

Assignment # 1

1

1.(a) The roots of the polynomial

$$Ax^2 + Bx + C = 0$$

are just given by the quadratic equation:

$$x = \frac{-B \pm \sqrt{B^2 - 4AC}}{2A}$$

(b) The roots will be real when $B^2 - 4AC > 0$.

2. Consider the complex valued function

$$z(t) = a e^{i\alpha} e^{i\omega t} + b e^{i\beta} e^{i\omega t}$$

which we would like to write in the form

$$z(t) = r e^{i(\omega t + \varphi)}$$

We can factor the $e^{i\omega t}$ so that

$$z(t) = (a e^{i\alpha} + b e^{i\beta}) e^{i\omega t}$$

in which case we just need to derive the conditions on r and φ such that

$$a e^{i\alpha} + b e^{i\beta} = r e^{i\varphi}.$$

$$\begin{aligned} \text{First, } r^2 &= |r e^{i\varphi}|^2 = |a e^{i\alpha} + b e^{i\beta}|^2 \\ &= (a e^{i\alpha} + b e^{i\beta})(a e^{-i\alpha} + b e^{-i\beta}) \\ &= a^2 + ab(e^{i(\alpha-\beta)} + e^{-i(\alpha-\beta)}) + b^2 \end{aligned}$$

Then, using the identity $\cos \theta = \frac{1}{2}(e^{i\theta} + e^{-i\theta})$ this can be written:

$$r^2 = a^2 + b^2 + 2ab \cos(\alpha - \beta).$$

Since r is assumed to be positive we can write

$$r = \sqrt{a^2 + b^2 + 2ab \cos(\alpha - \beta)}$$

The phase, φ , is determined as follows.

$$\text{Since } e^{i\varphi} = \cos \varphi + i \sin \varphi = \frac{1}{r} (a e^{i\alpha} + b e^{i\beta})$$

$$\begin{aligned} \text{we can write } \tan \varphi &= \frac{\sin \varphi}{\cos \varphi} = \frac{\text{Im } e^{i\varphi}}{\text{Re } e^{i\varphi}} \\ &= \frac{\text{Im}(a e^{i\alpha} + b e^{i\beta})}{\text{Re}(a e^{i\alpha} + b e^{i\beta})}. \end{aligned}$$

Expand $a e^{i\alpha} + b e^{i\beta}$ into real and imaginary parts:

$$\begin{aligned} a e^{i\alpha} + b e^{i\beta} &= a(\cos \alpha + i \sin \alpha) + b(\cos \beta + i \sin \beta) \\ &= a \cos \alpha + b \cos \beta + i(a \sin \alpha + b \sin \beta) \end{aligned}$$

$$\text{Therefore, } \tan \varphi = \frac{a \sin \alpha + b \sin \beta}{a \cos \alpha + b \cos \beta}.$$

It is necessary to keep track of which quadrant φ lies in by examining the signs of the numerator and denominator.

3. (a) When two springs with spring constants k_1 and k_2 are connected in parallel, a force, F will stretch both springs by the same distance, Δx .

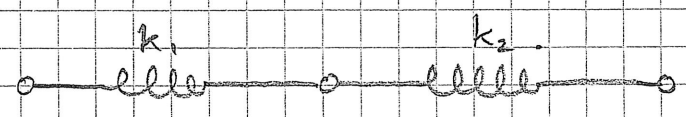
But the force F is the sum of the forces exerted by the springs, so

$$\begin{aligned}
 F &= F_1 + F_2 = -k_1 \Delta x - k_2 \Delta x \\
 &= -(k_1 + k_2) \Delta x \\
 &= -k_p \Delta x.
 \end{aligned}$$

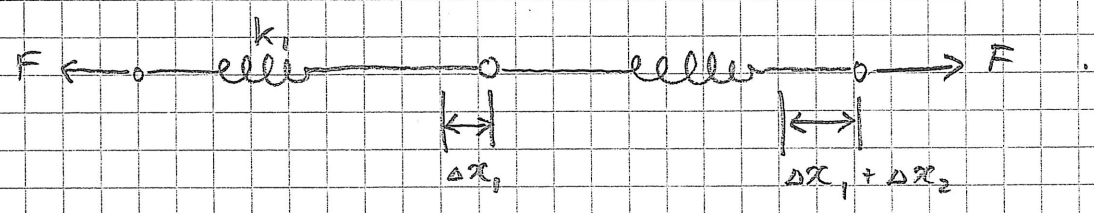
Thus, the effective spring constant for springs connected in parallel is

$$k_p = k_1 + k_2.$$

(b) Consider two springs that are connected in series as shown:



When a force F is applied to both springs, they will both stretch by amounts Δx_1 and Δx_2 :



The force applied to each spring individually is still F , so

$$F = -k_1 \Delta x_1 = -k_2 \Delta x_2 = -k_s (\Delta x_1 + \Delta x_2).$$

$$\text{Thus, } \Delta x_1 = -\frac{F}{k_1}, \quad \Delta x_2 = -\frac{F}{k_2}$$

$$\text{and } \Delta x_1 + \Delta x_2 = -F \left(\frac{1}{k_1} + \frac{1}{k_2} \right) = -\frac{F}{k_s}.$$

$$\text{Therefore, } k_s = \left(\frac{1}{k_1} + \frac{1}{k_2} \right)^{-1}.$$

4. Consider an object of length L with cross sectional area given by the function $A(l)$ where $0 \leq l \leq L$.

If the object has uniform elastic modulus, Y then a thin slice of thickness Δl , located at a position l will behave like a spring with spring constant k given by

$$k = \frac{YA(l)}{\Delta l}$$

so that it will stretch a distance x when subjected to a force F :

$$F = -\frac{YA(l)}{\Delta l}x = -kx$$

(This result was derived in class or consider Eq 3-10 in the text.)

The entire object can be treated as many springs in series.

If the object is divided into N slices then, using the result from question 3, the effective spring constant will be

$$\frac{1}{k_{equiv}} = \frac{1}{k_1} + \frac{1}{k_2} + \dots + \frac{1}{k_N}$$

In the limit as $\Delta l \rightarrow 0$ and $N \rightarrow \infty$ we can write this as an integral:

$$\frac{1}{k_{equiv}} = \int_0^L d\left(\frac{1}{k}\right)$$

$$\text{where } d\left(\frac{1}{k}\right) = \frac{dl}{YA(l)}$$

$$\text{Therefore, } \frac{1}{k_{equiv}} = \frac{1}{Y} \int_0^L \frac{dl}{A(l)}$$

Whether or not this integral can be evaluated depends on the function $A(l)$. It could be done numerically of course.

5. One end of the spring is located at position

$$X(t) = l + vt$$

while the end with the mass attached to it is located at position $x(t)$. The force acting on the mass is then

$$F = -k(x(t) - X(t) + l)$$

The mass satisfies Newton's second law:

$$F = -k(x(t) - X(t) + l) = m \frac{d^2 x}{dt^2}$$

$$\text{Thus, } m \frac{d^2 x}{dt^2} + kx(t) = k(X(t) - l) = kuvt.$$

Written in standard form this is

$$\frac{d^2 x}{dt^2} + \omega_0^2 x = \frac{k}{m} vt \quad \text{where } \omega_0 = \sqrt{\frac{k}{m}}$$

The homogeneous equation is

$$\frac{d^2 x}{dt^2} + \omega_0^2 x = 0$$

and has solutions $x_1(t) = A \cos(\omega_0 t + \varphi)$.

A particular solution to the nonhomogeneous equation is

$$x_2(t) = \frac{1}{\omega_0^2} \cdot \frac{kvt}{m} = vt$$

Therefore, the complete solution to the differential equation is

$$x(t) = x_1(t) + x_2(t) = A \cos(\omega_0 t + \varphi) + vt$$

At time $t=0$ the mass is located at $x=0$.

Therefore we must have $\varphi = \pi/2$ and we can rewrite the solution in the form

$$x(t) = A \sin(\omega_0 t) + vt.$$

Also, at time $t=0$, the velocity of the mass is

$$\dot{x}(0) = 0$$

But, $\dot{x}(t) = A\omega_0 \cos \omega_0 t + v$

and at $t=0$ we must have

$$A\omega_0 + v = 0.$$

Therefore, $A = -v/\omega_0$ and the complete

solution can be written

$$x(t) = -\frac{v}{\omega_0} \sin(\omega_0 t) + vt$$

Out of curiosity, we can check the direction of the acceleration.

$$\frac{dx}{dt} = -v \cos(\omega_0 t) + v$$

$$\text{and } \frac{d^2x}{dt^2} = v\omega_0 \sin(\omega_0 t)$$

which is positive for small $t \ll 1/\omega_0$.

Analyzing the problem this way is equivalent to recognising that $\bar{x}(t)$ could be used to represent the motion of the spring about its equilibrium position, which will be moving at constant velocity v .