

Sec. 6.10 Transformation properties of $\vec{E}, \vec{B}, \Phi, \vec{A}, \rho, \vec{J}, \dots$

Observe - For the equations interrelating various quantities like $\vec{J}, \vec{B}, \frac{\partial \vec{E}}{\partial t}$, etc to make sense, the terms in an equation must ALL behave the same under the operations of

- coordinate rotations
- coordinate inversions
- time-reversal

A) Rotations $\Rightarrow$ These are orthogonal transformations that replace the original coordinates x_β by transformed coordinates x'_α

$$\Rightarrow \boxed{x'_\alpha = \sum_\beta a_{\alpha\beta} x_\beta} \quad (*)$$

where orthogonality is expressed by

$$\underline{a}^T \underline{a} = \underline{I}, \text{ or } \sum_\beta (\underline{a}^T)_{\alpha\beta} a_{\beta\gamma} = \delta_{\alpha\gamma}$$

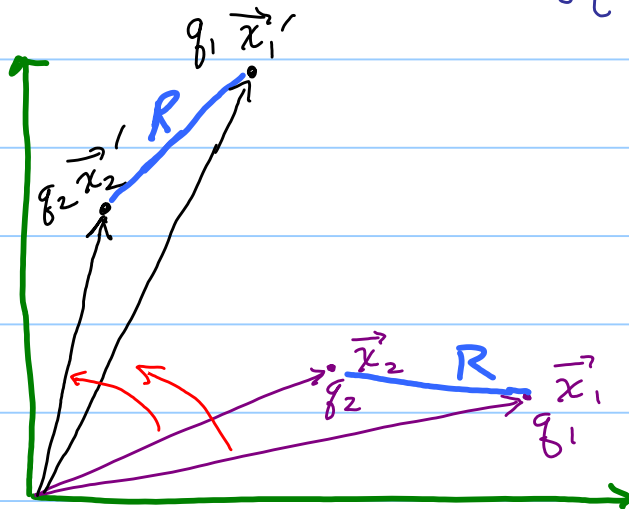
which ensures that $\vec{x}'^2 = \vec{x}^2$

And if $\det(\underline{a}) = +1$, the rotation can be replaced by a sequence of infinitesimal steps and is a "PROPER ROTATION"

While if $\det(\underline{a}) = -1$, it is an "IMPROPER ROTATION" which requires a reflection plus a rotation.

- Positions $\vec{x}_i$, Velocities $\vec{v}_i$, momenta $\vec{p}_i$ all transform according to (*) and these are **VECTORS** = tensors of rank 1
- Scalar products like $\vec{x}_1 \cdot \vec{x}_2$ or $\vec{x}_1 \cdot \vec{v}_2$ are invariant under rotations and are called **SCALARS** = tensors of rank 0
- The Maxwell Stress Tensor is an example of a rank 2 tensor, i.e.

$$\overleftrightarrow{T}'_{\alpha\beta} = \sum_{\gamma\delta} a_{\alpha\gamma} \overleftrightarrow{T}_{\gamma\delta} (a^T)_{\delta\beta}$$



e.g. an active rotation of all physical vectors through the same angle preserves scalars such as $\frac{1}{|\vec{x}'_1 - \vec{x}'_2|}$

Behavior of differential operators under rotation:

$\nabla \Phi = \text{vector}$, assuming $\Phi = \text{scalar}$

$\nabla \cdot \vec{V} = \text{scalar}$, if $\vec{V} = \text{vector}$

$\nabla \cdot \overleftrightarrow{T} = \text{vector}$, if $\overleftrightarrow{T} = 2^{\text{nd}}$ -rank tensor

$\nabla^2 \Phi = \text{scalar}$

$\nabla^2 \vec{A} = \text{vector}$, etc....

B) Spatial inversions and reflections

→ **INVERSION (I)** changes the sign of all 3 coordinates of a position vector, i.e.

$$\vec{r} = x\hat{x} + y\hat{y} + z\hat{z}$$

$$\text{and } I \vec{r} = -x\hat{x} - y\hat{y} - z\hat{z}$$

REFLECTION (R) reflects through one chosen plane, i.e.

$$R_y \vec{r} = x\hat{x} - y\hat{y} + z\hat{z}$$

⇒ Neither R nor I can be accomplished by simple (i.e. **PROPER**) rotations.

Polar vectors (or vectors) — these transform like (*) under proper rotations, i.e. $V'_\alpha = \sum_{\beta} a_{\alpha\beta} V_\beta$

AND they change sign under inversion, i.e. $\vec{V}' = I \vec{V} = -\vec{V}$

Axial vectors (or pseudovectors) - these rotate the same way (*) but they are SYMMETRIC under inversion

$$\text{i.e. } \vec{L}' = \mathbf{I} \vec{L} = +\vec{L}$$

e.g. a cross product of 2 polar vectors,

$$\mathbf{I} (\vec{x} \times \vec{p}) = (-\vec{x}) \times (-\vec{p}) = \vec{x} \times \vec{p}$$

Note that pseudovectors and pseudoscalars imply a "TWIST" somewhere in the physics!

Pseudoscalars are ODD under inversion:

$$\begin{aligned} \mathbf{I} [\vec{x}_1 \cdot (\vec{x}_2 \times \vec{p}_2)] &= [(-\vec{x}_1) \cdot (-\vec{x}_2) \times (-\vec{p}_2)] \\ &= -\vec{x}_1 \cdot (\vec{x}_2 \times \vec{p}_2) \end{aligned}$$

C) TIME REVERSAL - the classical microscopic laws of physics are invariant if the direction of time is reversed, i.e. if time runs backwards, $\Rightarrow T' dt = -dt'$

Electromagnetic entities

<u>Entity</u>	<u>Proper rotations</u>	<u>Inversions</u>	<u>Time Reversal</u>	<u>Overall</u>
charge	scalar	even	even	true scalar
E-field ($\nabla \cdot \vec{E} = \rho/\epsilon_0$)	vector	odd	even	polar vector
B-field ($\nabla \times \vec{E} + \frac{\partial \vec{B}}{\partial t} = 0$)	vector	even	odd	pseudovector
Current $\vec{J} = \rho \vec{v}$	vector	odd	odd	polar vector

see Table 6.1, p 271 for more

Question: What are the most general relations among $\vec{E}, \vec{B}$ and their sources, consistent with $\left\{ \begin{array}{l} \text{rotational invariance} \\ \text{inversions} \\ \text{time-reversal invariance} \end{array} \right.$?

Answer: suppose the relationships are linear in $\vec{E}, \vec{B}$ and first derivatives w.r.t. $\vec{x}, t$ and linear in the sources

Notation: call R = rotations

I = inversions

T = time-reversal

Behavior under,

R	I	T
vector	odd	even

vector	odd	odd
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pseudovector	even	odd
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pseudovector	even	even
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scalar	even	even
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pseudoscalar	odd	odd
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entities

$\vec{E}$

$\frac{\partial \vec{E}}{\partial t}, \nabla \times \vec{B}, \vec{J}$

$\vec{B}$

$\frac{\partial \vec{B}}{\partial t}, \nabla \times \vec{E}$

$\nabla \cdot \vec{E}, \rho$

$\nabla \cdot \vec{B}$

$\vec{J}_m = \rho_m \vec{v}$
magnetic
monopole
current density
if ever found

(and ρ_m = magnetic
monopole charge if
ever found)

This table gives deeper insight into why Maxwell's Equations are what they are!