

THE EXPONENTIAL FUNCTION e^x

10

As with most mathematical functions there are many ways to introduce/derive e^x . For our purposes a simple way (among many!) is to define e^x by the series

$$e^x \equiv \sum_{n=0}^{\infty} \frac{x^n}{n!} = 1 + \frac{x}{1!} + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots \quad (1)$$

Evidently the number $e = e^1 = 1 + \frac{1}{1!} + \frac{1}{2!} + \frac{1}{3!} + \dots \quad (2)$

$$\underbrace{1 + 1 + \frac{1}{2} + \frac{1}{6} + \dots}_{2.666\dots} = 2.718281828\dots \quad (3)$$

(a transcendental number)
- like π

Why is this interesting? Consider the derivative of the series (1):

$$\begin{aligned} \frac{d}{dx} e^x &= 0 + 1 + \frac{2x}{2!} + \frac{3x^2}{3!} + \frac{4x^3}{4!} + \dots \\ &= 1 + \cancel{\frac{2x}{2\cdot 1}} + \cancel{\frac{3x^2}{3\cdot 2\cdot 1}} + \cancel{\frac{4x^3}{4\cdot 3\cdot 2\cdot 1}} + \dots = 1 + x + \frac{x^2}{2!} + \frac{x^3}{3!} + \dots \end{aligned} \quad (4)$$

Hence $\frac{d}{dx} e^x = e^x$ (5)

The function e^x can also be arrived at in other ways.
Consider the function $f(x)$ defined by

$$f(x) = \lim_{N \rightarrow \infty} \left(1 + \frac{x}{N}\right)^N \quad (6)$$

and use the chain rule: $\frac{d}{dx} f(u(x)) = \frac{df}{du} \cdot \frac{du}{dx} \quad (7)$

$$\text{Let } u(x) = \left(1 + \frac{x}{N}\right); u^N = \left(1 + \frac{x}{N}\right)^N \quad (8)$$

11

$$\text{Then } \frac{df(x)}{dx} = \lim_{N \rightarrow \infty} \frac{d}{dx}(u^N) = \lim_{N \rightarrow \infty} (Nu^{N-1}) \frac{du}{dx} \xrightarrow{\left(\frac{1}{N}\right)} \frac{1}{N} \quad (9)$$

$$\frac{df(x)}{dx} = \lim_{N \rightarrow \infty} \left\{ (Nu^{N-1}) \left(\frac{1}{N}\right) = u^{N-1} = \frac{u^N}{u} \right\} \quad (10)$$

$$= \lim_{N \rightarrow \infty} \left\{ \frac{\left(1 + \frac{x}{N}\right)^N}{1 + \frac{x}{N}} \right\} \quad (11)$$

In the limit $N \rightarrow \infty$ we can approximate $1 + \frac{x}{N} \approx 1$. Hence (11) \Rightarrow

$$\boxed{\frac{df(x)}{dx} = \lim_{N \rightarrow \infty} \left(1 + \frac{x}{N}\right)^N = f(x)} \quad (12)$$

Comparing (12) & (5) we see that this is the same differential equation as for e^x . Hence it must be that $f(x) = e^x + \text{constant}$.

$$\text{But } e^0 = 1 \text{ and } f(0) = 1 \text{ also } \Rightarrow f(x) = e^x \quad (13)$$

INTERPRETATION/APPLICATION OF (12) & (13):

Suppose you deposit $d\$\#$ into your bank account on January 1. Let x be the interest rate, for example $x = 8\% = 0.08$. If interest is evaluated at the end of the year the value of your account is

$$V(\text{Dec. 31}) = d(1.08) \quad (14)$$

Suppose, however that the interest is credited twice/year at the rate of $\frac{1}{2} \times 0.08 = 0.04$ each time. Then the value of your account would now be:

$$V(\text{Dec. 31}) = \$d(1.04)(1.04) = \$1(1+2 \times 0.04 + 0.0016)$$

$$= \$d(1.0816) \quad (15)$$

12

So you are now slightly better off due to the more frequent compounding. Continuing, suppose interest is compounded quarterly. Then,

$$V(\text{Dec. 31}) = \$d(1.02)(1.02)(1.02)(1.02) = \$d1.0824 \quad (16)$$

$$= \$d\left(1+\frac{0.08}{4}\right)^4$$

So more frequent compounding helps!!

However, this does not go on forever. In the limit as compounding happens infinitely often then,

$$V(\text{Dec. 31}) = \$d \lim_{N \rightarrow \infty} \left(1 + \frac{x}{N}\right)^N \xrightarrow{x=0.08} \$d \cdot e^x = \$d \cdot e^{0.08}$$

$$= \$d(1.08329)$$

THE "RULE OF 72"

This is a simple mnemonic to calculate how long it will take to double how much money you have, given an assumed constant rate of growth r ($\%$ per year).

At $t=0$ you have $\$d$. After 1 year you will then have $\$d(1+r)$, After n years you will then have $\$d(1+r)^n$.

Hence we want to find n such that

$$\$d(1+r)^n = 2\$d \Rightarrow (1+r)^n = 2 \Rightarrow n \ln(1+r) = \ln 2 \quad (1)$$

If we assume a conventionally modest value of $r \ll 1$, we can then expand,

$$\ln(1+r) \approx r \Rightarrow n \ln(1+r) \approx nr = \ln 2 \approx 0.693 \quad (2)$$

Hence

$$n = \frac{0.693}{r} = \frac{69.3}{100r} \approx \frac{72}{R} \quad R = \text{growth rate in } 100 \times 0.0 \quad (69\% \rightarrow 6)$$

Final result:

$$n \approx \frac{72}{R} = \text{"rule of 72"} \quad (4)$$

Examples:

$$r = 4\% \Rightarrow R = 4$$

$$\text{Rule (4)}$$

$$n = 18$$

$$\text{Exact (1)}$$

$$17.7$$

$$r = 6\% \Rightarrow R = 6$$

$$n = 12$$

$$11.9$$

$$r = 8\% \Rightarrow R = 8$$

$$n = 9$$

$$9.006$$

$$r = 9\% \Rightarrow R = 9$$

$$n = 8$$

$$8.04$$

$$r = 12\% \Rightarrow R = 12$$

$$n = 6$$

$$6.11$$

APPLICATION TO RADIOACTIVE DECAY AND POPULATION GROWTH

-13.1

Natural radioactivity was discovered in 1896 by Henri Becquerel. It was indicated by subsequent experiments that if one started with a sample of N_0 atoms at $t=0$ then the number remaining at a time $t > t=0$ is given by

$$N(t) = N_0 e^{-\lambda t} \Rightarrow \frac{dN(t)}{dt} = -\lambda \underbrace{N_0 e^{-\lambda t}}_{N(t)} = -\lambda N(t) \quad (1)$$

$$\Rightarrow \boxed{-\frac{dN(t)}{N(t)} = \lambda dt} \quad (2) \quad \lambda \equiv \text{decay constant (units of 1/time)}$$

Half Life $\equiv T_{1/2}$; this is the time at which $N(t)/N_0 = 1/2$.

$$\text{From (1) \& (2): } \frac{N(t)}{N_0} = e^{-\lambda t} \Rightarrow \underbrace{\ln(1/2)}_{\ln(1)-\ln(2)} = -\lambda t \quad \left. \begin{array}{l} \ln(1)=0 \\ \ln(2)=-\ln(2) \end{array} \right\} -\ln(2) = -\lambda t \equiv -\lambda T_{1/2}$$

$$\Rightarrow \boxed{T_{1/2} = \frac{\ln 2}{\lambda} \approx \frac{0.693}{\lambda}} \quad (3)$$

- Half-lives of radioactive isotopes range from fractions of a second to tens of billions of years.
- There are recent indications that λ itself (and hence $T_{1/2}$) is not a constant. See "THROUGH THE WORMHOLE WITH MORGAN FREEMAN", July 13, 2011 story.
- For population growth Eq.(1) also applies with a change in sign: From (2) this gives

$$\boxed{\frac{dN(t)}{N(t)} = +\lambda dt} \quad (4)$$

INTRODUCTION TO $\ln(x)$

(natural logarithms of x)

As with other functions, $\ln(x) \equiv \ln x$ can be introduced in a number of ways. We choose to define $\ln x$ as the inverse of e^x as defined by

$$x = e^{\ln x} \quad (1)$$

$x \rightarrow \boxed{\ln x} \rightarrow \boxed{e^{\ln x}} \rightarrow x$

Then for $x=e \Rightarrow e=e^1=e^{\ln e} \Rightarrow \ln e=1$. We know how to differentiate e^x , and we can use this to differentiate $\ln x$:

Start with the chain rule $\frac{d}{dx} f(u(x)) = \frac{df}{du} \cdot \frac{du}{dx}$ "cancel"

Then let $f(u) = e^u$; $u = \ln x$

From (1) & (2)

$$\underbrace{\frac{d}{dx} \{x\}}_{\text{1}} = \frac{d}{dx} \{e^{\ln x}\}$$

$$1 = \underbrace{e^{\ln x}}_x \cdot \frac{d \ln x}{dx} = x \cdot \frac{d \ln x}{dx} \Rightarrow \boxed{\frac{d \ln x}{dx} = \frac{1}{x}} \quad (3)$$

Continuing in this way we find:

$$\frac{d}{dx} \ln x \equiv D \ln x = \frac{1}{x}; \quad \frac{d^2}{dx^2} \ln x = \frac{-1}{x^2} \equiv D^2 \ln x; \quad D^3 \ln x = \frac{2}{x^3}, \dots \quad (4)$$

$$D^n \ln x = \frac{(-1)^{n+1} (n-1)!}{x^n} \quad (5)$$

MORE ON LOGARITHMS ("logs")

14.1

From our previous definitions of $\ln a$:

$$a = e^{\ln a} \quad b = e^{\ln b} \quad (1)$$

$$\Rightarrow ab = e^{\ln a} \cdot e^{\ln b} = e^{\ln a + \ln b} \equiv e^{\ln(ab)} \quad (2)$$

But, whatever ab is, it is just another number so that

Hence it must be that

$$\ln(ab) = \ln(a) + \ln(b) \quad (3)$$

↳ this is how a slide rule works!

From (1) we can also write : $a^x = (e^{\ln a})^x = e^{x \ln a} \quad (4)$

This will be important later in the semester when we discuss complex variables: If $z = x+iy$ we can determine the properties of e^z and $\ln z$. Using (4) we can then write

$$z_1^{z_2} = e^{z_2 \ln z_1} \quad (5)$$

For example $i^i = e^{i \ln i} = e^{-\pi/2} = \text{REAL!} \quad (6)$

CHANGE OF BASE:

$$\rightarrow 1000 = 10^3 \text{ etc.}$$

e = base for natural logs. 10 = base for common logs

2 = base for binary numbers. For any base b

$$a = b^{\log_b a} \quad (7)$$

$$1000 = 10^{\log_{10}(1000)} = 10^3$$

$$3 = \log_{10}(1000)$$

A number in one base can be re-expressed in terms of another base:

$$y = e^{\ln y} \equiv b^{\log_b y} \quad b = \text{another base}$$

$$\downarrow b = e^{\ln b}$$

$$\Rightarrow y = (e^{\ln b})^{\log_b y} = e^{\ln b \cdot \log_b y}; \ln \equiv \log_e \quad (8)$$

$$\text{But } y = e^{\ln y} \Leftrightarrow \ln y = \ln b \cdot \log_b y \Rightarrow \log_b y = \frac{\ln y}{\ln b}$$

The two most common bases are e and 10

$$= \frac{\log_e y}{\log_e b} \quad (9)$$

\ln is the shorthand for \log_e = "natural logs"

\log_{10} = "common logs"

Derivative of (Any Number)^{any power}:

$$\frac{d}{dx} x^p = \frac{d}{dx} (e^{\ln x})^p = \frac{d}{dx} e^{p \ln x} = e^{p \ln x} \cdot \underbrace{\frac{d}{dx}(p \ln x)}_p \quad (10)$$

$$\Rightarrow \frac{d}{dx} x^p = \frac{p}{x} \cdot \underbrace{e^{p \ln x}}_{x^p} = p x^{p-1} \quad (11)$$

This is the usual result, but now we have shown that it holds for any x and any p . Later we will show that it also holds when $x \rightarrow z = \text{complex}$, and p is also complex.

This illustrates a general technique for generalizing simpler results,

TAYLOR SERIES

15

Suppose that we know the value of a function $f(x)$ at some point x , but we would like to be able to compute the value at a nearby point $(x+a)$. This can be done via a Taylor Series which can be written in a number of equivalent ways.

A formula (useful in Quantum Mechanics - see below!!) is

INCREMENT FORM

$$f(x+a) = e^{a \frac{d}{dx}} f(x) = \left(1 + a \frac{d}{dx} + \frac{a^2}{2!} \frac{d^2}{dx^2} + \dots \right) f(x) \quad (1)$$

Equivalently by interchanging $x \leftrightarrow a$,

$$f(a+x) = e^{x \frac{d}{da}} f(a) = \left(1 + x \frac{d}{da} + \frac{x^2}{2!} \frac{d^2}{da^2} + \dots \right) f(a) \quad (2)$$

If we now shift x by replacing $x \rightarrow x-a$ then (2) \Rightarrow

$$f(a+x-a) = f(x) = \left(1 + \frac{(x-a)}{1!} \frac{d}{da} + \frac{(x-a)^2}{2!} \frac{d^2}{da^2} + \dots \right) f(a) \quad (3)$$

TEXT
(1.3, 44)

$$\therefore f(x) = f(a) + \frac{(x-a)}{1!} f'(a) + \frac{(x-a)^2}{2!} f''(a) + \dots \quad (4)$$

Both of these forms are useful in different contexts.

Connection to Quantum Mechanics: ($\hbar = \text{Planck's constant}$)

In (1) write $a \frac{d}{dx} = \underbrace{\frac{ia}{\hbar}}_{\text{momentum operator} \equiv P} \cdot \underbrace{\frac{\hbar}{i} \frac{d}{dx}}_{= P} = \frac{iaP}{\hbar}$ (5)

Hence (1) & (5) \Rightarrow

$$e^{iaP/\hbar} f(x) = f(x+a) \quad (6)$$

The momentum operator induces spatial translations

The fact that the momentum operator "moves" (translates) the function $f(x) \rightarrow f(x+a)$ is intimately associated with the

HEISENBERG UNCERTAINTY PRINCIPLE

$$\Delta p \Delta x \gtrsim \hbar \quad (7)$$

The close connection between momentum conservation and spatial translations is part of NOETHER'S THEOREM.

APPLICATIONS OF TAYLOR SERIES:

$$(1) \quad f(x) = \frac{1}{1-x} = \frac{1}{u} \quad (8)$$

$$f(x) = f(a) + \frac{(x-a)}{1!} f'(a) + \frac{(x-a)^2}{2!} f''(a) + \dots \quad \text{choose } a=0 \quad (9)$$

$$f'(a) = f'(a) = \left. \frac{df}{du} \cdot \frac{du}{dx} \right|_{x=0} = \left. -\frac{1}{u^2} \cdot (-1) \right|_{x=0} = \left. \frac{1}{(1-x)^2} \right|_{x=0} = +1 \quad (10)$$

$$f''(a) = f''(a) = \left. \frac{-2}{(1-x)^3} \cdot (-1) \right|_{x=0} = +2 \quad (11)$$

$$\therefore f(x) = \frac{1}{1-x} = \underbrace{f(0)}_1 + \frac{(x-0)}{1!} \cdot 1 + \frac{(x-0)^2}{2!} \cdot 2 + \dots = 1+x+x^2+\dots \quad (12)$$

We will show later that for $|x| < 1$ this series converges.

Hence for $|x| < 1$ we can write

$$\frac{1}{1-x} = 1+x+x^2+\dots = \sum_{n=0}^{\infty} x^n \quad |x| < 1 \quad (13)$$

$$\text{Also } \frac{1}{1+x} = 1-x+x^2-x^3+\dots = \sum_{n=0}^{\infty} (-1)^n x^n \quad |x| < 1$$

[2]

$$f(x) = \ln x$$

17

$$f(x) = f(a) + \frac{(x-a)}{1!} f'(a) + \frac{(x-a)^2}{2!} f''(a) + \dots \quad (14)$$

Again the Taylor series starts by knowing a value a at which we know the function. Here we can choose $x=1 \Rightarrow \ln 1 = 0$. Also from our previous results, for any value a

$$D^n \ln a = \frac{(-1)^{n+1} (n-1)!}{a^n} \xrightarrow{a=1} (-1)^{n+1} (n-1)! \xrightarrow{0!=1} \quad (15)$$

Hence the first few terms in the series expansion of $\ln x$ are

$$\ln x = \underbrace{\ln 1}_0 + (x-1) - \frac{(x-1)^2}{2} + \frac{(x-1)^3}{3} - \dots \quad (16)$$

This series converges for $|x-1| < 1 \Rightarrow 0 < x < 2$. Another useful form for the expansion of $\ln x$ follows from the replacement $y = x-1 \Rightarrow x = 1+y$. Then (16) \Rightarrow

$$\ln(1+y) = y - \frac{y^2}{2} + \frac{y^3}{3} - \frac{y^4}{4} + \dots \quad (17)$$

This formula is used in our derivation of The RULE OF 72: $\ln(1+y) \approx y$

[3]

$f(x) = (1+x)^n$ Typical applications are for small x near 1 so we expand about $f(a)=1$, which corresponds to $x=a=0$. Then

$$f(x) = \underbrace{f(a)}_{f(x=a)} + \frac{(x-a)}{1!} f'(x=a) + \frac{(x-a)^2}{2!} f''(x=a) + \dots \quad (18)$$

Note that (as a notation check) $f(x=a) = f(a) + 0 + 0 + 0 \dots$ which is as expected.

$$f'(x) = \frac{df}{du} \cdot \frac{du}{dx} ; u = 1+x ; \frac{du}{dx} = 1 \quad (19)$$

18

$$\Rightarrow f'(x) = n u^{n-1} \cdot 1 = n(1+x)^{n-1} \Rightarrow f'(x=a=0) = n \quad (20)$$

$$f''(x) = n(n-1) u^{n-2} \cdot 1 \Rightarrow f''(x=a=0) = n(n-1) \dots \quad (21)$$

$$\text{Hence } (1+x)^n = 1 + \frac{(x-0)}{1!} \cdot n + \frac{(x-0)^2}{2!} n(n-1) + \dots \quad (22)$$

$$\therefore \text{To leading order, } (1+x)^n \approx 1+nx \quad (23)$$

$$\text{More generally: } (1\pm x)^n = 1 \pm \frac{nx}{1!} + \frac{n(n-1)x^2}{2!} \pm \frac{n(n-1)(n-2)x^3}{3!} + \dots \quad (24)$$

This series converges for $x^2 < 1$.

$$\underline{\text{CHECK!}} \quad \text{For } n=-1, (1\pm x)^{-1} = \frac{1}{1\pm x} = 1 \mp x + x^2 \mp x^3 \quad (25)$$

This agrees with the previous results in (13),

The preceding results also hold if n is replaced by a rational power ϕ , such as $\phi = 1/2$. This is useful, for example in the theory of relativity where we often encounter the expressions

$\sqrt{1-v^2/c^2}$ and $1/\sqrt{1-v^2/c^2}$. Let $x = v^2/c^2$, then

$$\sqrt{1-v^2/c^2} \Rightarrow (1\mp x)^{1/2} = 1 \mp \frac{1}{2}x - \frac{1 \cdot 1}{2 \cdot 4} x^2 \mp \frac{1 \cdot 1 \cdot 3}{2 \cdot 4 \cdot 6} x^3 \dots \quad (26)$$

$$\therefore (1\mp x)^{1/2} = 1 \mp \frac{1}{2}x - \frac{1}{8}x^2 \mp \frac{1}{16}x^3 \dots \quad \boxed{\text{GR 1.112}} \quad (27)$$

$$\text{Also: } (1\mp x)^{-1/2} = \frac{1}{\sqrt{1\mp x}} = 1 \pm \frac{1}{2}x + \frac{3}{8}x^2 \pm \frac{5}{16}x^3 \dots \quad (28)$$

$$\text{Hence } \frac{mc^2}{\sqrt{1-v^2/c^2}} \approx mc^2 \left(1 + \frac{1}{2} \frac{v^2}{c^2}\right) \approx mc^2 + \frac{1}{2}mv^2 \quad (29)$$

COMPARISON OF THE TWO FORMS OF THE TAYLOR SERIES

19

INCREMENT
FORM

$f(x) = \text{Known then}$

$$f(x+a) = e^{\frac{a}{\Delta x}} f(x) = \left(1 + \frac{a}{1!} \frac{d}{dx} + \frac{a^2}{2!} \frac{d^2}{dx^2} + \dots\right) f(x) \quad (1)$$

$$f(x+a) = f(x) + \frac{a}{1!} f'(x) + \frac{a^2}{2!} f''(x) + \dots \quad (2)$$

In this form we assume that we know the function $f(x)$ at some point x , and this allows us to evaluate the r.h.s. exactly.

Then if we know $f(x)$ we can compute $f(x+a)$,

TAYLOR SERIES
FORM

$$f(x) = f(a) + \frac{(x-a)}{1!} f'(a) + \frac{(x-a)^2}{2!} f''(a) + \dots \quad (3)$$

In this form we again assume that we know the function $f(x)$, which now appears on the l.h.s. of (3). Here the main objective is to construct a series expansion of $f(x)$ in terms

of the (presumably small) quantities $(x-a)$, $(x-a)^2$, ...

A very useful compendium of formulas for series is:

I.S. Gradshteyn & I.M. Ryzhik

"TABLE OF INTEGRALS, SERIES AND PRODUCTS"
(Academic Press, 1980)