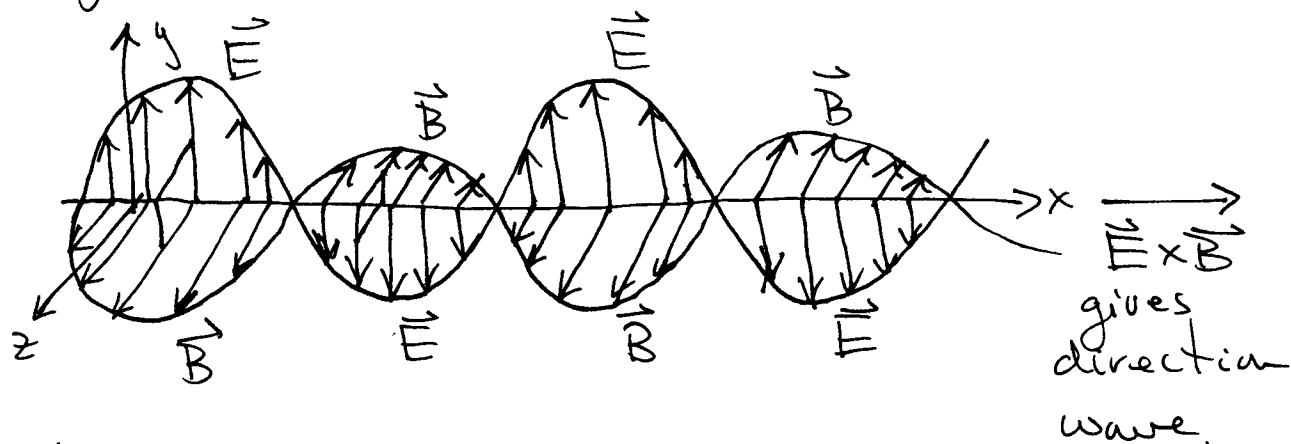
So

$$\boxed{\frac{\omega}{k} = c = \frac{\omega}{2\pi} \frac{2\pi}{k} = \lambda \lambda}$$

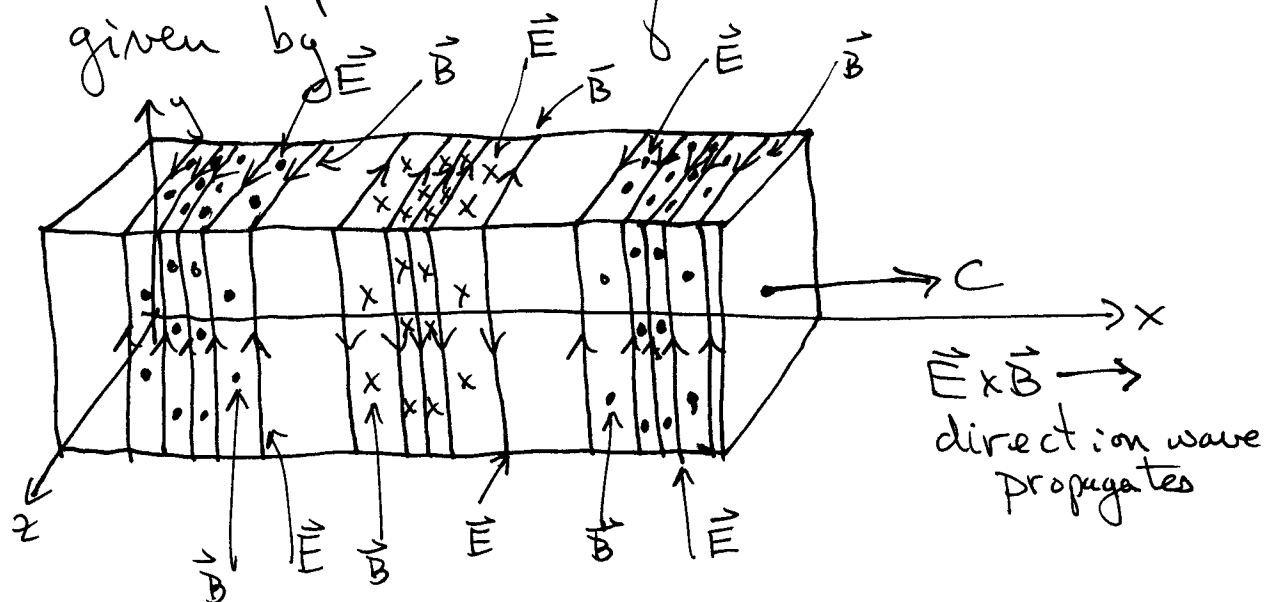
-IX. G) Linearly Polarized : $\vec{E} = E \hat{y}$
 $\vec{B} = B \hat{z}$

$\vec{E}$ oscillates along one direction only, as does $\vec{B}$.

For fixed t this looks like



Another representation of the wave is
 given by



IX. H) Maxwell's Equations: $\frac{\partial E}{\partial x} = -\frac{\partial B}{\partial t}$: Faraday's Law
 $\frac{\partial B}{\partial x} = -\frac{1}{c^2} \frac{\partial E}{\partial t}$: Ampere's Law

$\Rightarrow$ 1) E & B must be in phase

2) $E = cB$

I) Energy Transport: The Poynting Vector

$$\vec{S} = \vec{E} \times \vec{H}$$

$$= \frac{1}{\mu_0} \vec{E} \times \vec{B} \quad \text{in vacuum}$$

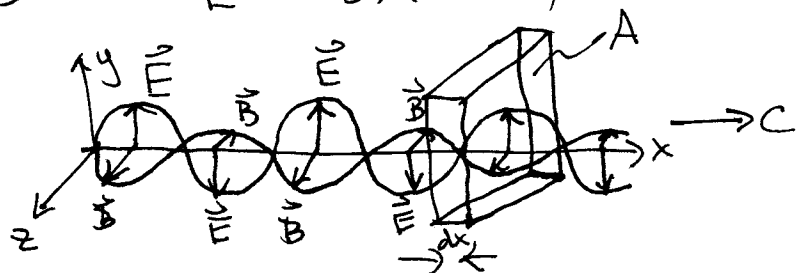
$\vec{S}$ = rate of energy flow per unit area in the direction of $\vec{S}$
 = electromagnetic power per unit area.

1) For above plane wave $\vec{S} = \frac{1}{\mu_0} EB \hat{i} \equiv S \hat{i}$
 $\uparrow$ gives direction of propagation

2) $dU = dU_E + dU_B = (u_E + u_B)(A dx)$

$$u_E = \frac{1}{2} \epsilon_0 E^2$$

$$u_B = \frac{1}{2\mu_0} B^2$$



-IX.I.)2) $dU = \frac{EBA}{\mu_0} \frac{dx}{c}$ but $dt = \frac{dx}{c}$ -128-

This is amount of energy that passes through the box in time dt , the energy flow per unit time thru area

$$A : \frac{dU}{dt} = \frac{1}{\mu_0} EBA = S \cdot A$$

So $S = \frac{dU}{dt A} = \frac{1}{\mu_0} EB$ is the energy flows per unit time per unit area.

3) Intensity $I = \bar{S}$ = time average of S

$$I = \bar{S} = \frac{1}{2\mu_0} E_m B_m$$

$$= \frac{1}{\mu_0} E_{rms} B_{rms}$$

