# Modeling Thermal & Non-Thermal Signatures from the GC

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Workshop on Relativistic Plasma Astrophysics: Wednesday, May 11<sup>th</sup>

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### Properties of Sgr A\*

- Lies in a low-luminosity X-ray emission state:  $L_X \approx 2.4 \times 10^{33} \ erg \ s^{-1}$  (Baganoff et al. 2003)
- Frequent to short, X-ray flares with timescales ~ minutes to hours at ~ 10  $R_g$  (Degenaar et al. 2013, Barrière et al. 2014)
- Our work explores possible monthlong X-ray flaring events at ~ 10<sup>3</sup> R<sub>g</sub> (Giannios & Sironi 2013, Giannios & Lorimer 2016, Christie et al. 2016)



Composite X-ray image (Y.Bai. et al).

### Gas Density around Sgr A\*

- Sgr A\* is an ideal place to study quiescent accretion and properties of a geometrically thick disk.
- ◆ At  $R_b \sim 10^5 R_g$ , Chandra resolves X-ray, thermally emitting gas with density  $n \sim 100 \ cm^{-3}$  (Baganoff et al. 2003)
- ♦ At ~ 10 R<sub>g</sub>, through modeling and Faraday rotation measurements, constrain  $n \sim 10^7 cm^{-3}$  (Marrone et al. 2007, Moscibrodzka et al. 2009)
- Our work aims on bridging this gap by exploring probes of the disk density.

### Stars in GC: S-Cluster & S2

- \* Massive, B-type stars  $\longrightarrow$  winds of  $\nu \sim 10^8 cm/s$
- ♦ The brightest star, S2, is characterized by a close pericenter ~10<sup>3</sup> R<sub>g</sub> and mass loss rate  $\dot{M}_w \leq 10^{-7} M_{\odot} yr^{-1}$  (Martins et al. 2008)
- These stars, specifically S2, are ideal probes of the accretion disk of Sgr A\* (Giannios & Sironi 2013)



# Stellar Wind – Accretion Disk Interactions

 These interactions, during the pericenter passage of a star, lead to the formation of a bow shock in the stellar winds

3.0 2.5  $v_d$ 1.5  $\mathbf{R}(\boldsymbol{\theta})$ Shocked Wind Material 1.0 0.5 Shocked Disk  $v_{w}$ 00 Material -2 -3 2 -1  $R_0$ Z

Sketch of Bow Shock Model: Rest Frame of Star

# Determining Shape of Termination Shock

- Assumptions for constructing semianalytical model:
  - i. System has reached a steady state
  - ii. Shocked wind region falls within thin-shell limit
- Follow a modified analysis of momentum supported bow shocks (Wilkin,1996)
- \* Include thermal pressure of disk:  $P_{therm} = \alpha P_{ram}$

#### Shape of Termination Shock



### Properties of Shocked Stellar Wind

• Using shock jump conditions  $\longrightarrow$  derivations of  $T_{sw}(\theta), \rho_{sw}(\theta), \text{ and } H(\theta) \longrightarrow$  thermal bremsstrahlung power



# Testing Model Through Hydro-Simulations

#### Used to make comparisons of:

- i. Our estimates of the termination shock and contact discontinuity
- ii. Thermal bremsstrahlung power produced from the shell
- Large back region beyond termination shock is dominated by Kelvin-Helmholtz instabilities

\* 
$$\dot{M}_w = 10^{-7} M_{\odot} yr^{-1}$$
,  $n_d = 10^4 cm^{-3}$ ,  
 $v_d = 8000 \ km \ s^{-1}$ 



 $\alpha = 0.1$ 

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# Angular Profiles of Shock Surface

- The contact discontinuity becomes quickly prone to instabilities
- \* Estimates for the termination shock in the wind and contact discontinuity are valid to  $\theta \approx \pi/2$





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# Thermal Bremsstrahlung Emission from Shell

- Computation of the thermal, radiated power is permitted up to  $\frac{\pi}{2}$  using semianalytical results
- Stellar wind composition of hydrogen with and without solar abundancies
- Appropriate scalings are derived for shell luminosity:
  - i. Hydrogen:  $L_{shell} \propto \dot{M}_w^{3/2} n_d^{1/2} (1 + 1.6 \, \alpha^{1.4})$
  - ii. SM:  $L_{shell} \propto \dot{M}_w^{3/2} n_d^{1/2} (1 + 1.2 \alpha^{1.2})$



# Role of Kelvin-Helmholtz Instabilities on Thermal Luminosity

 Back region composed of mixed fluids contributes large fraction of total luminosity from simulation (~ 10 × L<sub>shell</sub>)

- Simulations with different M
  <sub>w</sub> allow for the derived scaling factor of the total simulation luminosity:
  - i. SM:  $L_{tot} \propto \dot{M}_w^{1.3} n_d^{0.52} (1+5.5 \alpha)$

Cumulative Plot of Total Luminosity from Simulations



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### **Time-Dependent Simulation**

- Time-dependent density used to mimic transit of S2 through pericenter and study the evolution of the back, mixing region.
- Month-long flare with luminosities above quiescent emission is expected.



### Limitations to S2 Application

No specific accretion disk model was adopted for time-dependent simulation.

 Tidal effects and acceleration of S2 during pericenter passage was ignored. Such effects are expected to compress and elongate the back, mixing region — increase in thermal emission from shocked wind (Ballone et al. 2013, Gillessen et al. 2014, Ballone et al. 2016)

# Pulsar Population in GC

- The discovery of a magnetar, spin period ~ 3.76 s, within the inner ~ 0.1 pc of the GC was achieved by Swift and NuSTAR. (Degenaar et al. 2013, Mori et al. 2013, Rea et al. 2013)
- An upper limit of ~ 200 detectable pulsars was estimated to orbit around Sgr A\*. (Chennamangalam & Lorimer 2014)
- These stars are ideal probes of properties of the SMBH and its accretion disk. (Cordes & Lazio 1997)

#### Sample of 1982 pulsars from ANTF



# Comparison of Applications

S2 & Accretion Disk :

- Non-relativistic stellar winds
- Thermal X-ray emission (bremsstrahlung cooling)
- Flaring on timescale of ~month
- Quickly prone to hydrodynamic instabilities (KHI)

GC Pulsars & Accretion Disk:

- Relativistic stellar winds
- Non-thermal emission: X-ray & Radio (synchrotron)
- Flaring on timescale of several months
- KHI grows on much larger timescales

### Summary

- The formation of a bow shock structure in stellar winds through interactions is accurately described through a modified prescription of momentum-supported bow shocks.
- Application to S2's pericenter passage month-long thermal X-ray flare with luminosities comparable to quiescent emission of Sgr A\*.
- Application to pulsar-disk interactions non-thermal, X-ray flare on timescales of months with luminosities comparable to quiescent emission.

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