

Modeling of Pulsar Magnetospheres and other EMF*

(* ElectroMagnetic Fun)

Anatoly Spitkovsky (Princeton)

(with J. Arons, X. Bai, J. Li, A. Philippov, A. Tchekhovskoy)

Outline

- * Pulsar magnetosphere: open questions
- * Pulsar models: pros, cons and fails.
- * High energy emission models: emerging role of reconnection
- * Conclusions and outlook

Pulsars: observationally driven



Pulsar theory:



Open questions:

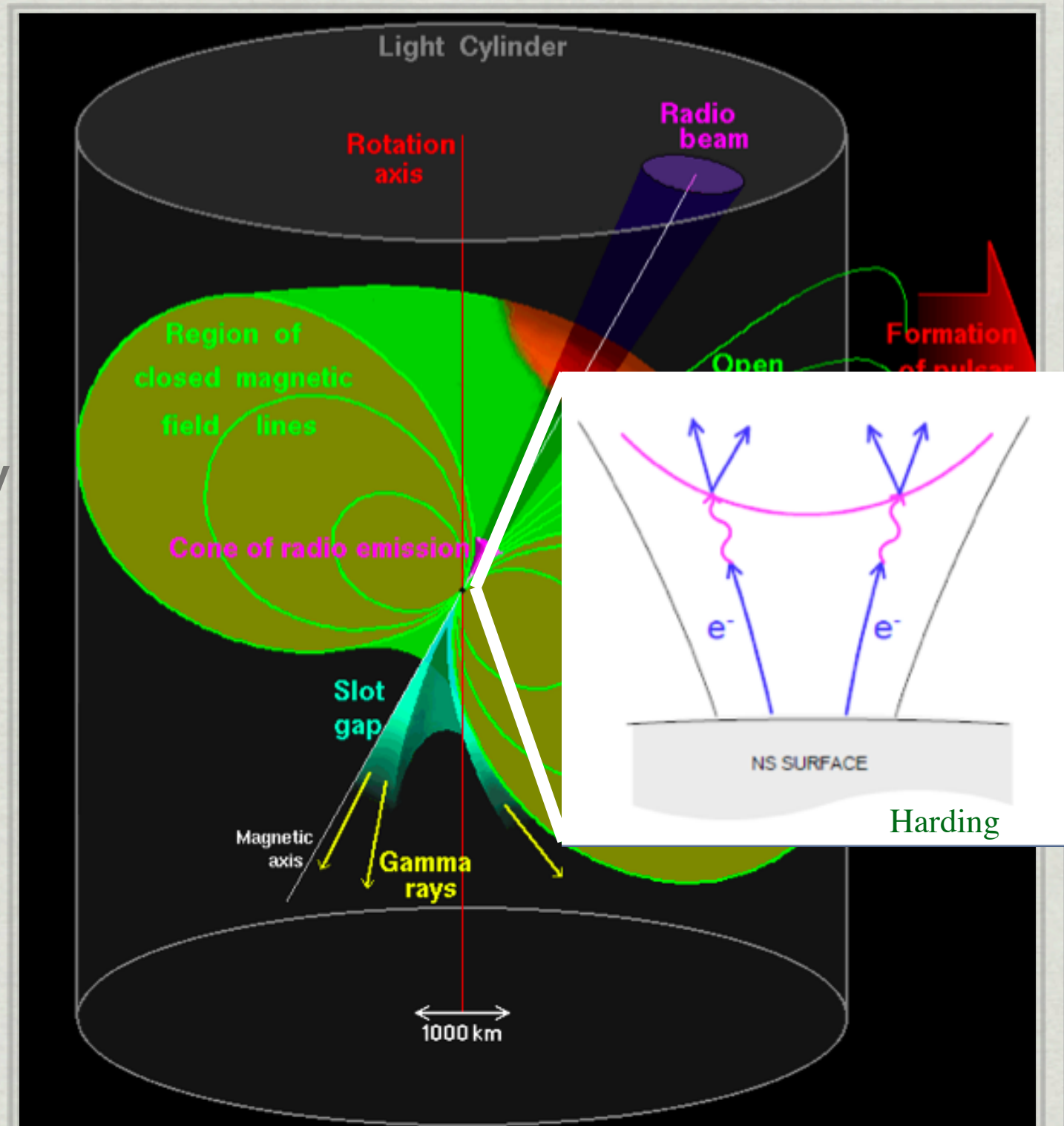
- ✱ What is the structure of pulsar magnetosphere and how do pulsars spin down?
- ✱ What are the properties of the wind near pulsar? In the nebula?
- ✱ What causes pulsed emission?
- ✱ How are observed spectra generated? (how particles are accelerated?)

Open questions:

- ✱ The structure of the magnetosphere is the primary question, as all emission physics must be done in the context of proper magnetospheric geometry.
- ✱ Related question is the nature of the spin-down: most energetic, but mostly invisible, process in normal pulsars.

Magnetospheric cartoon

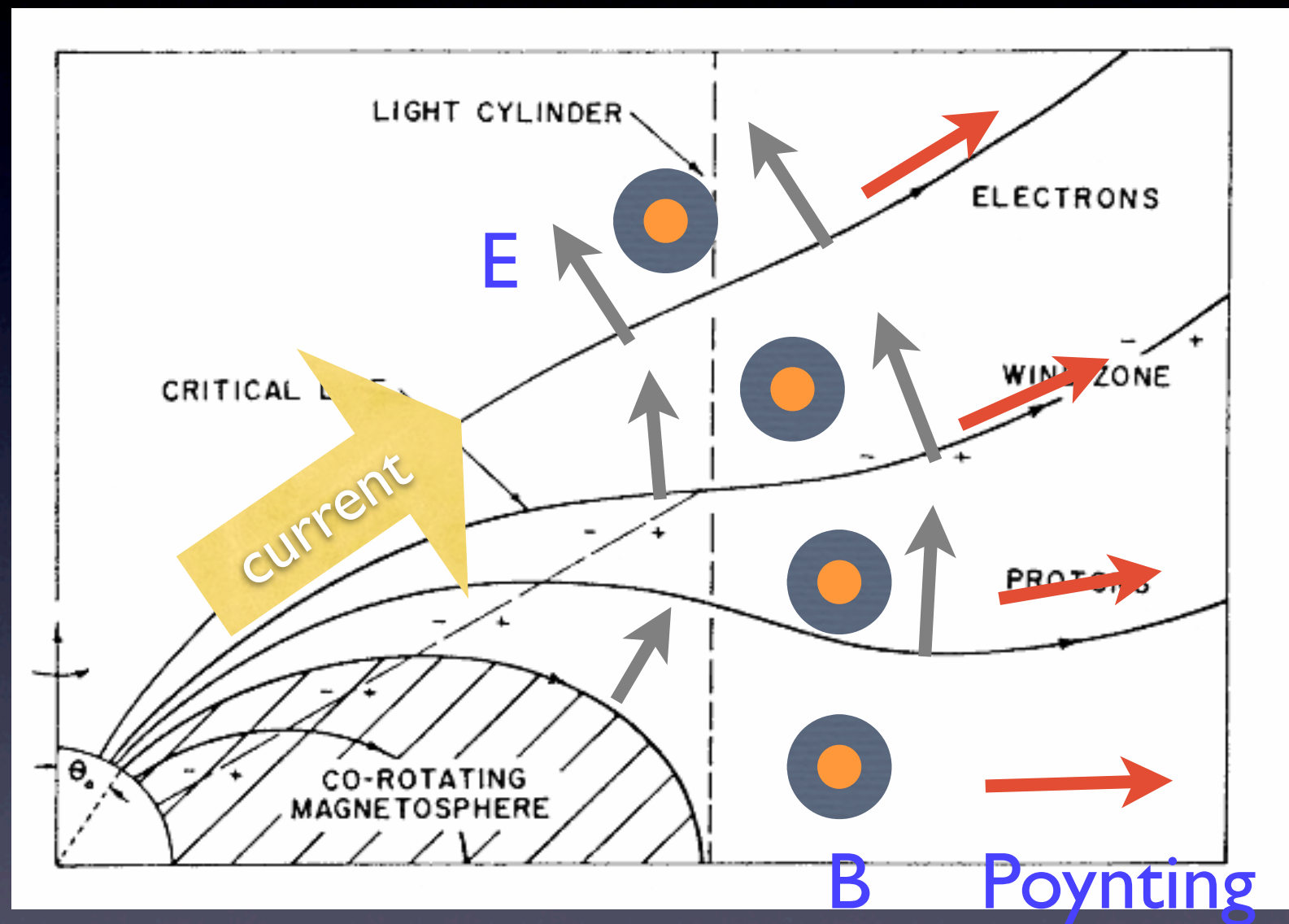
- * Open & closed (corotating) zones.
- * Minimal (Goldreich-Julian) charge density
- * Light cylinder
- * Sweepback
- * Plasma is born in discharges



Pulsars: energy loss

$$\rho_{GJ} = -\frac{\vec{\Omega} \cdot \vec{B}}{2\pi c}$$

- Corotation electric field
- Sweepback of B field due to poloidal current
- $\mathbf{E} \times \mathbf{B} \rightarrow$ Poynting flux
- Electromagnetic energy loss



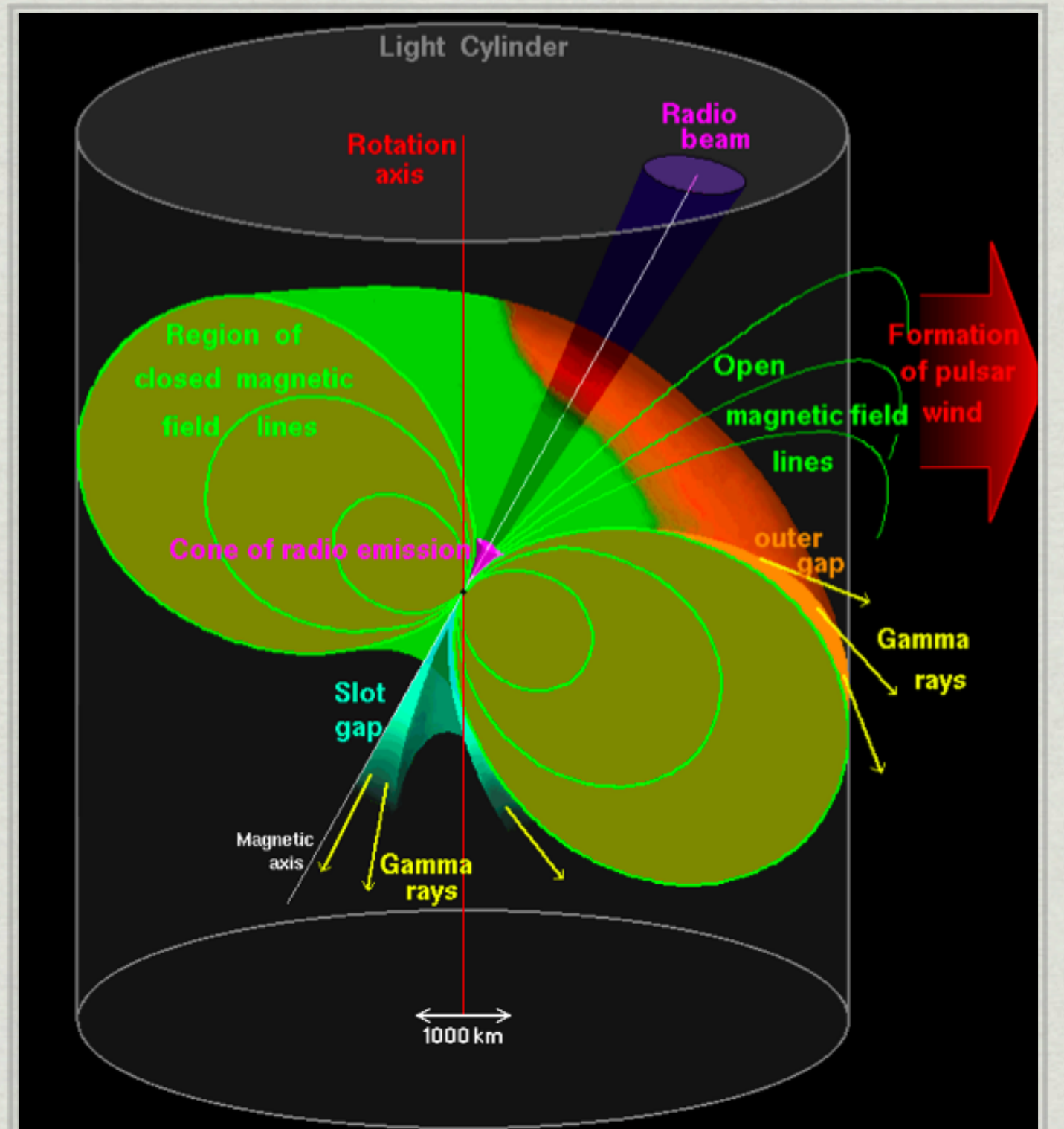
Goldreich & Julian 1969

Only gamma-ray output of pulsars is an interesting (<10%) fraction of the main spin-down energy flux.

Magnetospheric shape

Several ways of modeling, depending on charge supply:

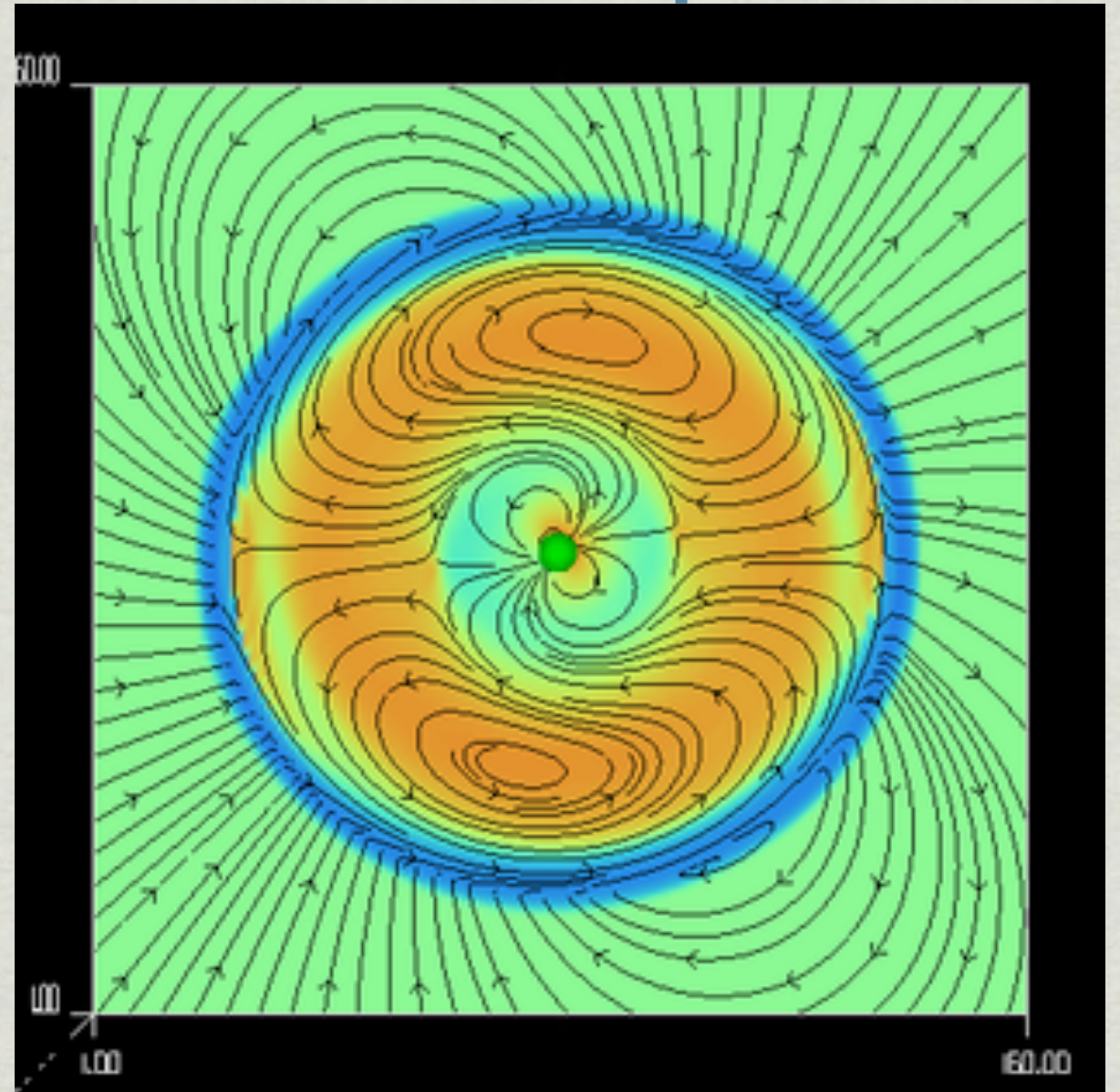
- * Vacuum rotator
- * Ab-initio particle
- * Full RMHD
- * Force-free variants
- * “Pulsar equation”



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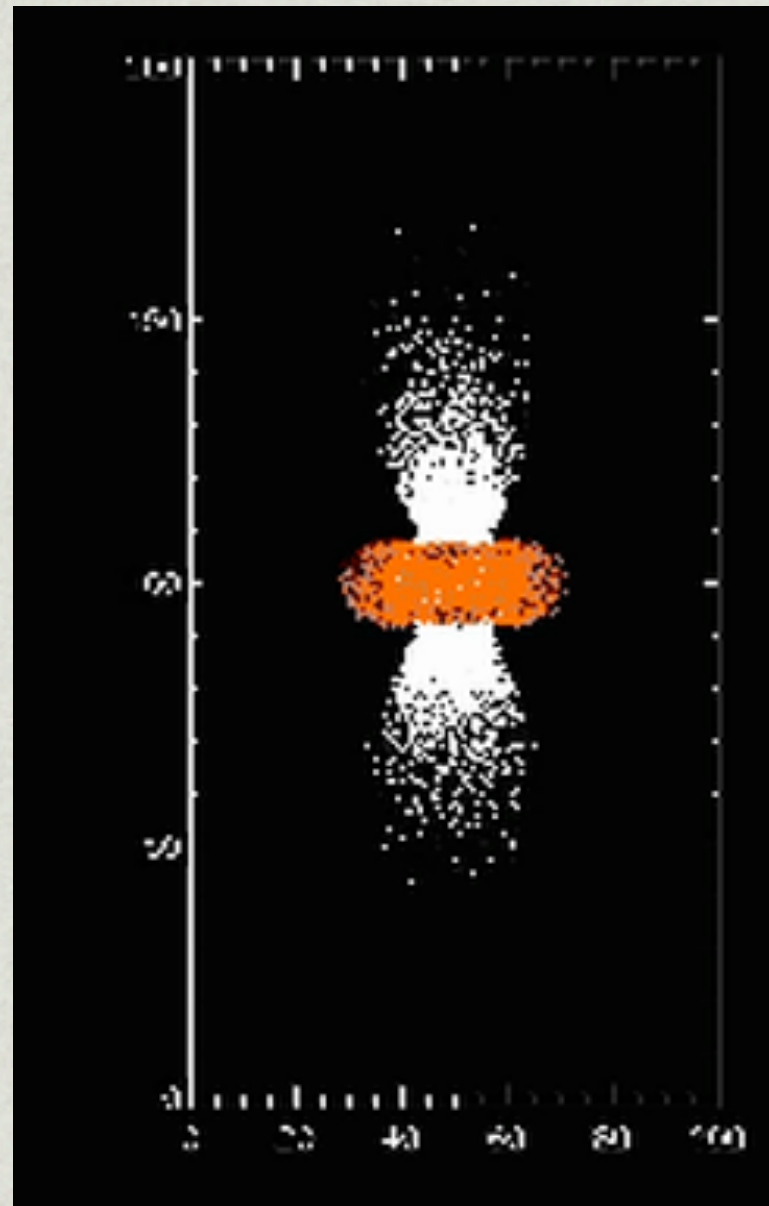


Analytic field line shape and spin down power formula

Magnetospheric shape

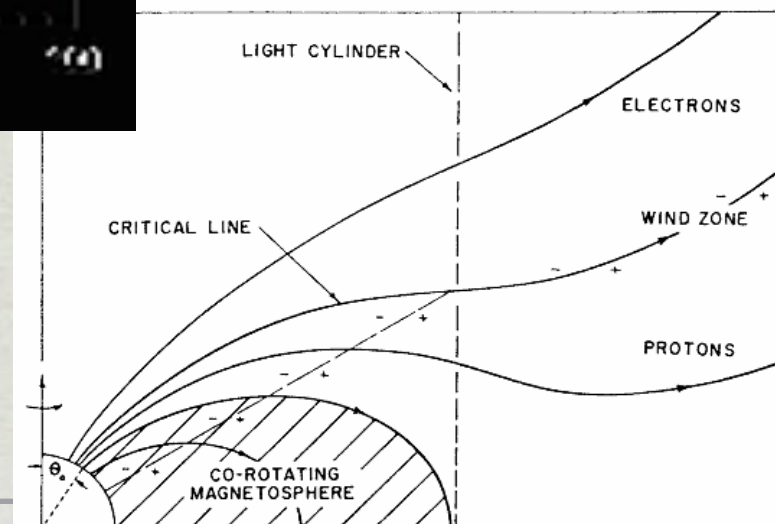
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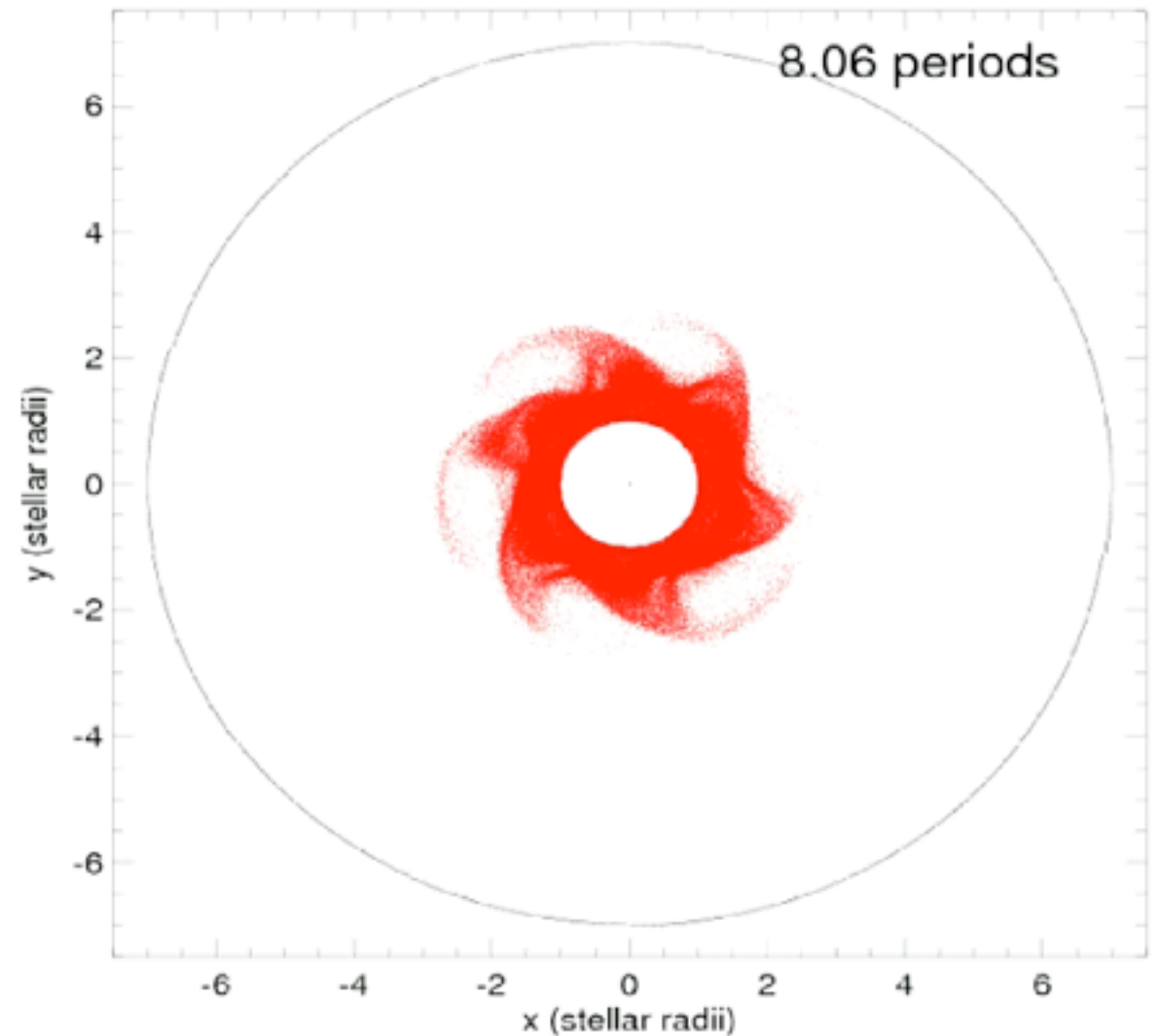
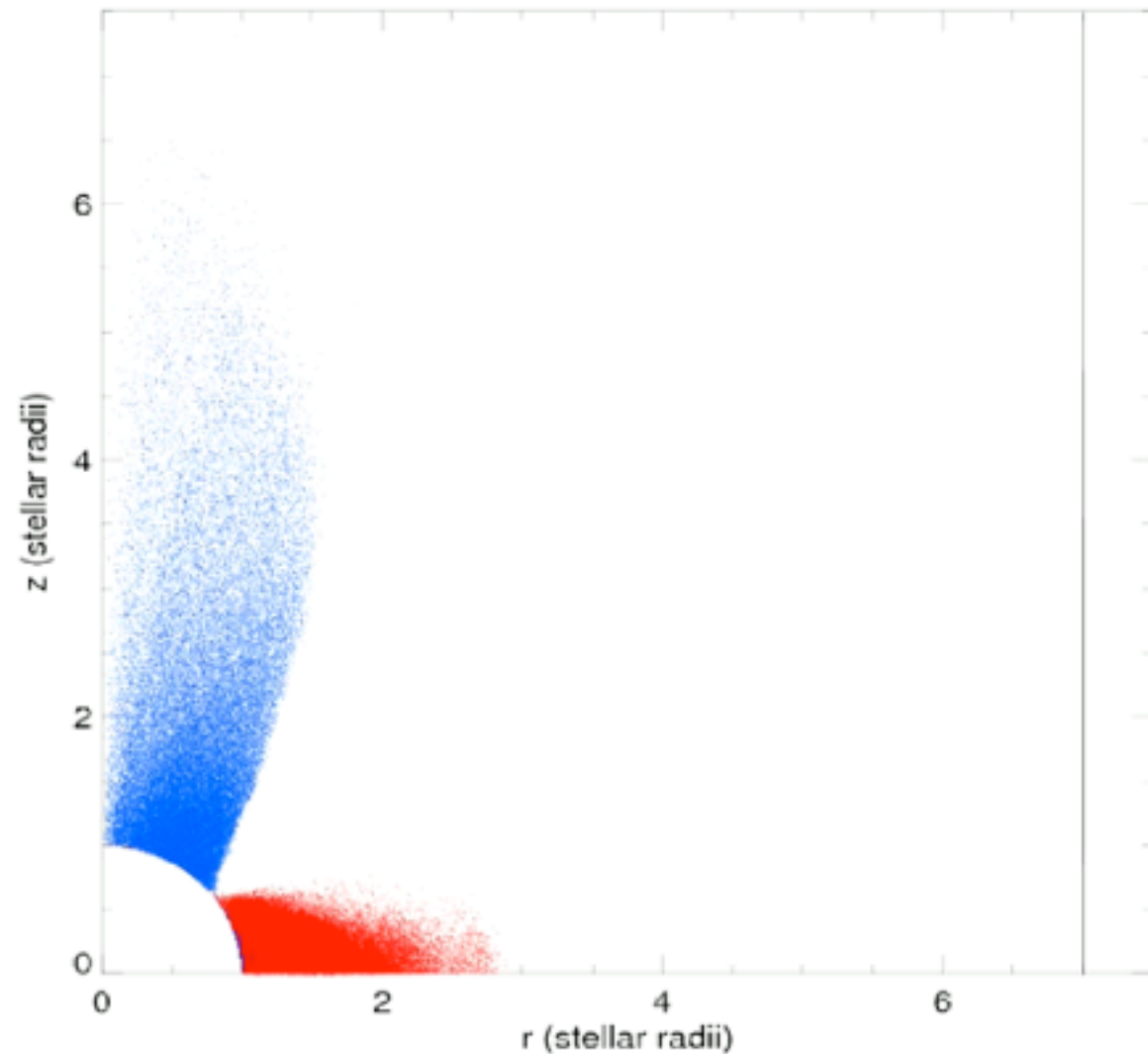


AS & Arons 02;
Michel et al 84, 01

**Is the
charge-
separated
solution
dead?**



Non-axisymmetric instabilities



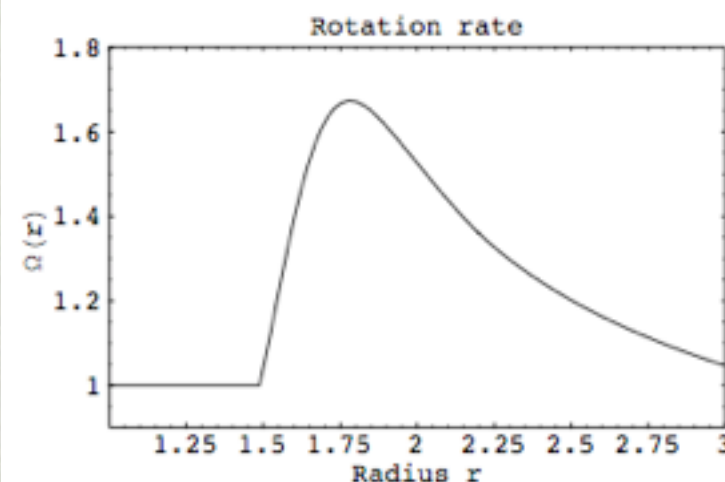
Disk-Torus Electrosphere

Michel et al '84-01

Diocotron instability

AS & Arons 02;

Petri et al 02-



Belyaev & AS, in prep

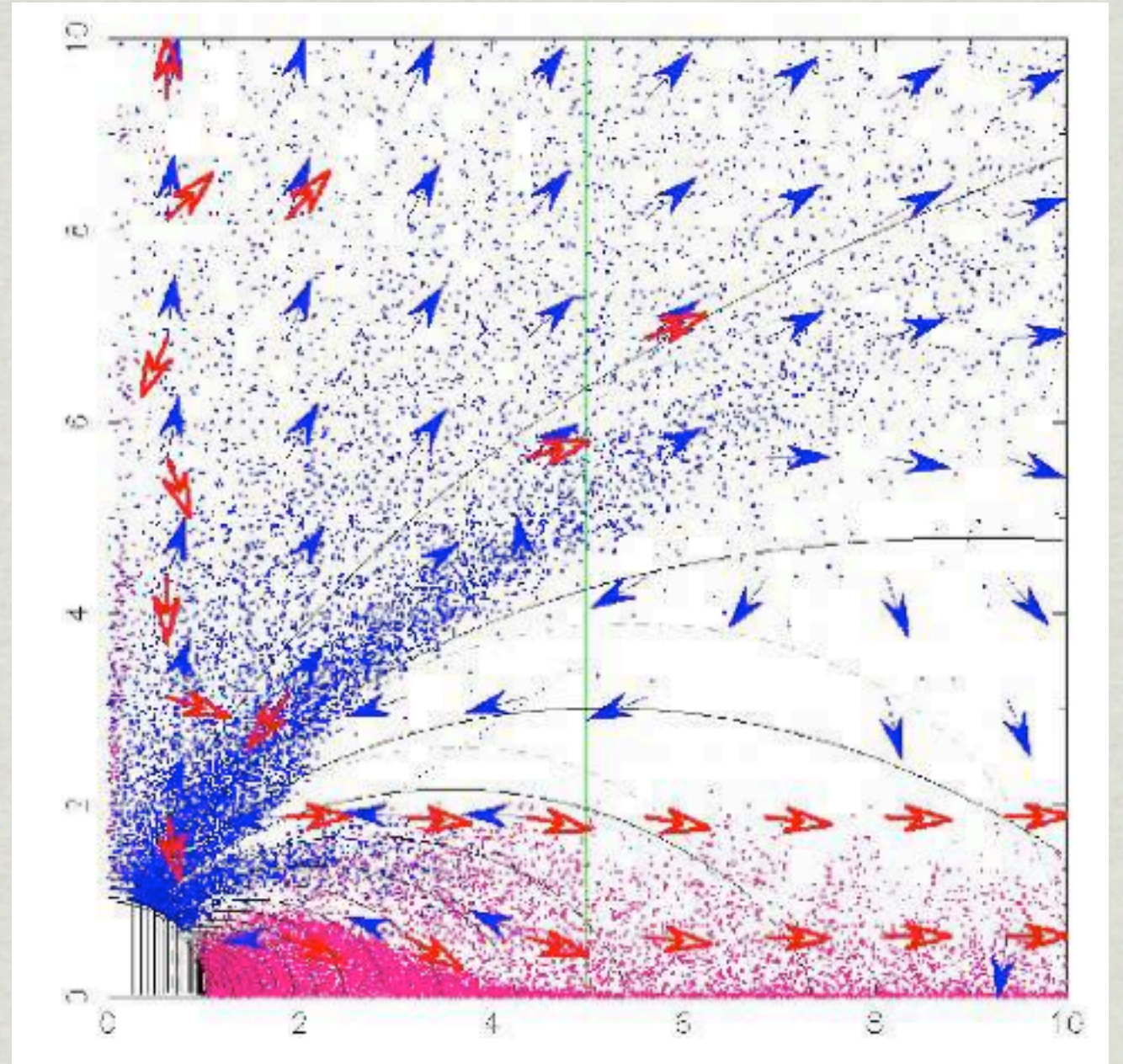
Possibility of radial current
Electrospheres are a curiosity
Add pairs?

Petri et al 02

Magnetospheric shape

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- ✱ Full RMHD
- ✱ Force-free variants
- ✱ “Pulsar equation”



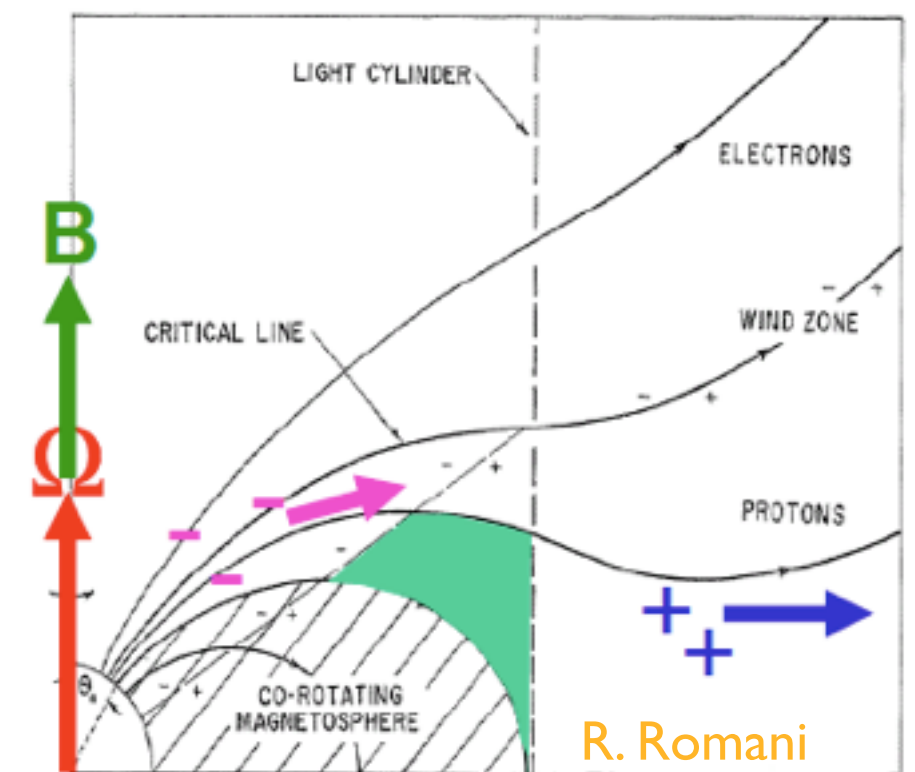
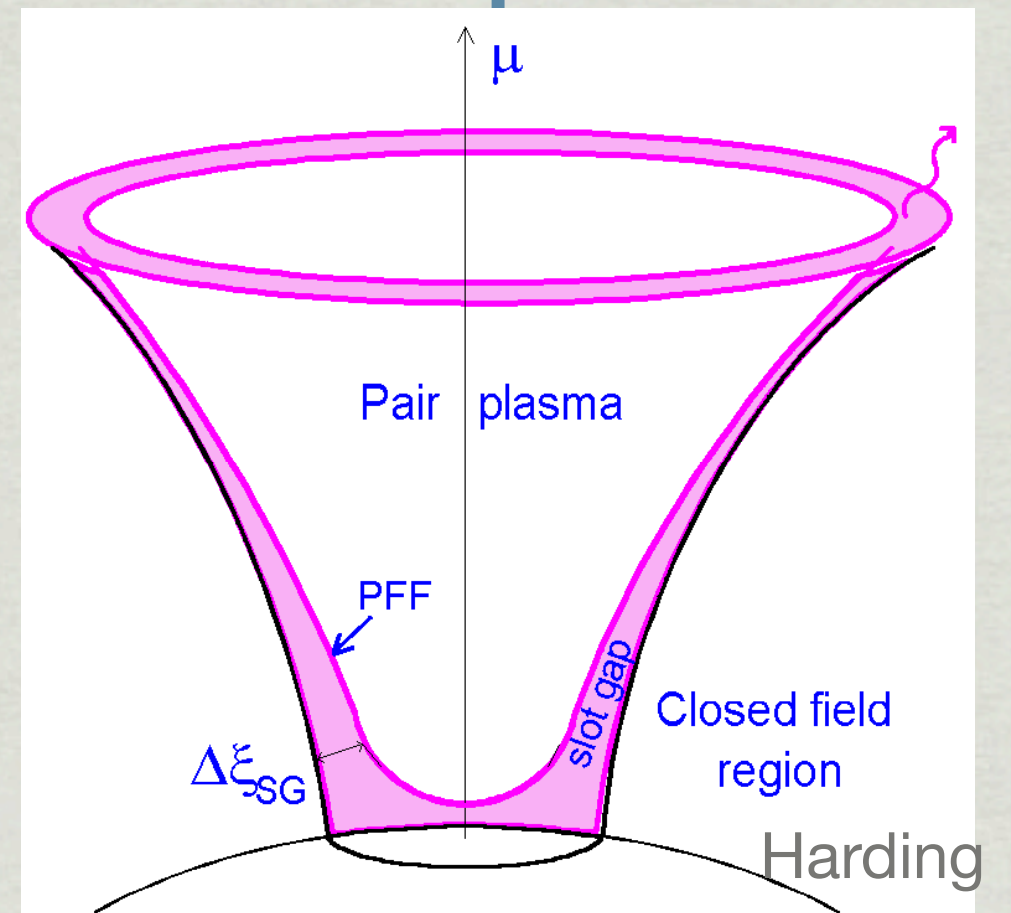
Not spinning down?

Wada & Shibata 11; Yuki & Shibata 12

Magnetospheric shape

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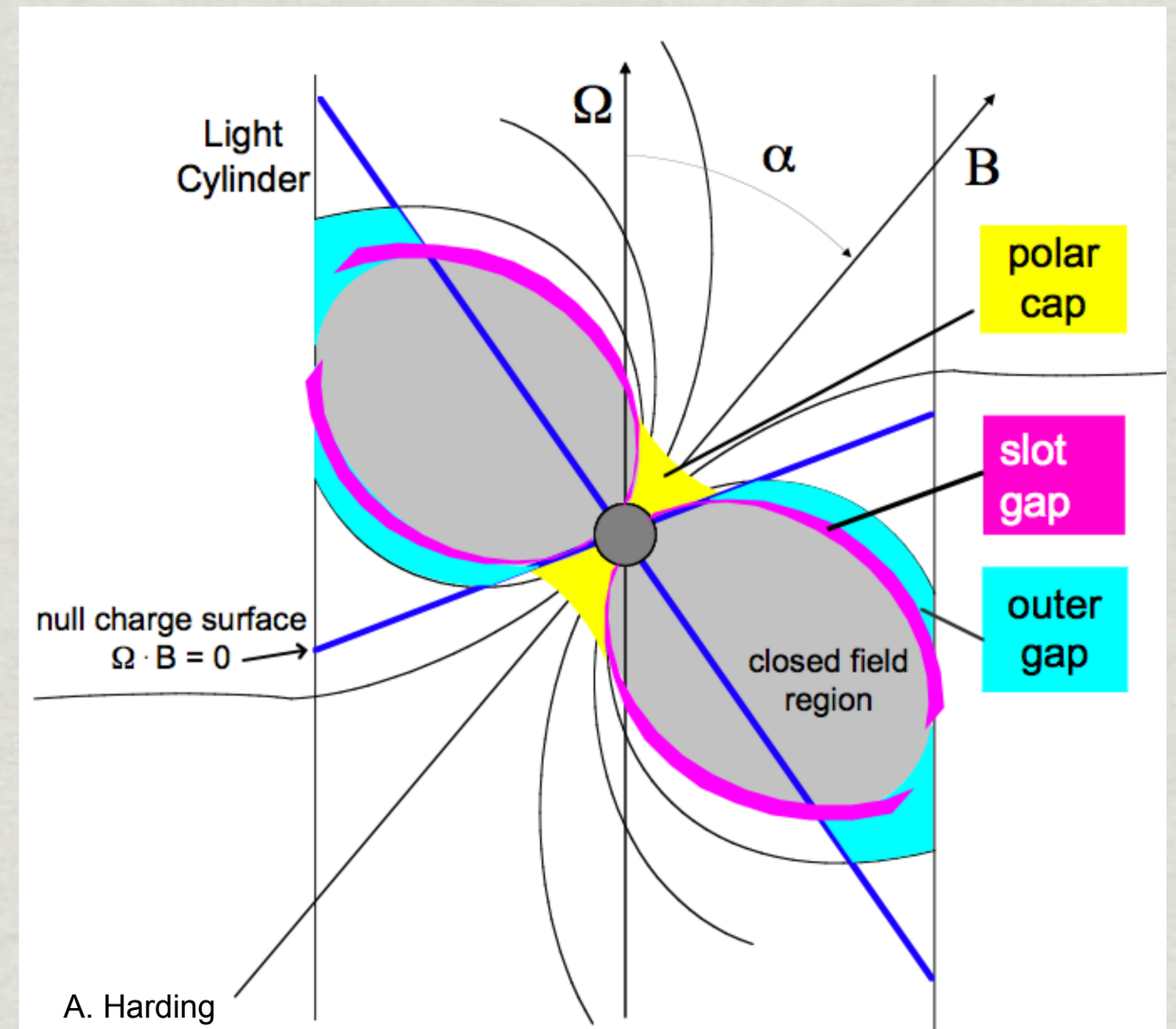
- * Vacuum rotator
- * **Particle/Vacuum+**
- * Full RMHD
- * Force-free variants
- * “Pulsar equation”



Magnetospheric shape

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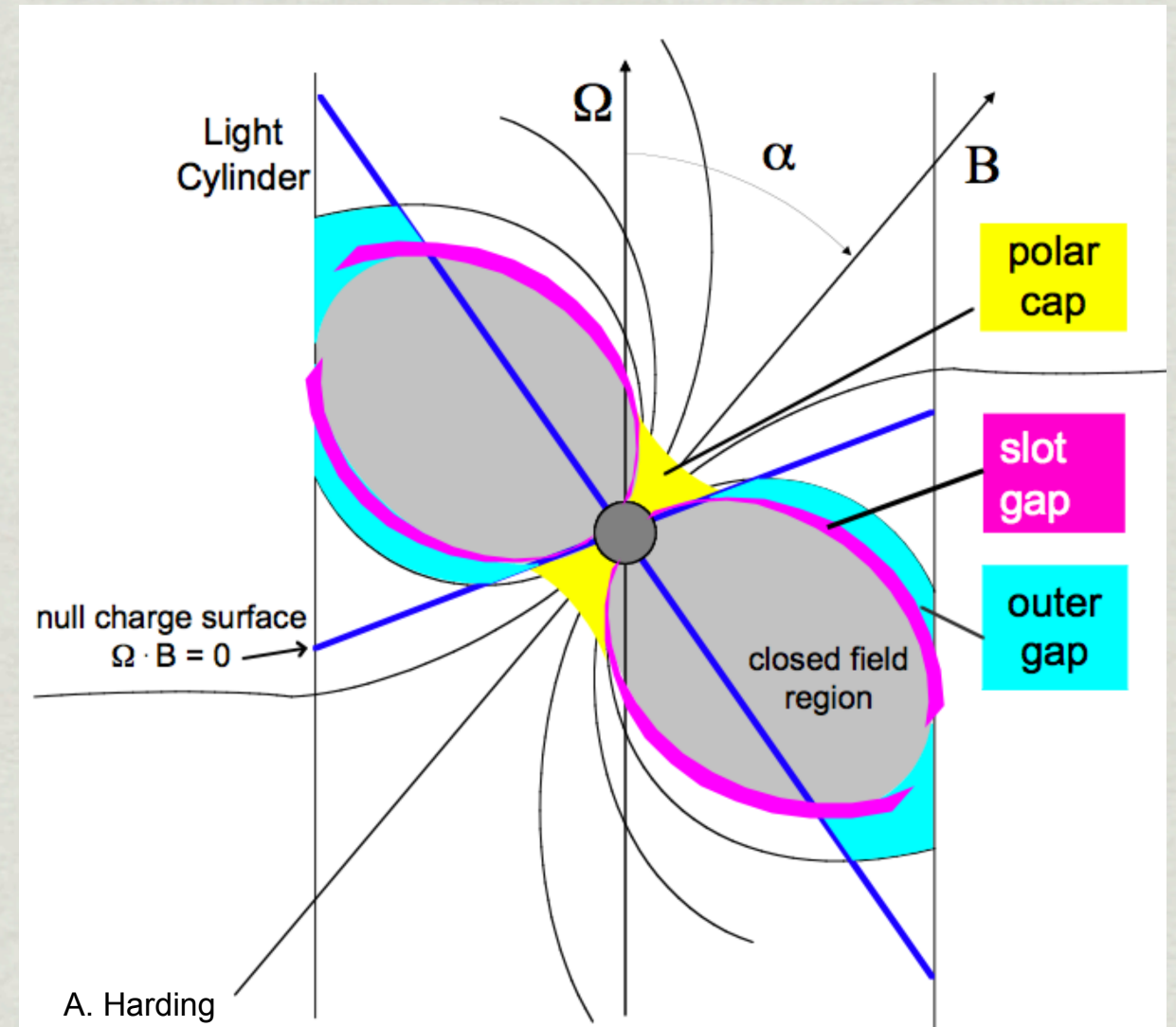


Harding

Space-charge limited flow + pairs + gaps. Curvature radiation from radiation reaction limited acceleration.

Gaps in understanding of gaps...

- ✱ Acceleration due to charge starvation
- ✱ Gaps imply space-charge separated background flow, even though pairs are created.
- ✱ PWNe require pair densities $\gg$ minimum charge separated density.



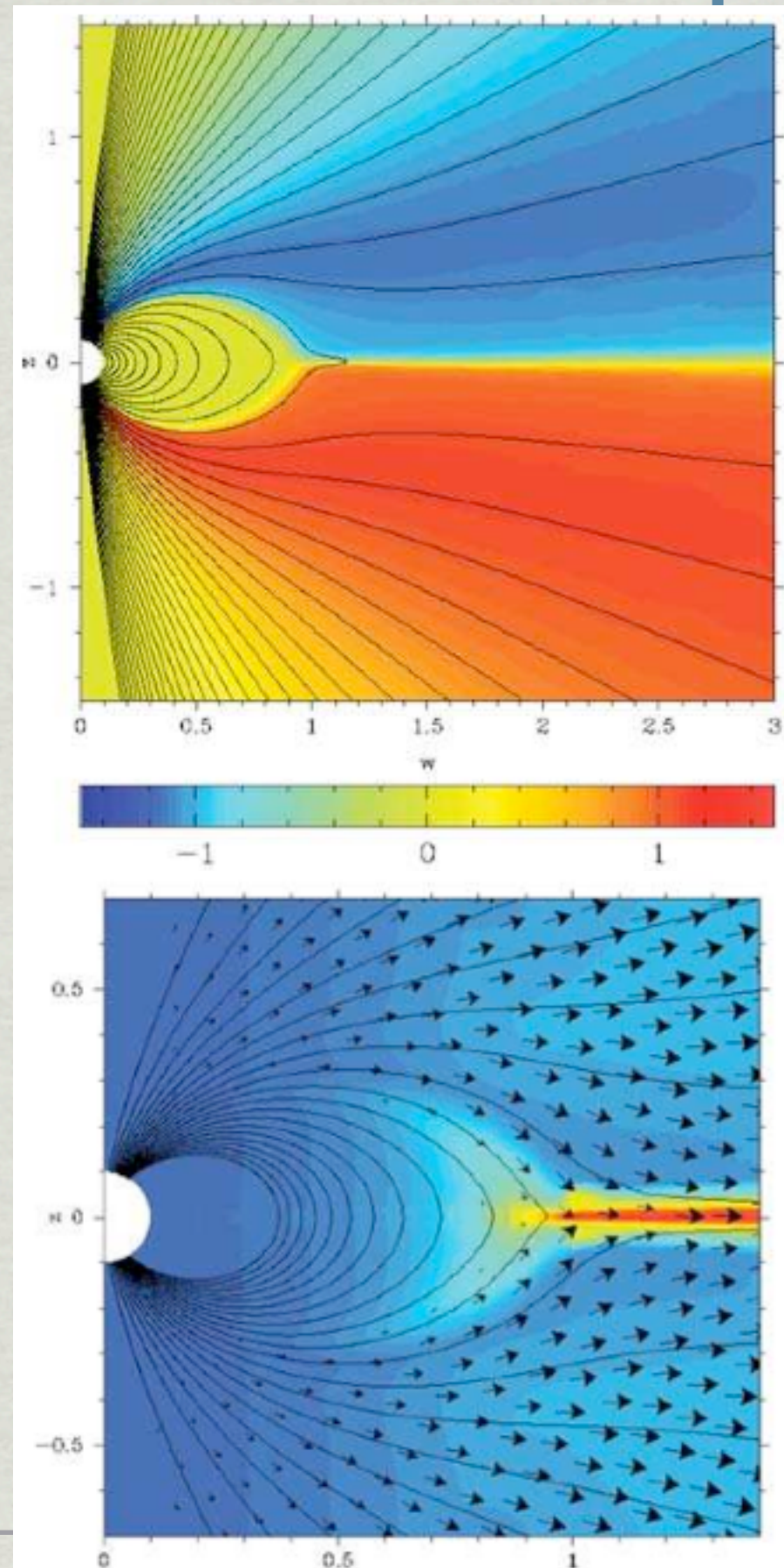
Harding

Gap models are best developed but are not self-consistent.

Magnetospheric shape

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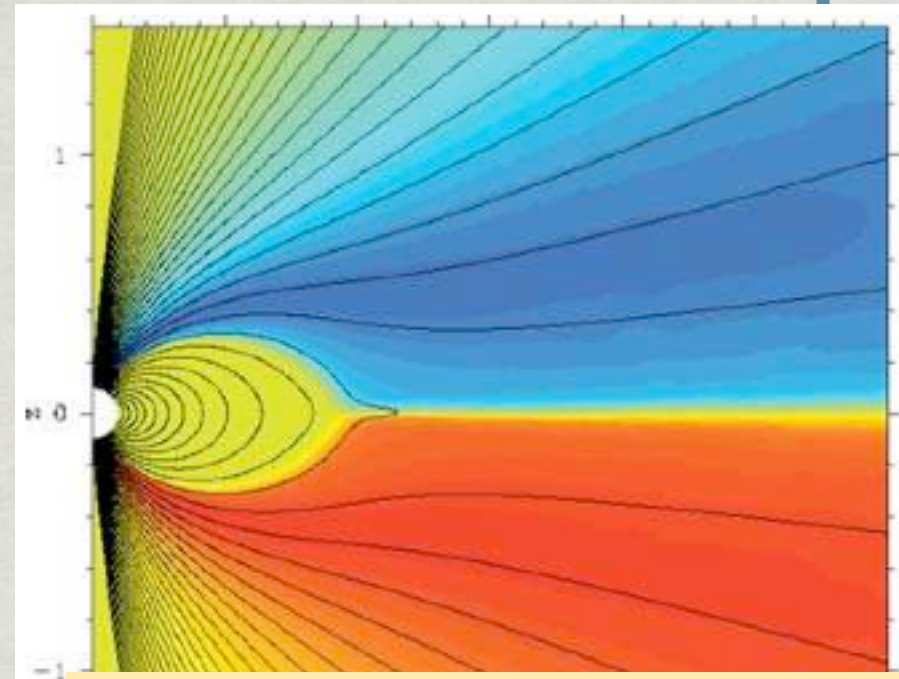


Komissarov 06

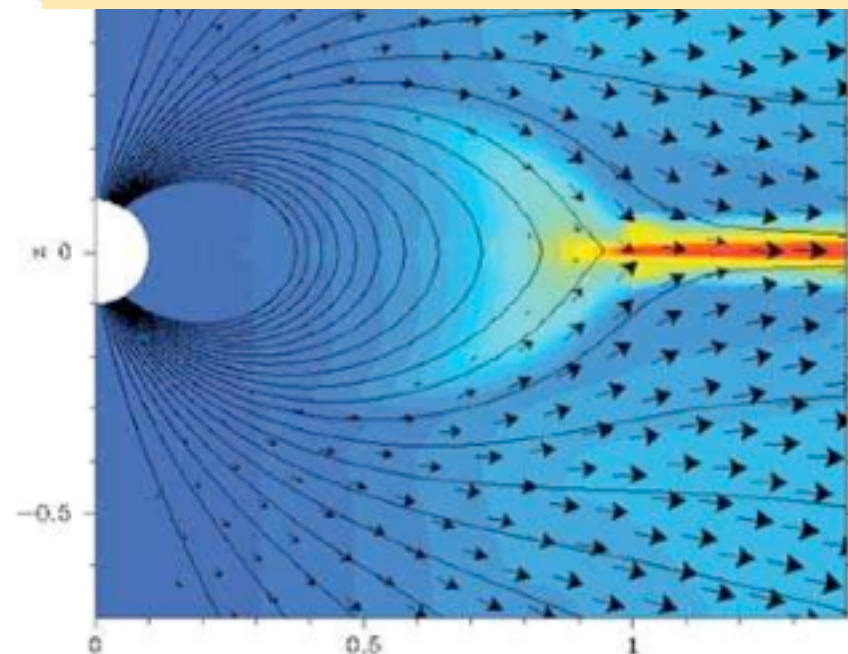
Magnetospheric shape

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High magnetization near the star is difficult to simulate with MHD. Require simplification: force-free



Komissarov 06

Magnetospheric shape

NS is immersed in massless conducting fluid with no inertia.

Several ways of modeling, depending on charge supply:

- * Vacuum rotator
- * Ab-initio particle
- * Full RMHD
- * **Force-free:**
- * “Pulsar equation”

$$\rho \mathbf{E} + (1/c) \mathbf{j} \times \mathbf{B} = 0$$

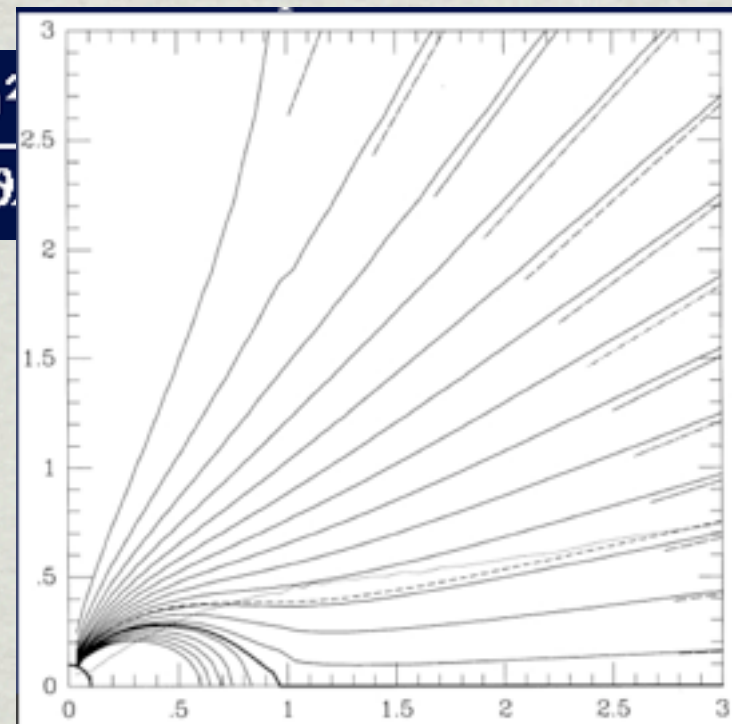
$$\frac{1}{c} \frac{\partial \mathbf{E}}{\partial t} = \nabla \times \mathbf{B} - \frac{4\pi}{c} \mathbf{j}, \quad \frac{1}{c} \frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E},$$

$$\mathbf{j} = \frac{c}{4\pi} \nabla \cdot \mathbf{E} \frac{\mathbf{E} \times \mathbf{B}}{B^2} + \frac{c}{4\pi} \frac{(\mathbf{B} \cdot \nabla \times \mathbf{B} - \mathbf{E} \cdot \nabla \times \mathbf{E}) \mathbf{B}}{B^2}$$

Gruzinov 99, Blandford 02

Time-independent version -- pulsar equation (Scharleman & Wagoner 73, Michel 73)

$$\frac{\partial^2 \Psi}{\partial x^2} + \frac{\partial^2 \Psi}{\partial y^2} =$$

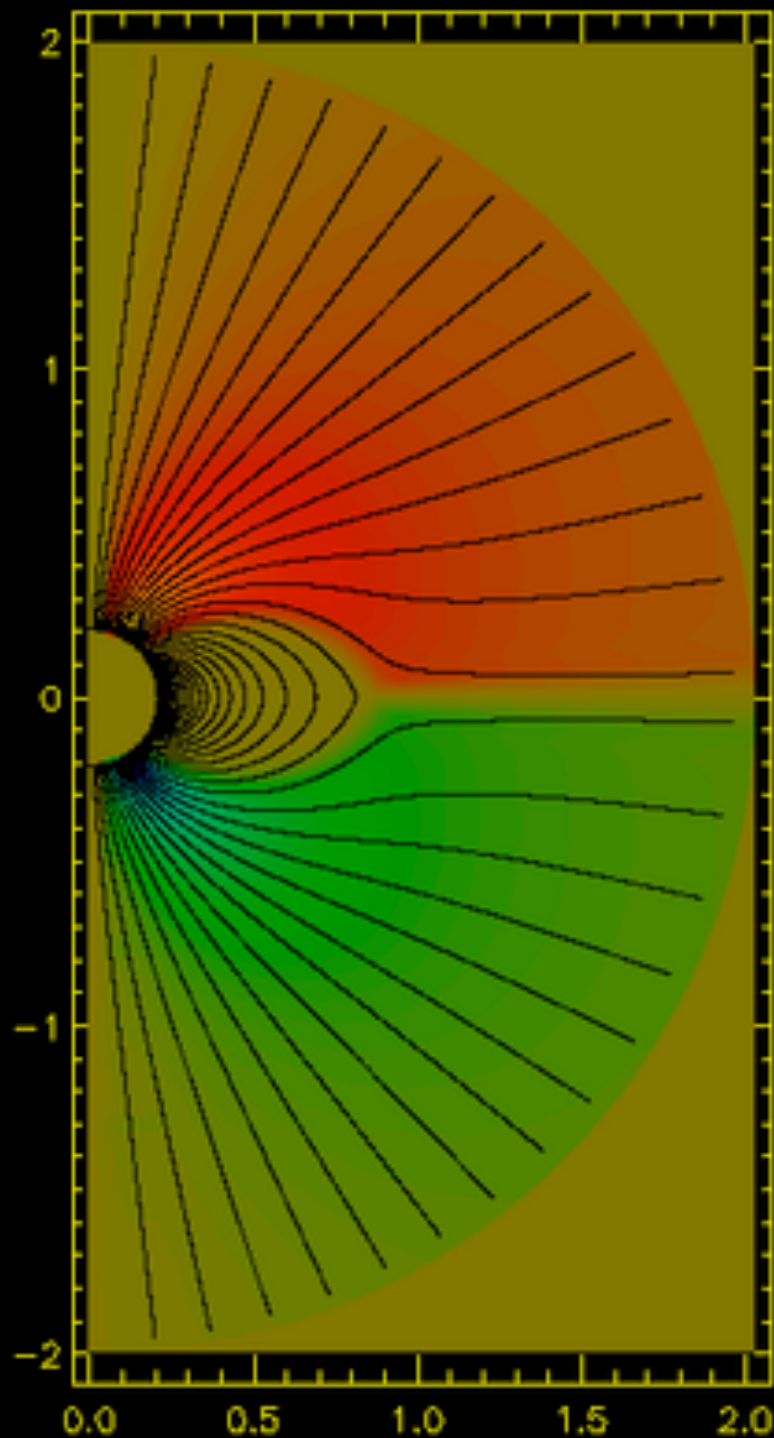
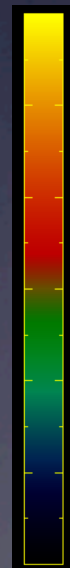


$$= \frac{\Psi I'(\Psi)}{x^2 (1-x^2)}$$

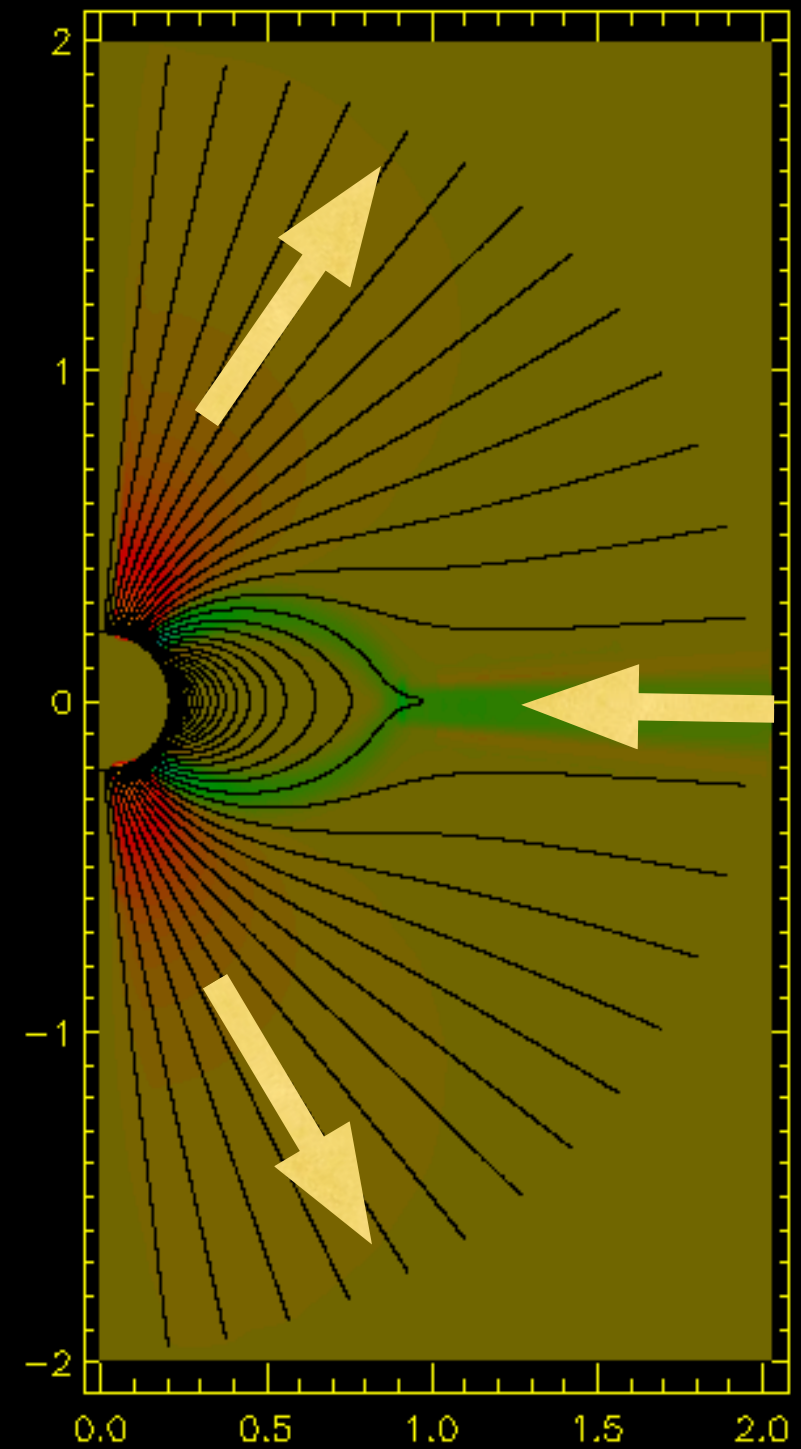
Contopoulos, Kazanas & Fendt 1999

Aligned rotator: plasma magnetosphere

Toroidal field

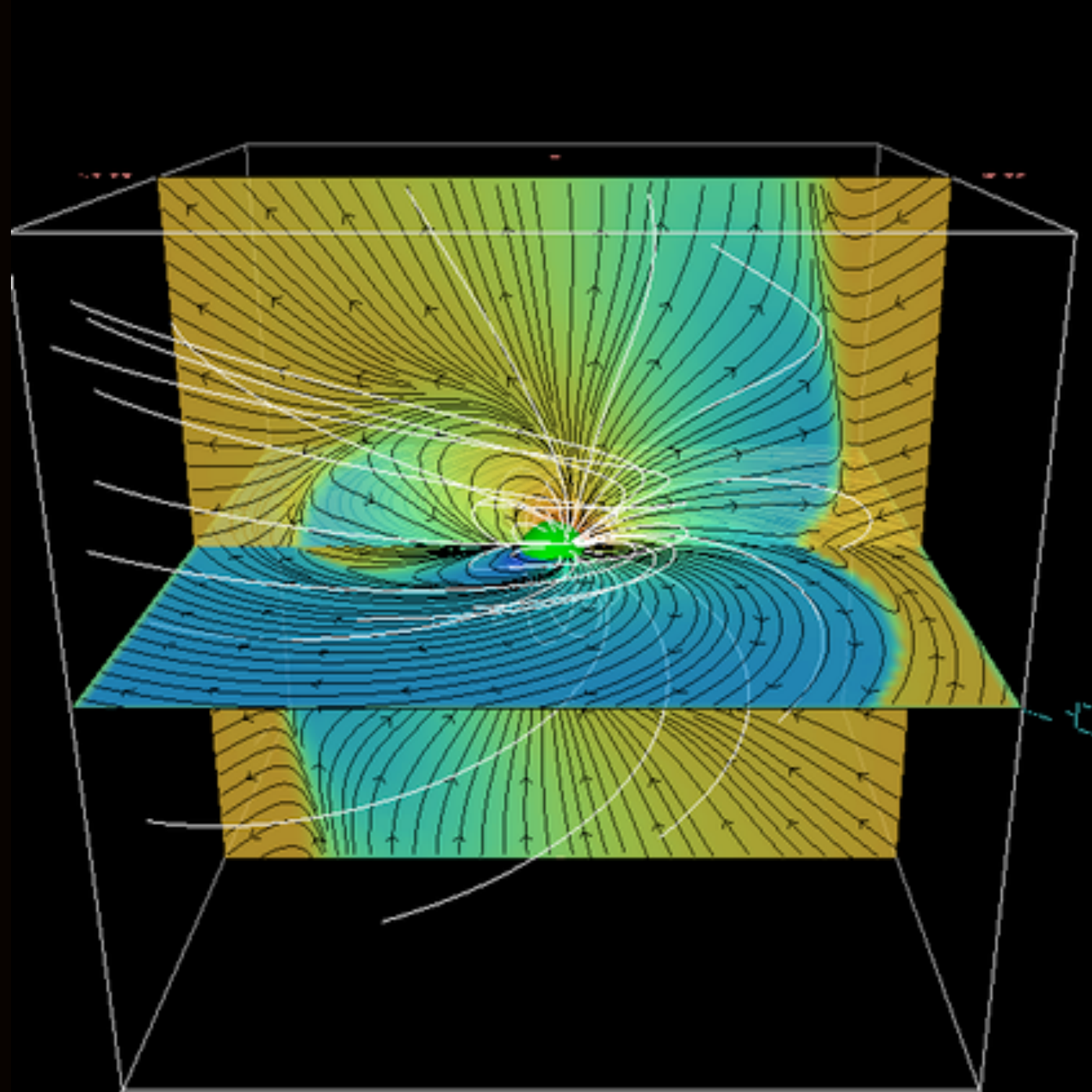


Current



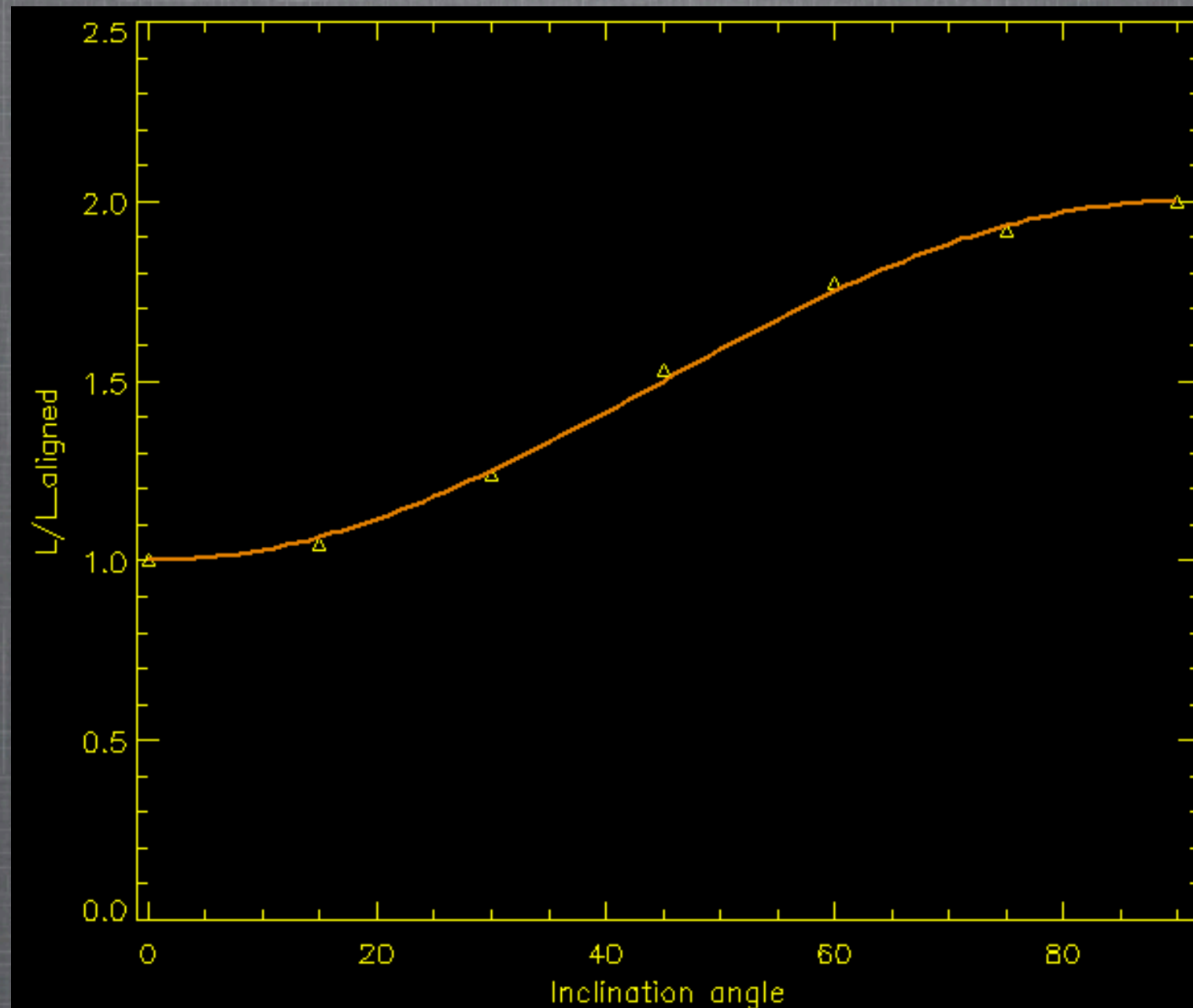
Properties: current sheet, split-monopolar asymptotics; closed-open lines; Y-point; null charge surface is not very interesting. Now at least 5 groups can do this (recently, Yu 11, Parfrey 11, Petri 12, Palenzuela 12 in addition to AS 06, McKinney 06, Kalapotharakis 09)

Oblique rotator: force-free



A.S. 2006

SPIN-DOWN POWER



Spin-down of oblique rotator

NB: this is a fit!

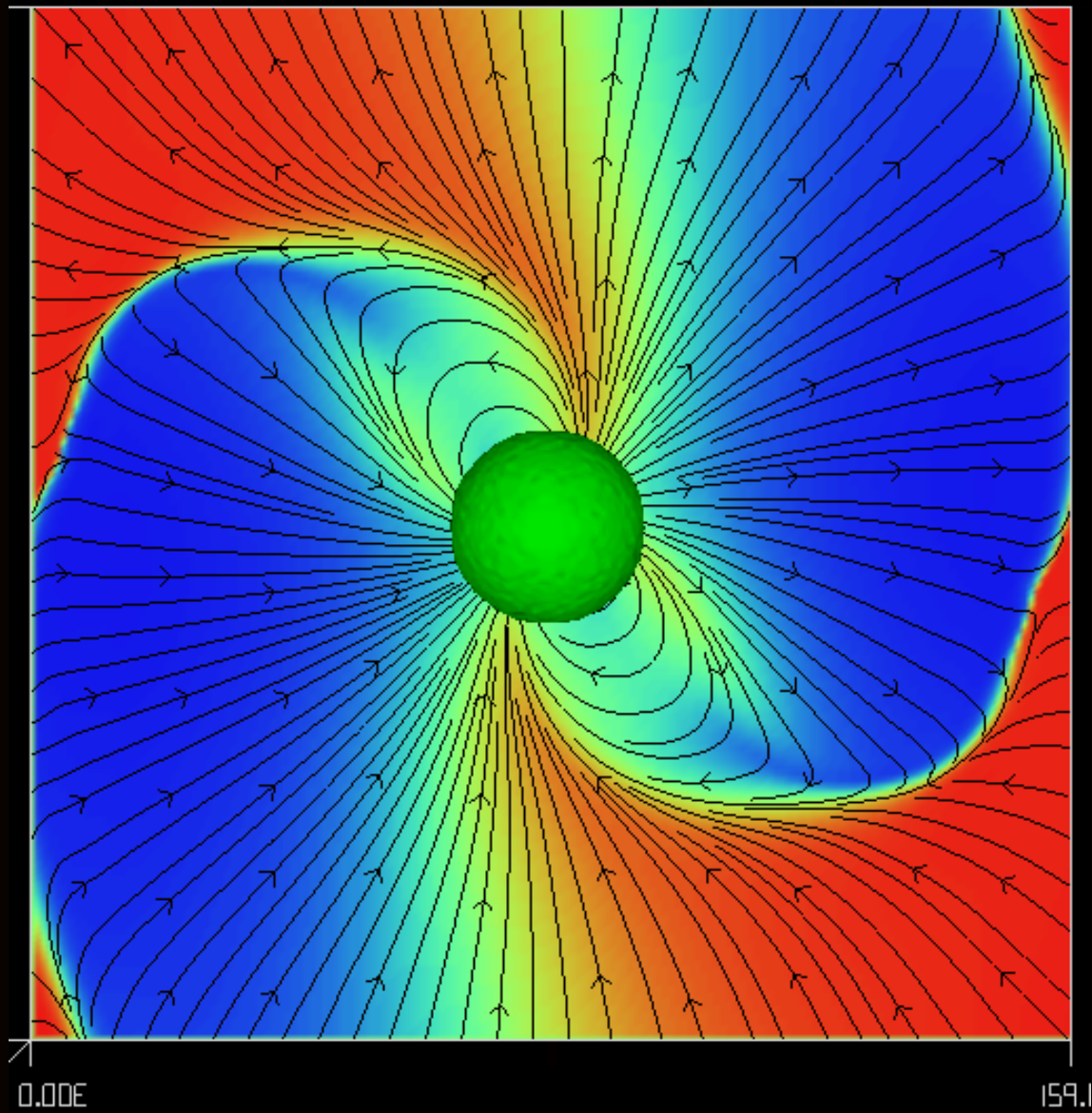
$$\dot{E} = \frac{\mu^2 \Omega^4}{c^3} (1 + \sin^2 \theta)$$

$$\dot{E}_{vac} = \frac{2}{3} \frac{\mu^2 \Omega^4}{c^3} \sin^2 \theta$$

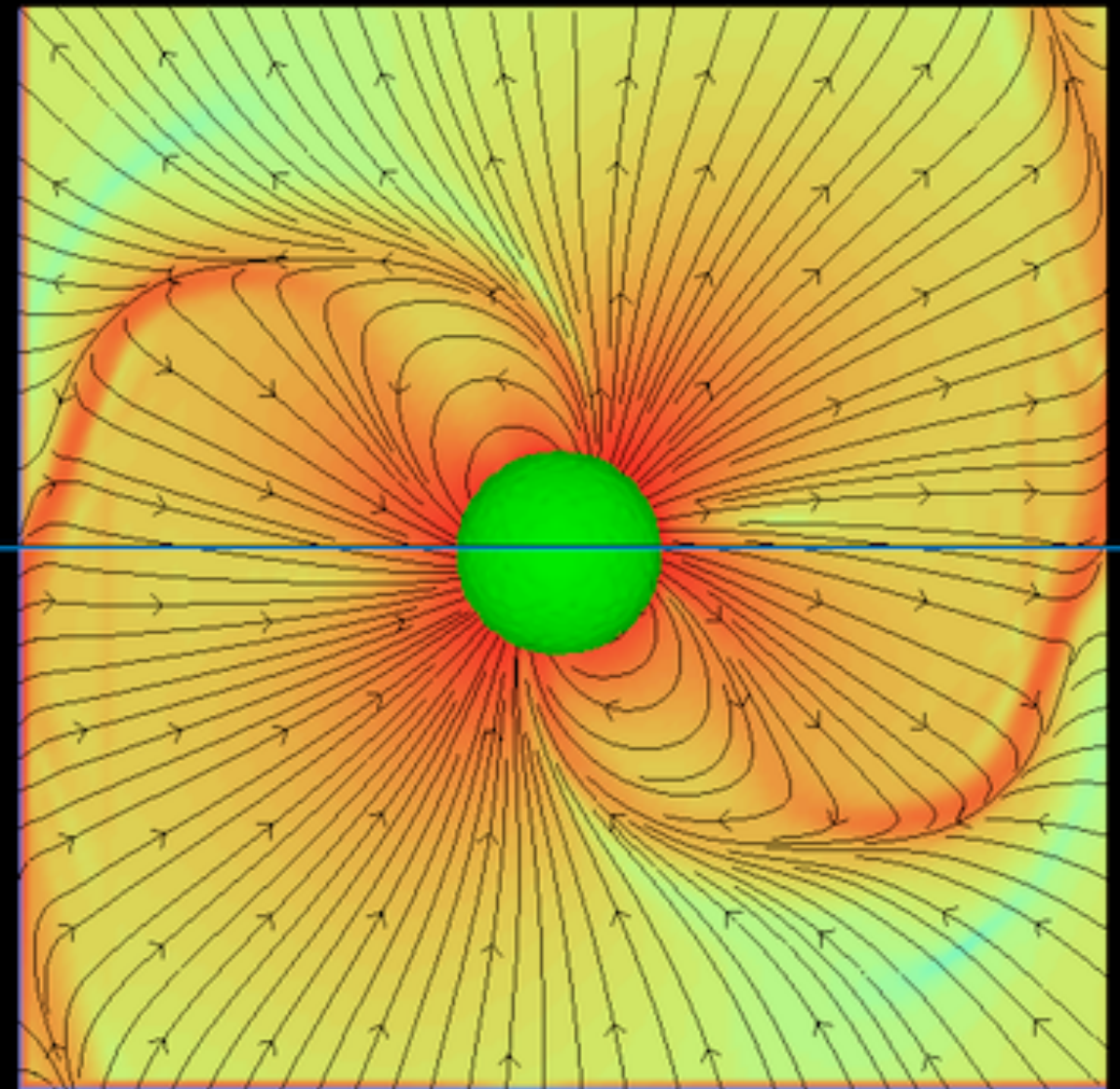
A.S.'06; also confirmed by Kalapotharakos & Contopoulos 09

IN COROTATING FRAME

60 degree inclination

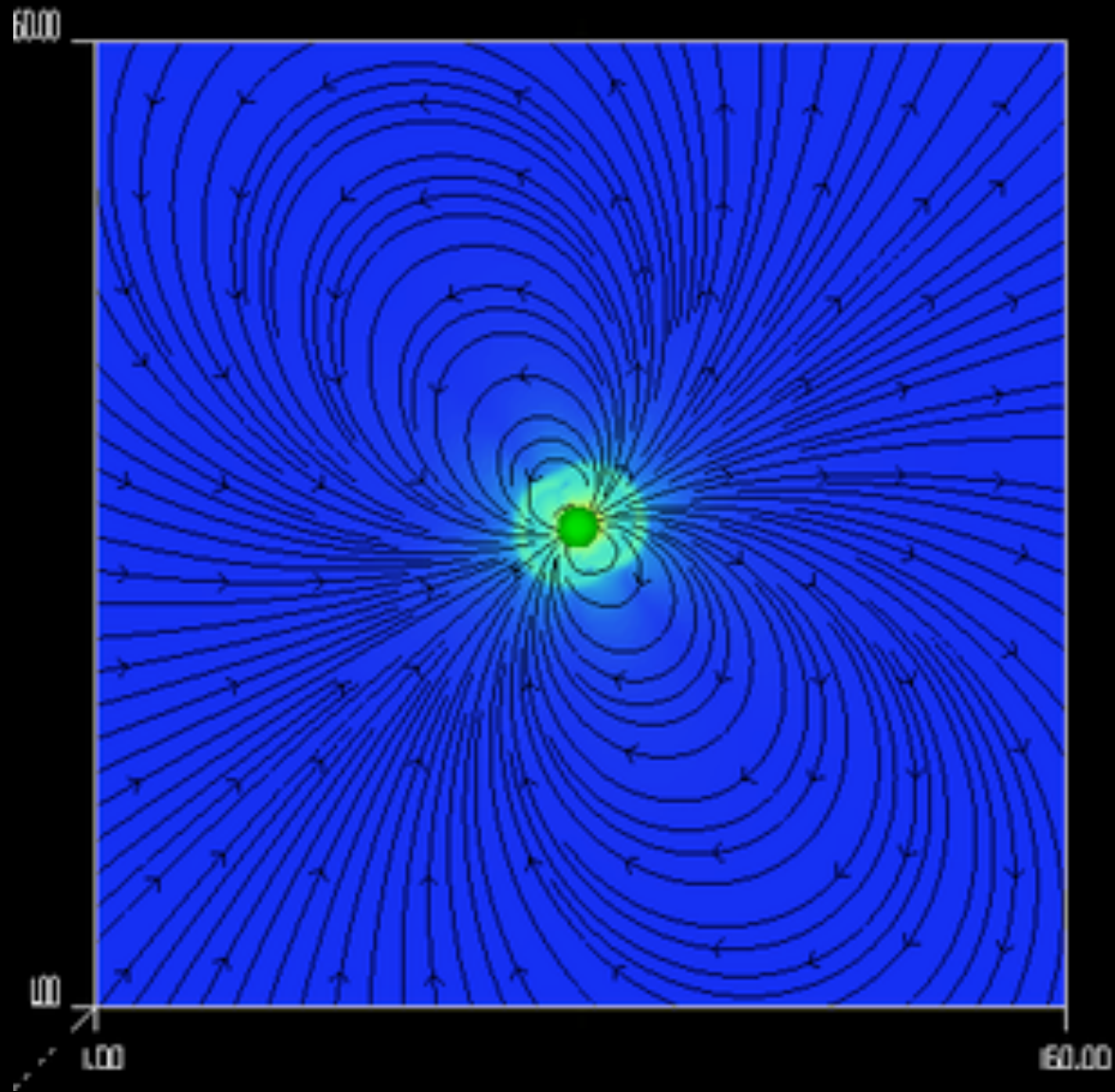


Force-free



Force-free current density

3D force-free magnetosphere: 60 degrees inclination



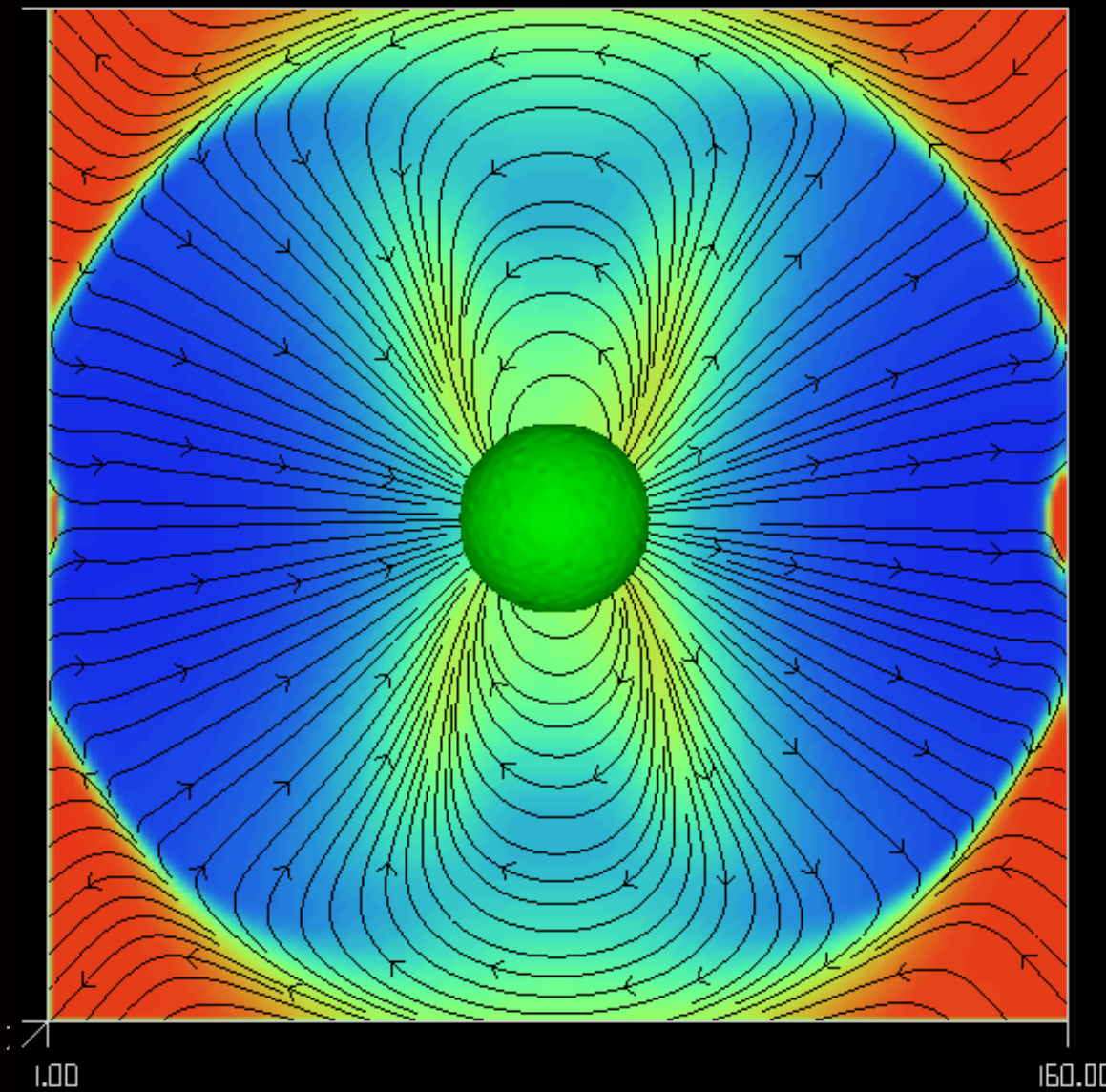
60 degrees force-free current



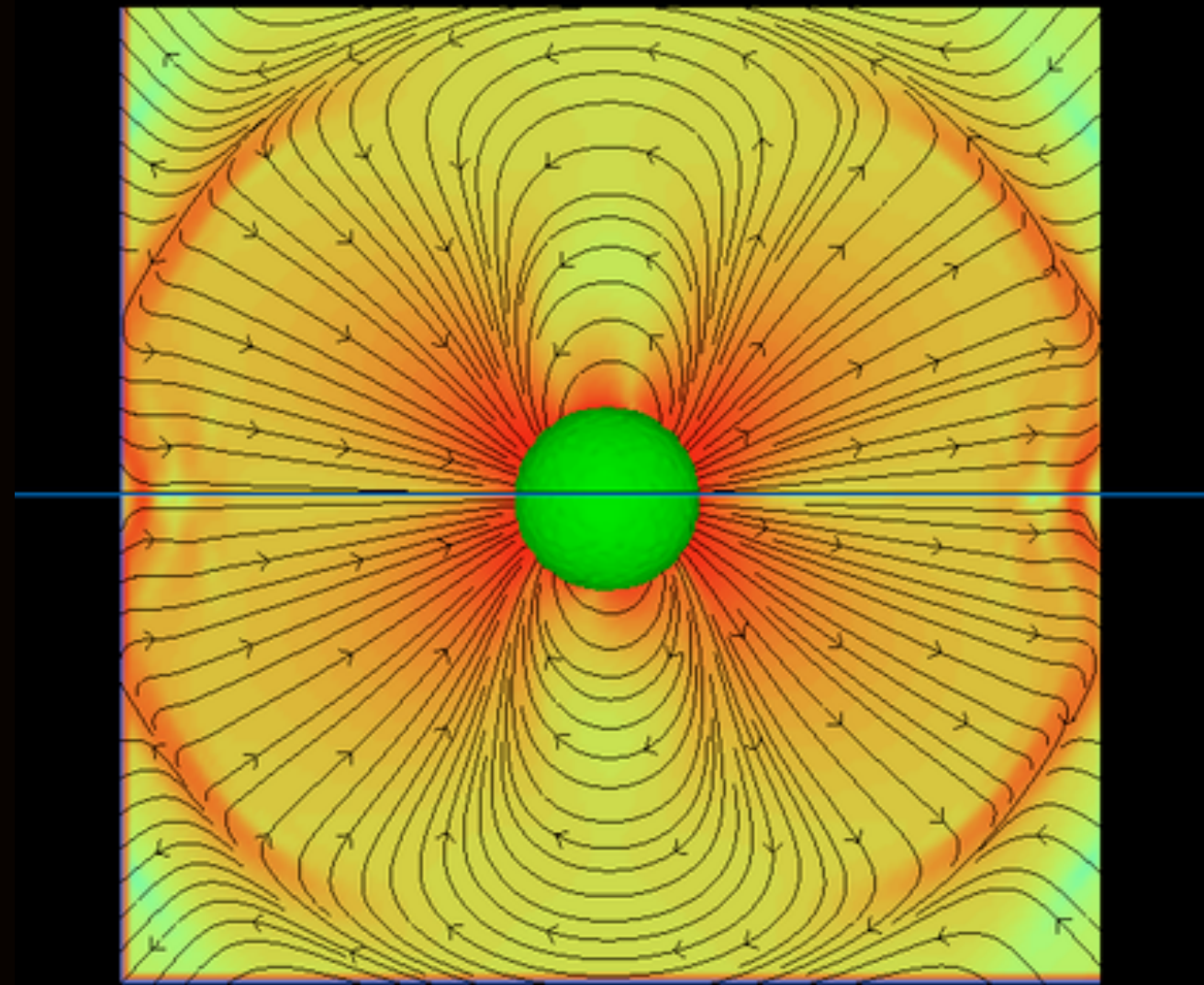
Similar to heliospheric current sheet

IN COROTATING FRAME

90 degree inclination

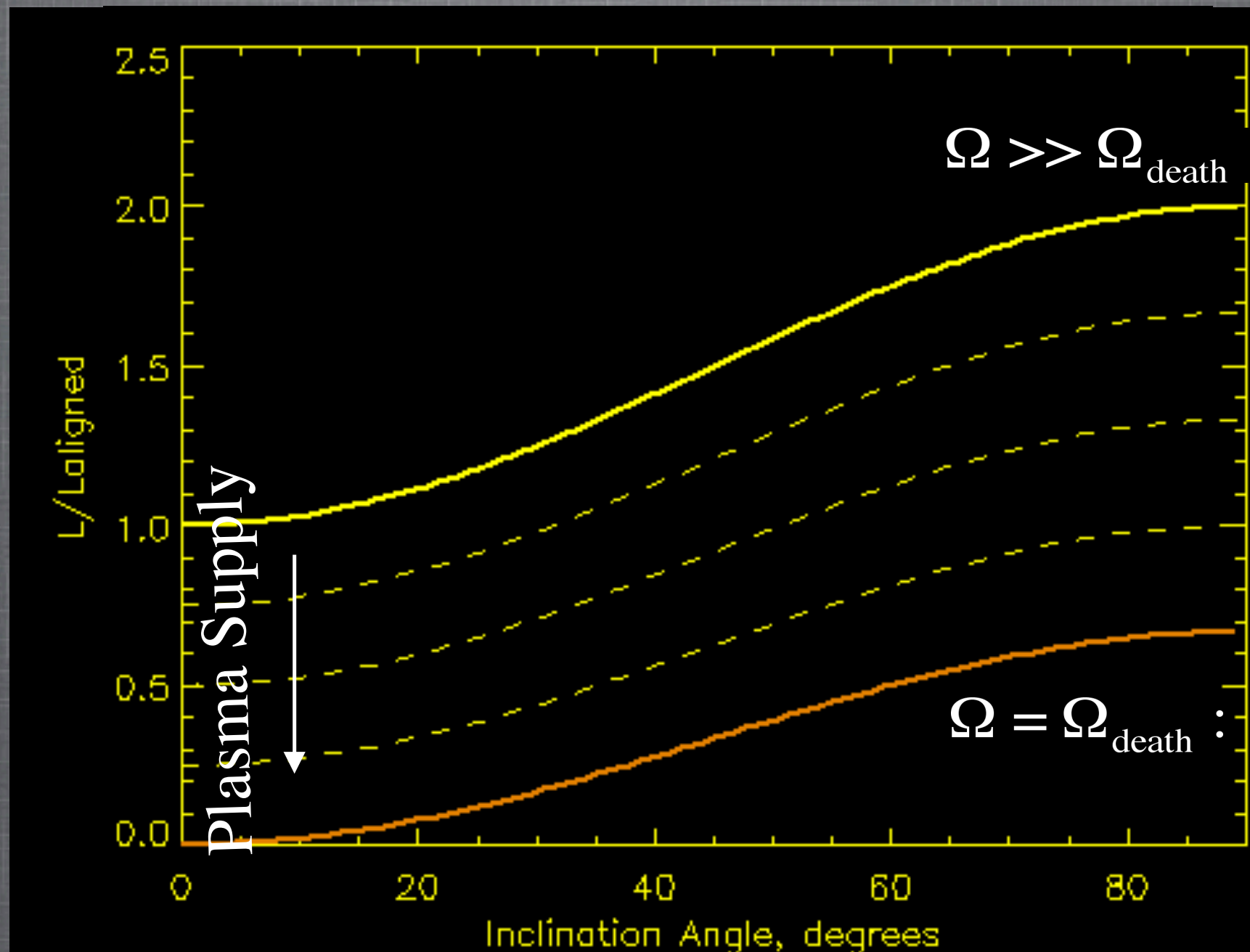


Force-free



Force-free current density

SPIN-DOWN POWER



There is a continuum of solutions that depend on plasma supply. These can be characterized by the presence of accelerating E field, or resistivity.

Resistive force-free

- ✳ There is a continuum of solutions between vacuum and ideal conducting force-free magnetosphere if plasma is not perfect everywhere.

- ✳ Can parameterize these with resistivity in the proper frame.

- ✳ Nice feature: re-emergence of parallel E field.

Ohm's law in the proper frame:

$$\vec{j}_{\text{fluid}} = \sigma \vec{E}_{\text{fluid}}$$

In lab frame:

$$\vec{j} = \frac{\rho c \vec{E} \times \vec{B}}{B^2 + E_0^2} + \frac{(-\beta_{||} \rho c + \sqrt{\frac{B^2 + E_0^2}{B_0^2 + E_0^2}} (1 - \beta_{||}^2) \sigma E_0) (B_0 \vec{B} + E_0 \vec{E})}{B^2 + E_0^2}$$

$$B_0^2 = \frac{\vec{B}^2 - \vec{E}^2 + \sqrt{(\vec{B}^2 - \vec{E}^2)^2 + 4(\vec{E} \cdot \vec{B})^2}}{2},$$
$$E_0 = \sqrt{B_0^2 - \vec{B}^2 + \vec{E}^2},$$
$$B_0 = \text{sign}(\vec{E} \cdot \vec{B}) \sqrt{B_0^2},$$

cf. Lyutikov 03
Gruzinov 07-11

Li, AS, Tchekhovskoy, 2011

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Minimal || velocity frame:

$$\vec{j} = \frac{\rho c \vec{E} \times \vec{B} + \sqrt{\frac{B^2 + E_0^2}{B_0^2 + E_0^2}} \sigma E_0 (B_0 \vec{B} + E_0 \vec{E})}{B^2 + E_0^2}.$$

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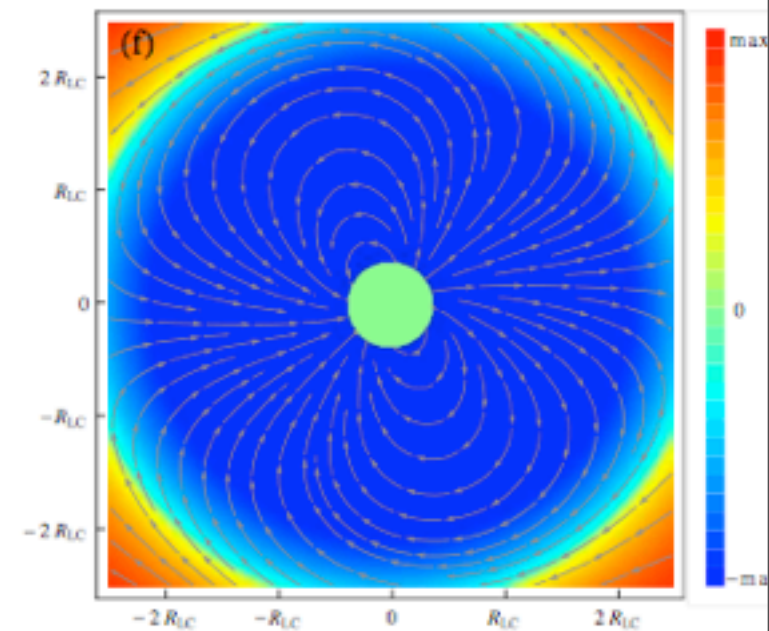
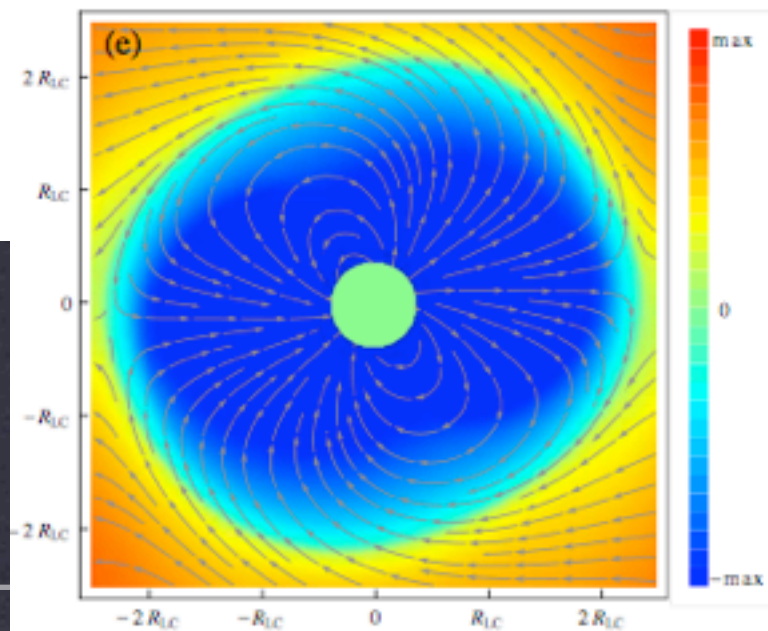
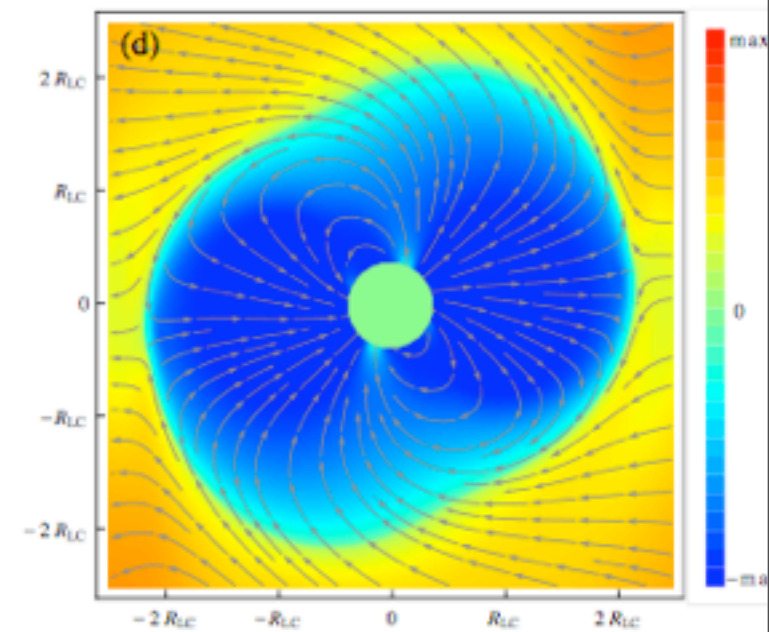
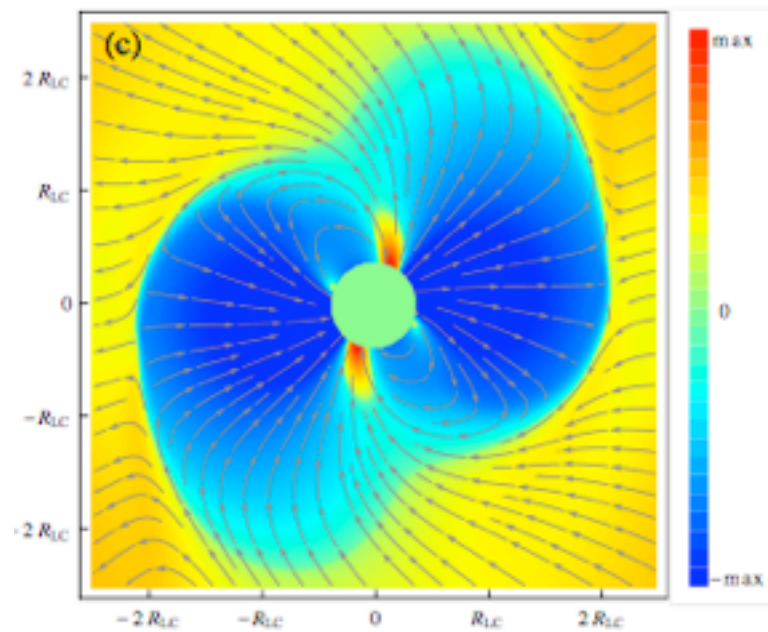
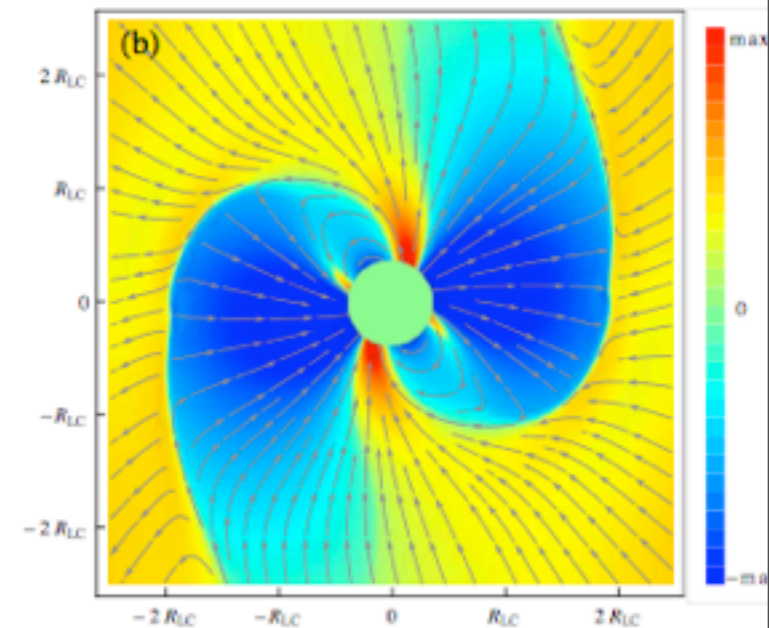
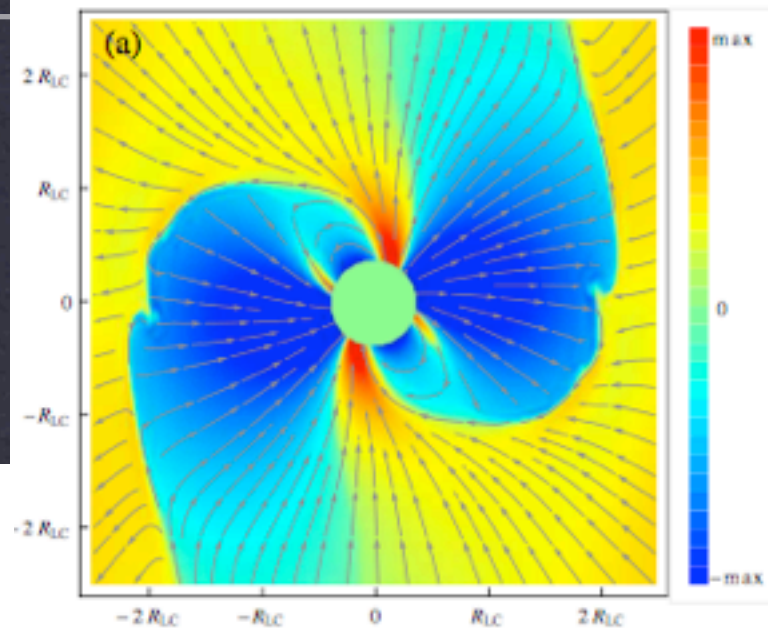
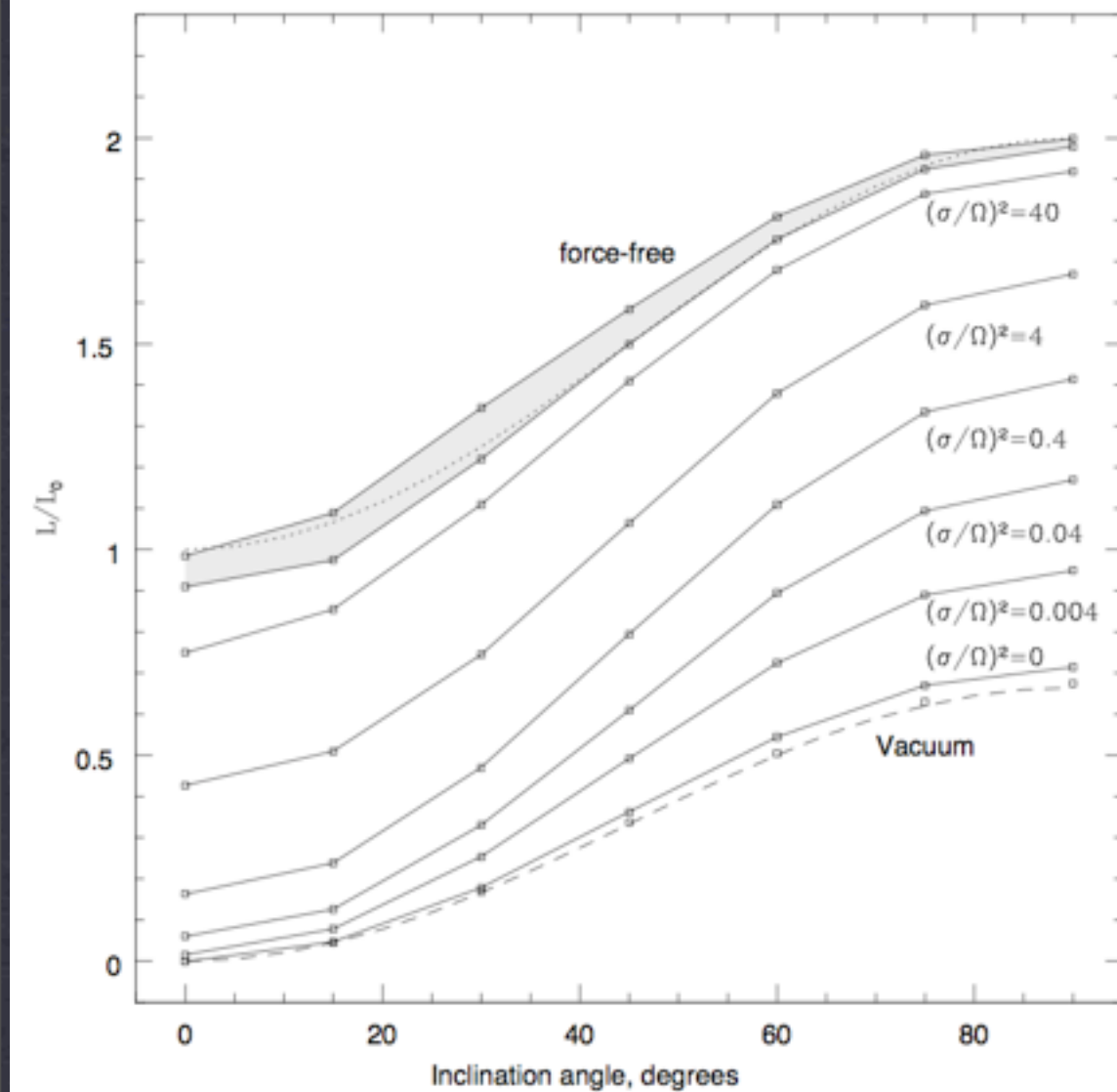
$$B_0 = \text{sign}(\vec{E} \cdot \vec{B}) \sqrt{B_0^2},$$

cf. Lyutikov 03
Gruzinov 07-11

Li, AS, Tchekhovskoy, 2011
also, Kalapotharakos et al 11

Resistive:

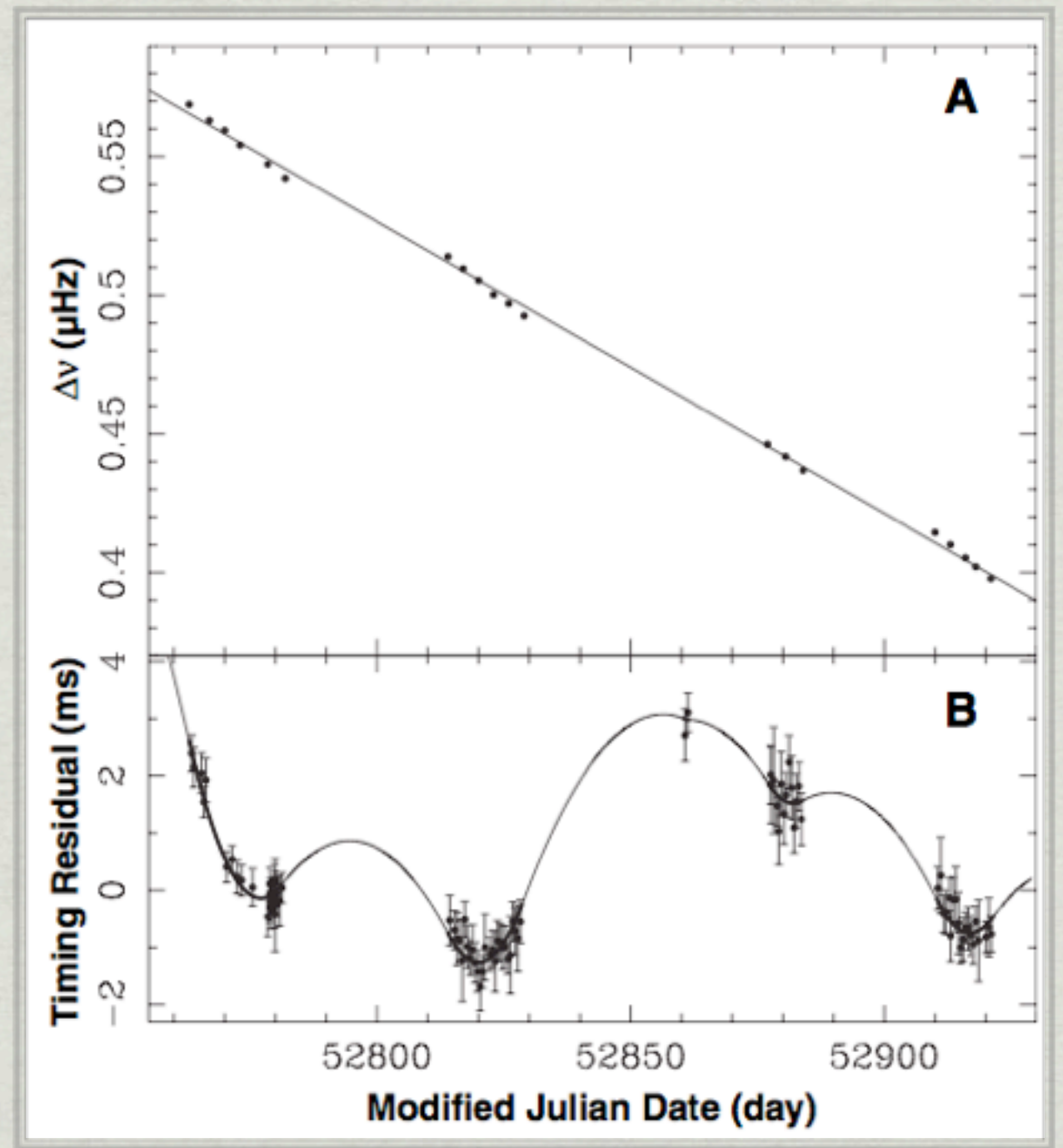
Vary σ/Ω



Spin-down power

Application: intermittent pulsars

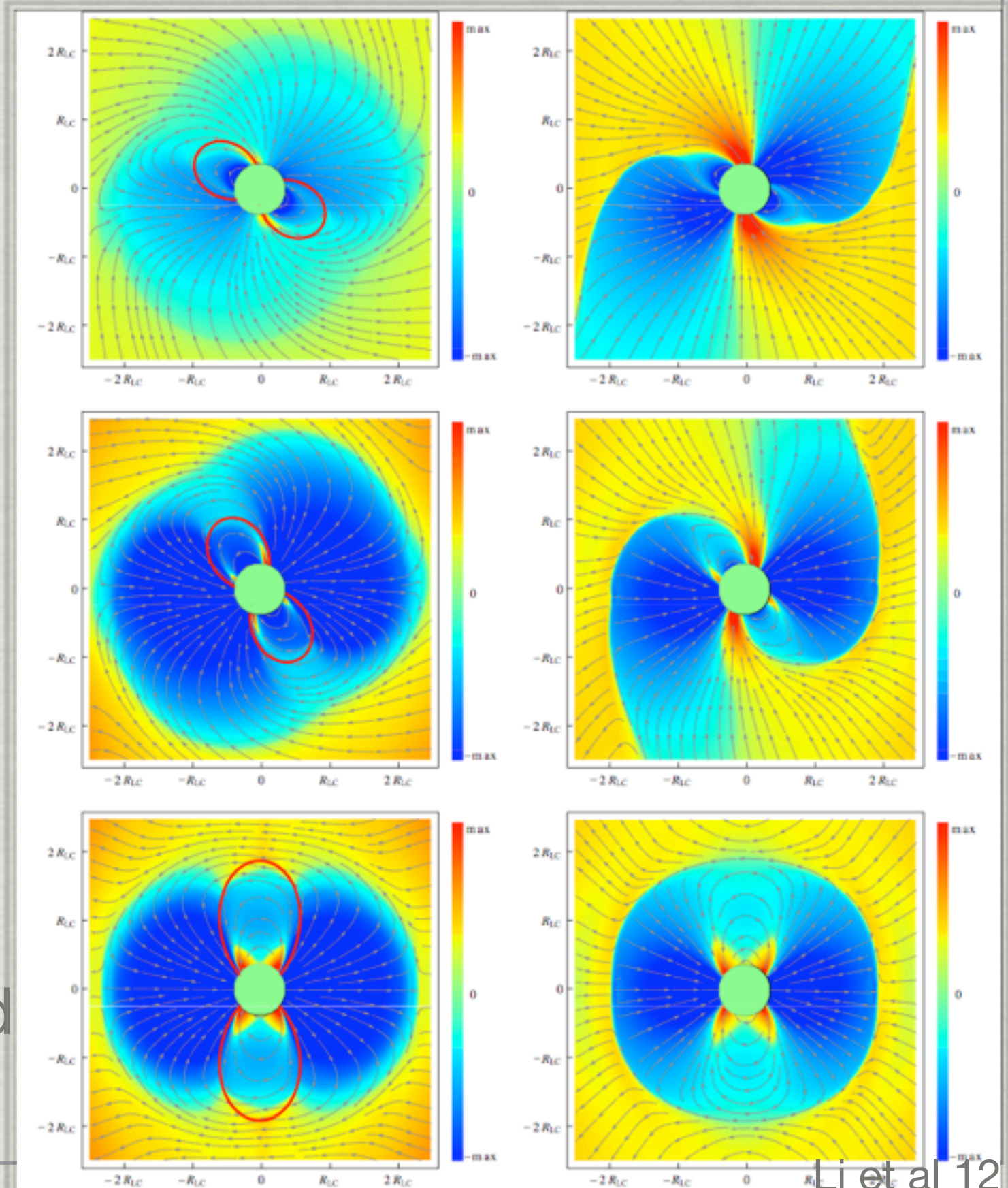
- * Intermittent pulsars display changes in spin-down power when they are ON and OFF in radio by factor >1.5
- * One possibility: conducting closed zone, vacuum-like open zone; Interrupted plasma production



Kramer et al 06

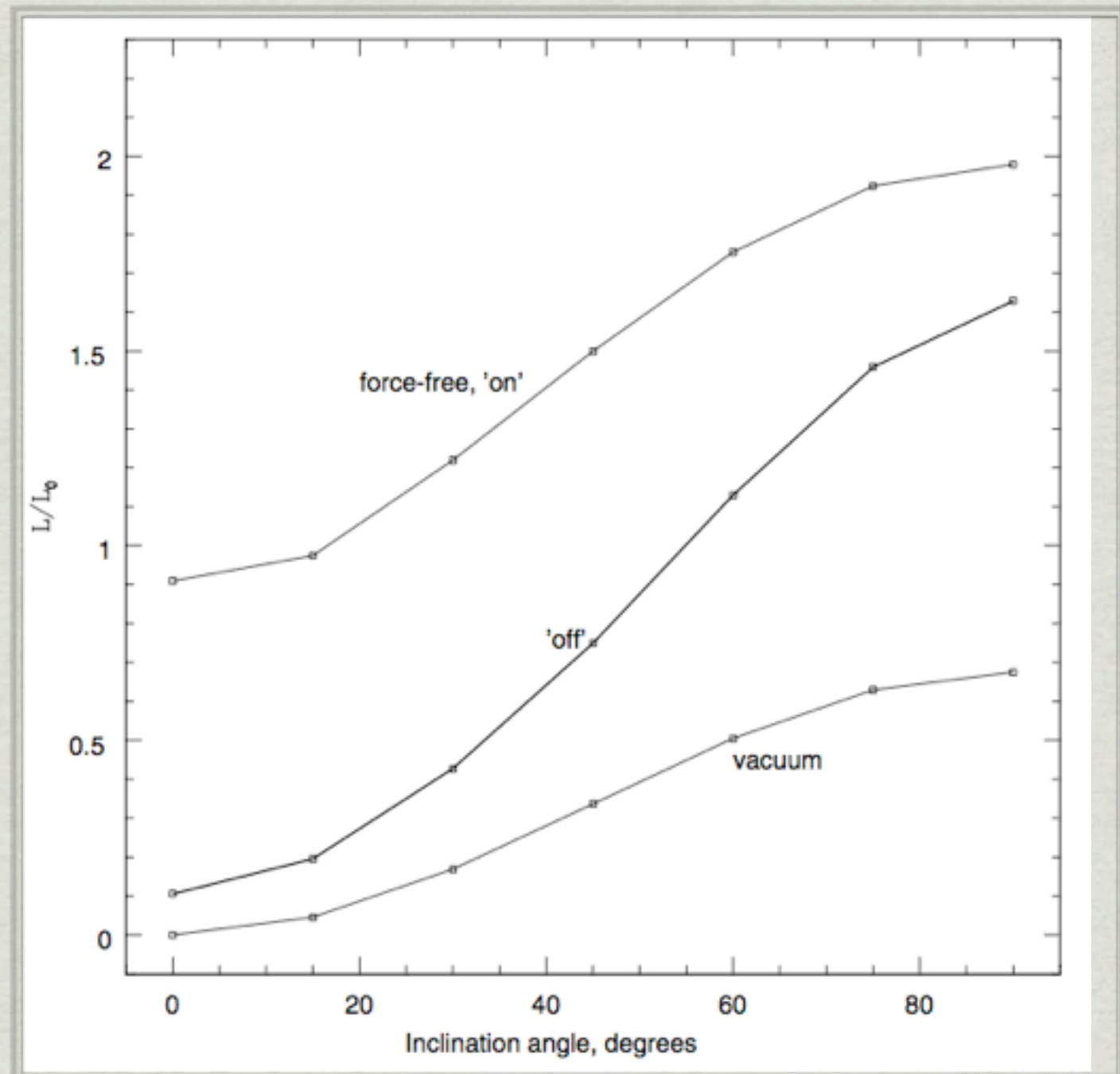
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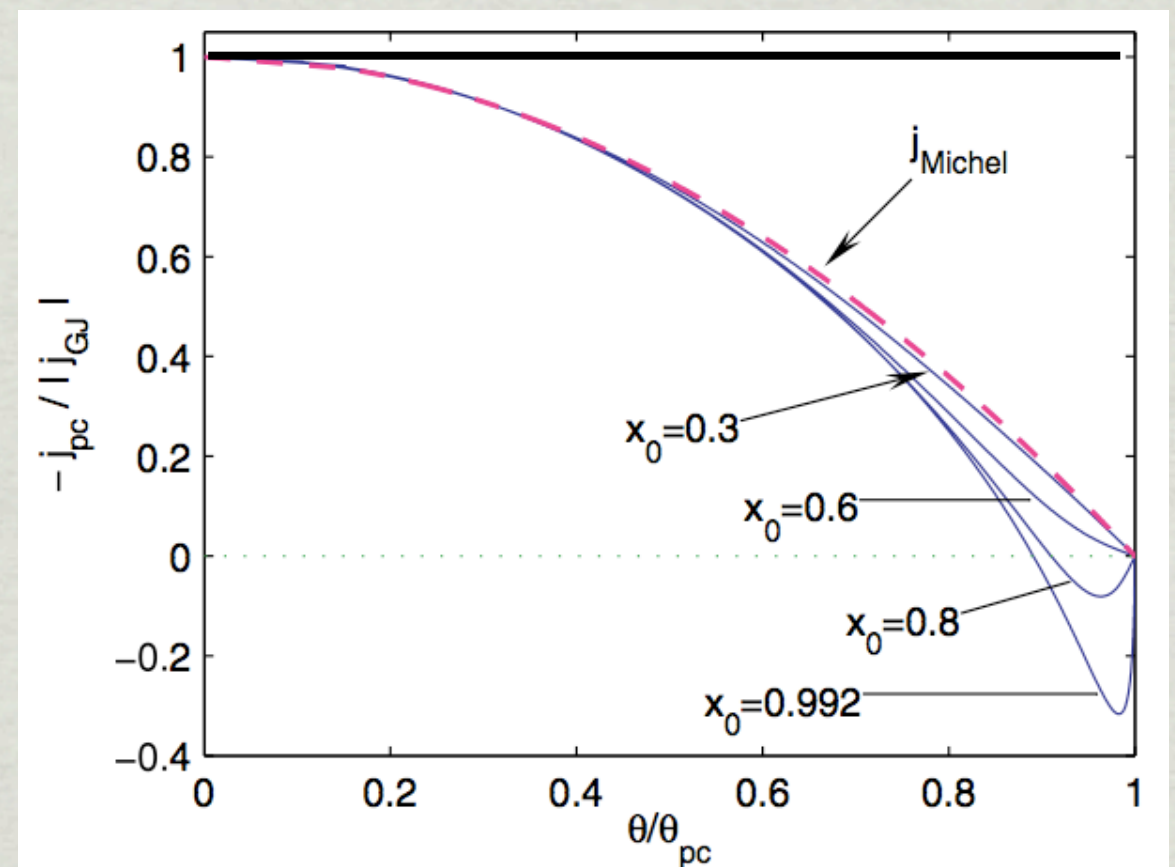
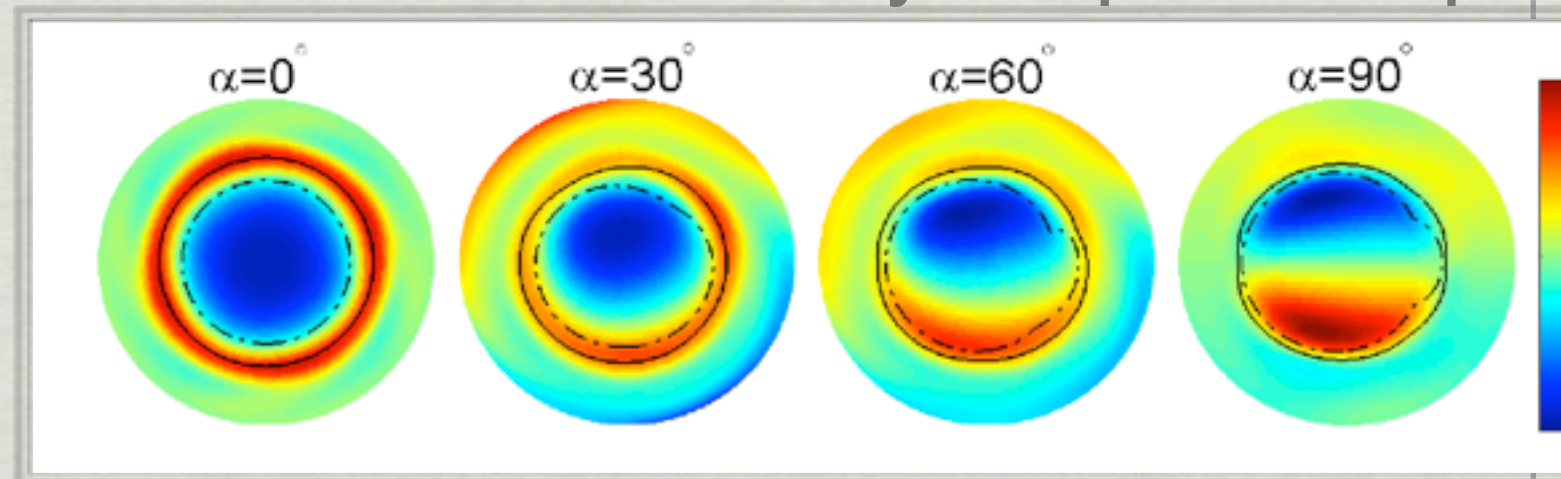
- * Factor of > 1.5 can be explained with “hybrid” vacuum-conducting magnetosphere.
- * The physical origin of switch is completely unclear.



Recent advances:

- ✳ Better understanding of the cascades and driving of current (see Timokhin's talk)
- ✳ Polar cascades can supply current both $>$ and $<$ than GJ current. Implies time-dependence. Interestingly, current $<$ GJ does not pair produce.

Current density on polar cap



$j < j_{GJ}$

Timokhin 06

Recent advances:

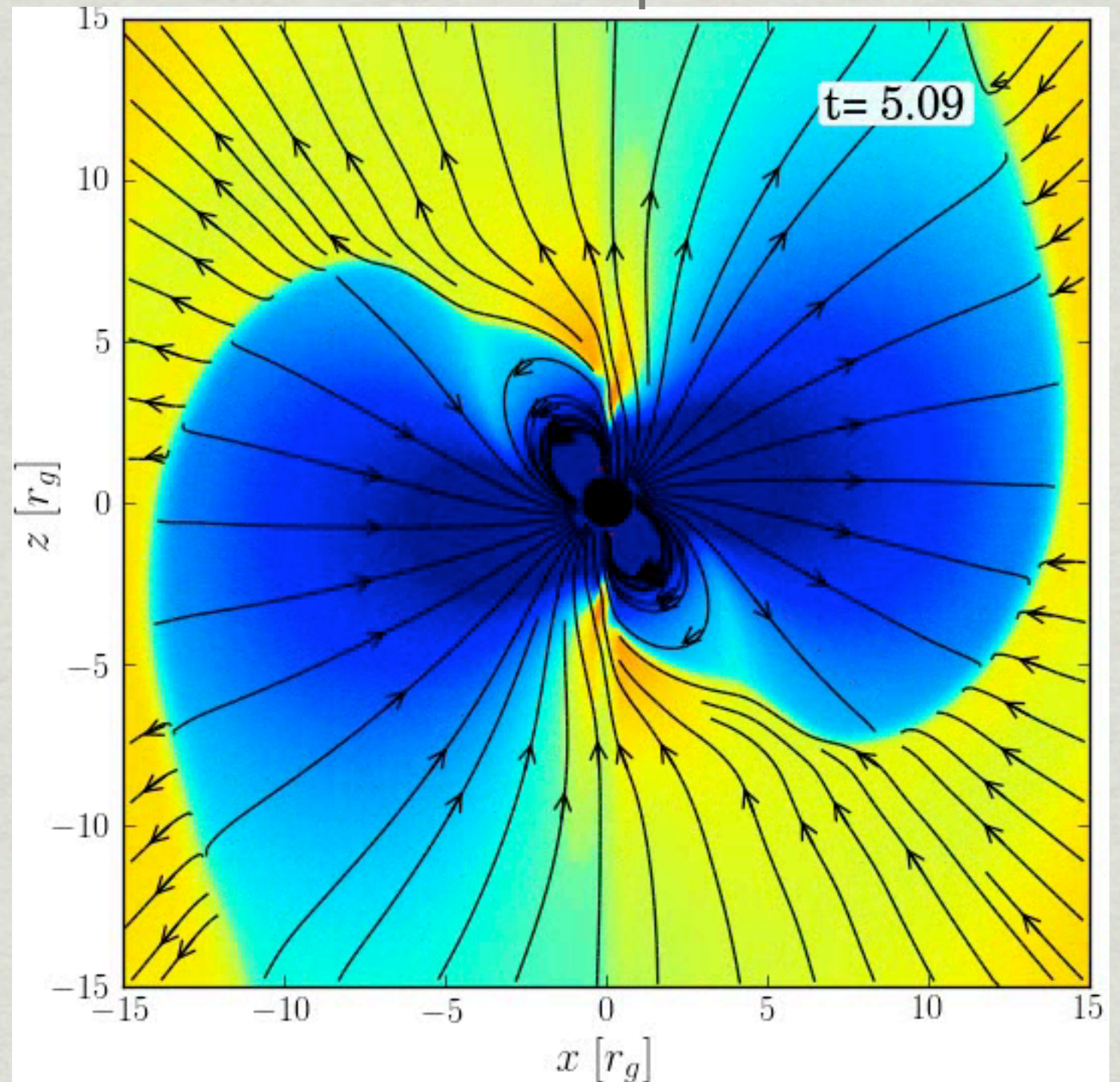
- ✱ Full RMHD is now in 3D!
- ✱ Oblique rotator can now be studied in ideal MHD
(Tchekhovskoy, AS, Li 2013)



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- ✱ Oblique rotator can now be studied in ideal MHD (Tchekhovskoy, AS, Li 2013)
- ✱ Spherical grid which allows non-axisymmetric solutions. Magnetization > 100 . Fixed magnetization inside 0.7 LC

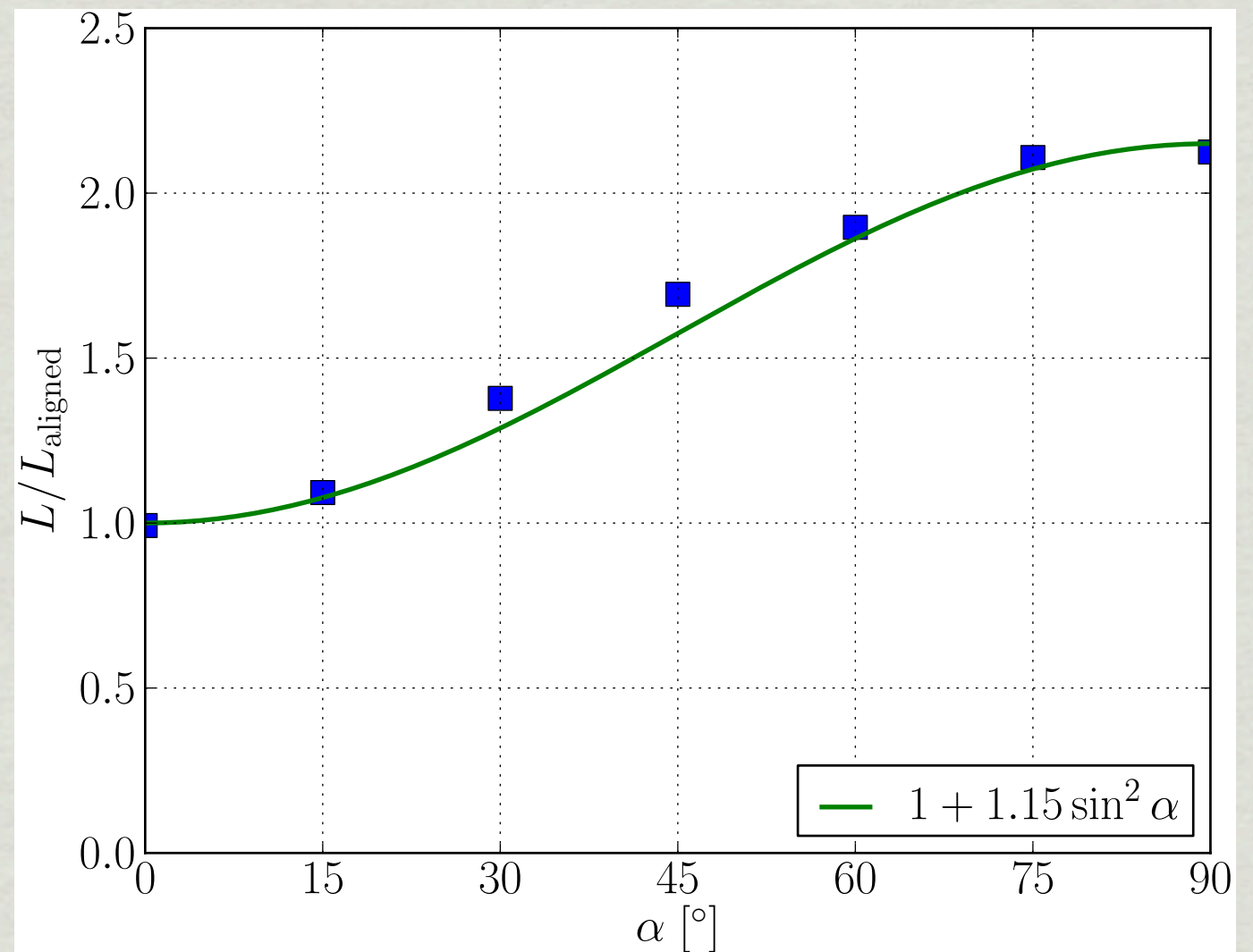
color: out of plane B field



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Spin down luminosity



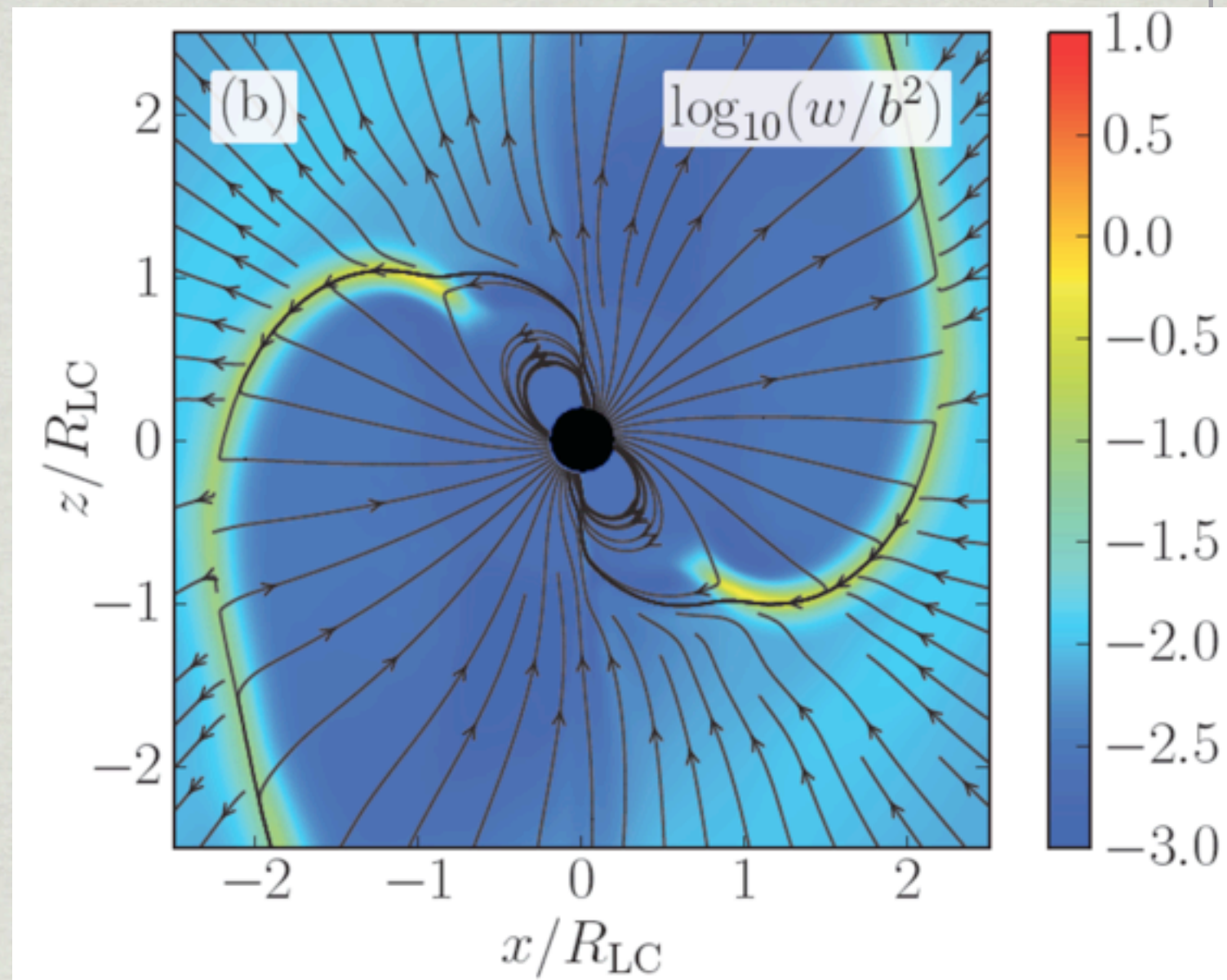
Obliqueness

Variation with angle is similar to force-free

Recent advances:

Magnetization

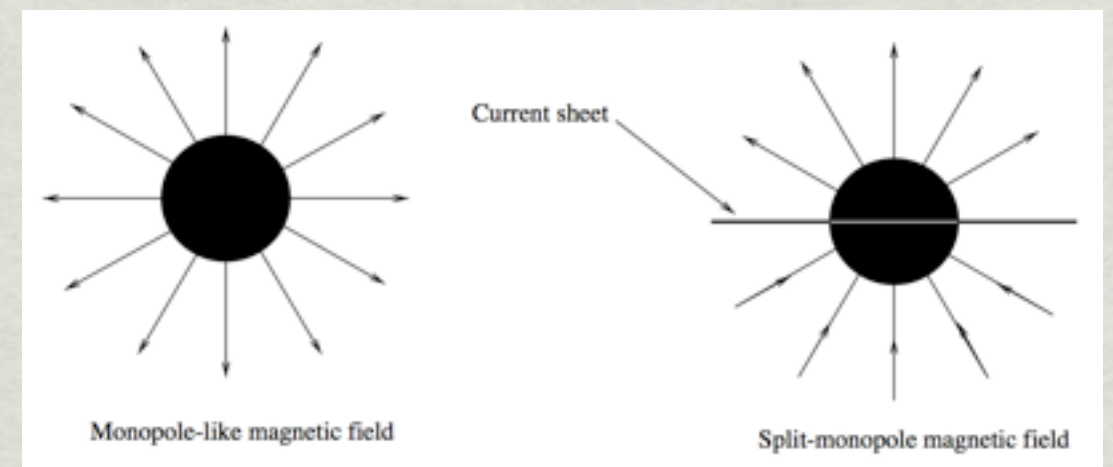
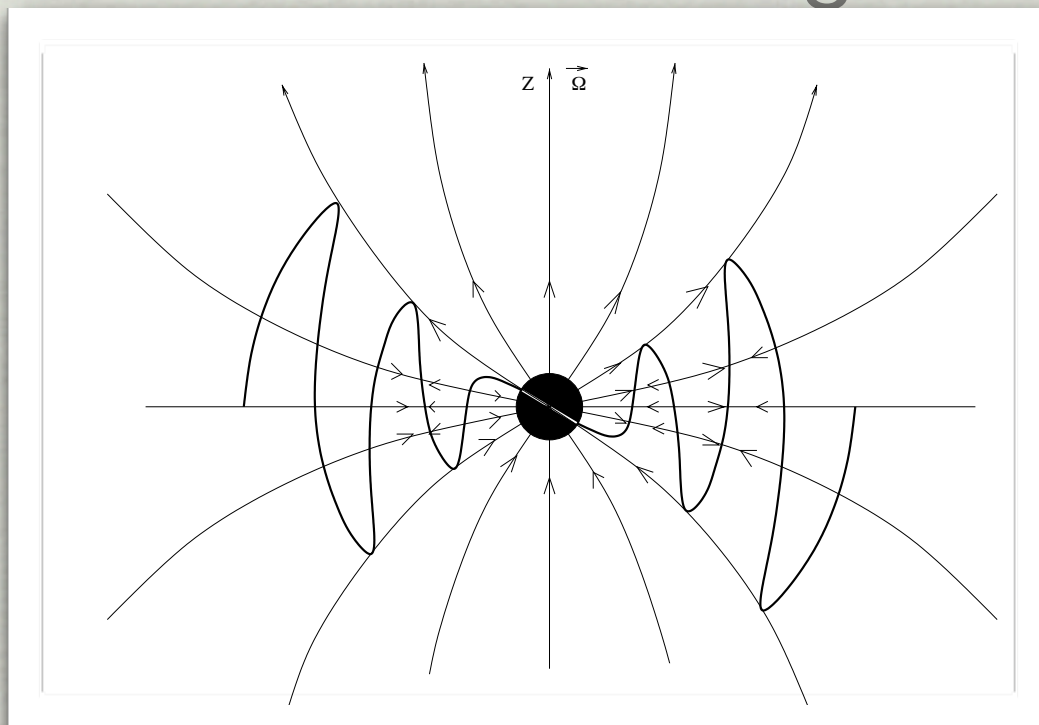
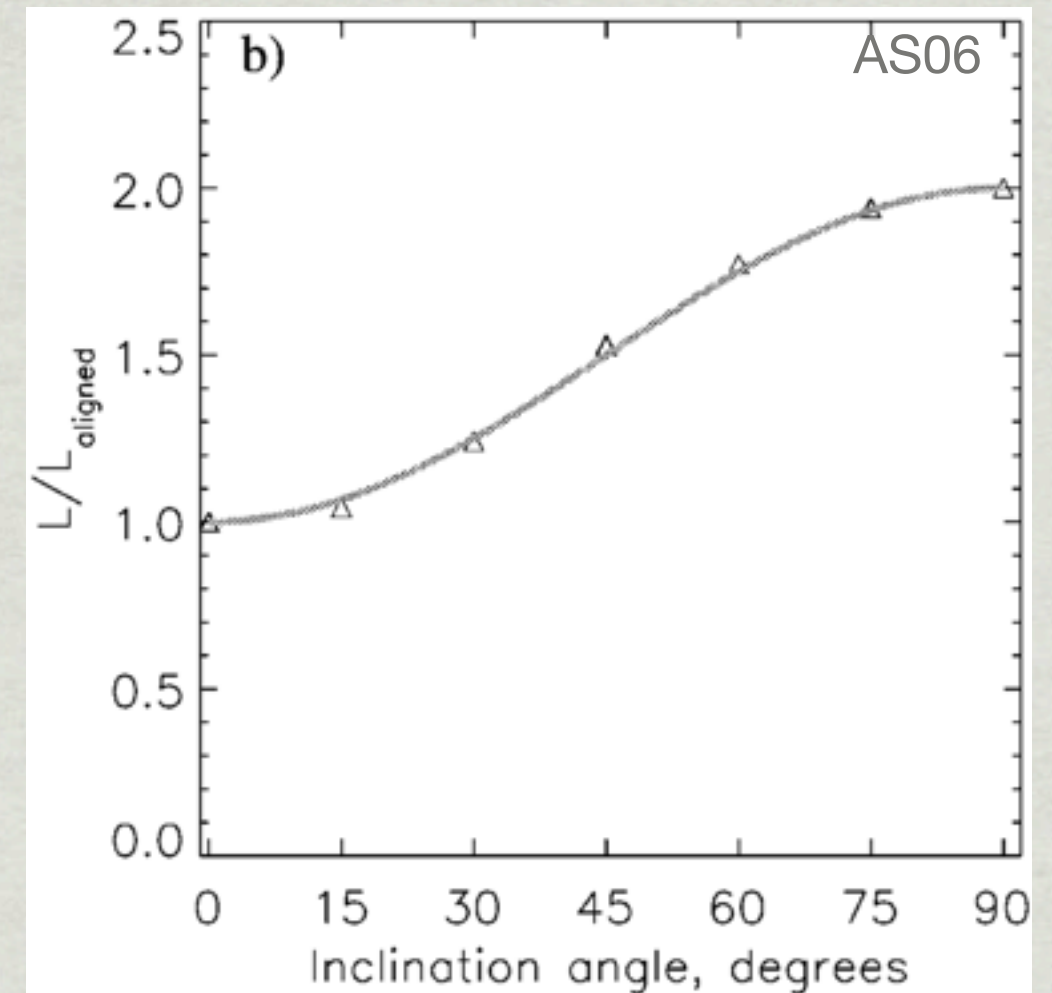
- * Thermal pressure support in the sheet
- * Low dissipation, $<10\%$ (cf. Gruzinov's 50%)
- * Big uncertainty: velocity along the field lines is not self-consistent (depends on cascades).



Current sheet is formed and has gas pressure in it

More on the magnetosphere

- ✳ Can we understand $1 + \sin^2 \alpha$ dependence of spin-down?
- ✳ Bogovalov 1999 split monopole: spin-down constant with angle!

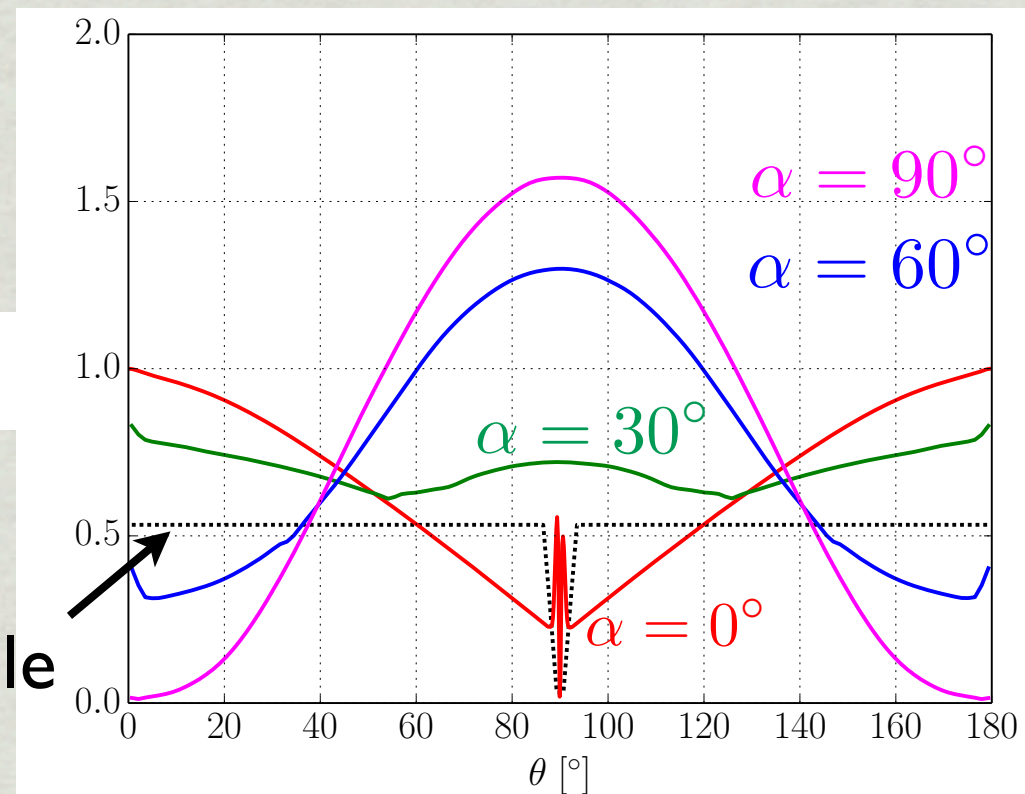


Are asymptotic field lines like split-monopole?

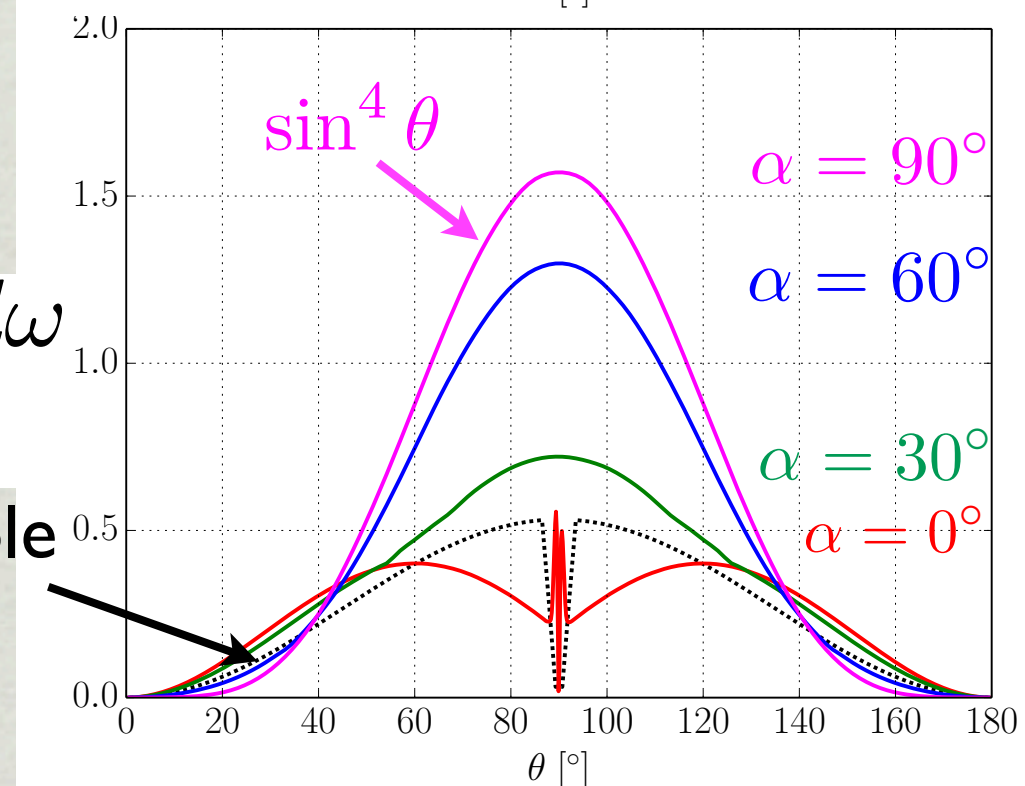
More on the magnetosphere

- ✱ B-field is equatorially-concentrated
- ✱ Wind luminosity is more equatorially concentrated than monopole
- ✱ This effect needs to be included for gamma-ray emission light curve calculation and PWN models.

$\langle B_r^2 \rangle$
split-monopole



$dL/d\omega$
split-monopole
 $\sin^2 \theta$



Tchekhovskoy, Spitkovsky, Li 2013,
Tchekhovskoy, Philippov, Spitkovsky in prep.

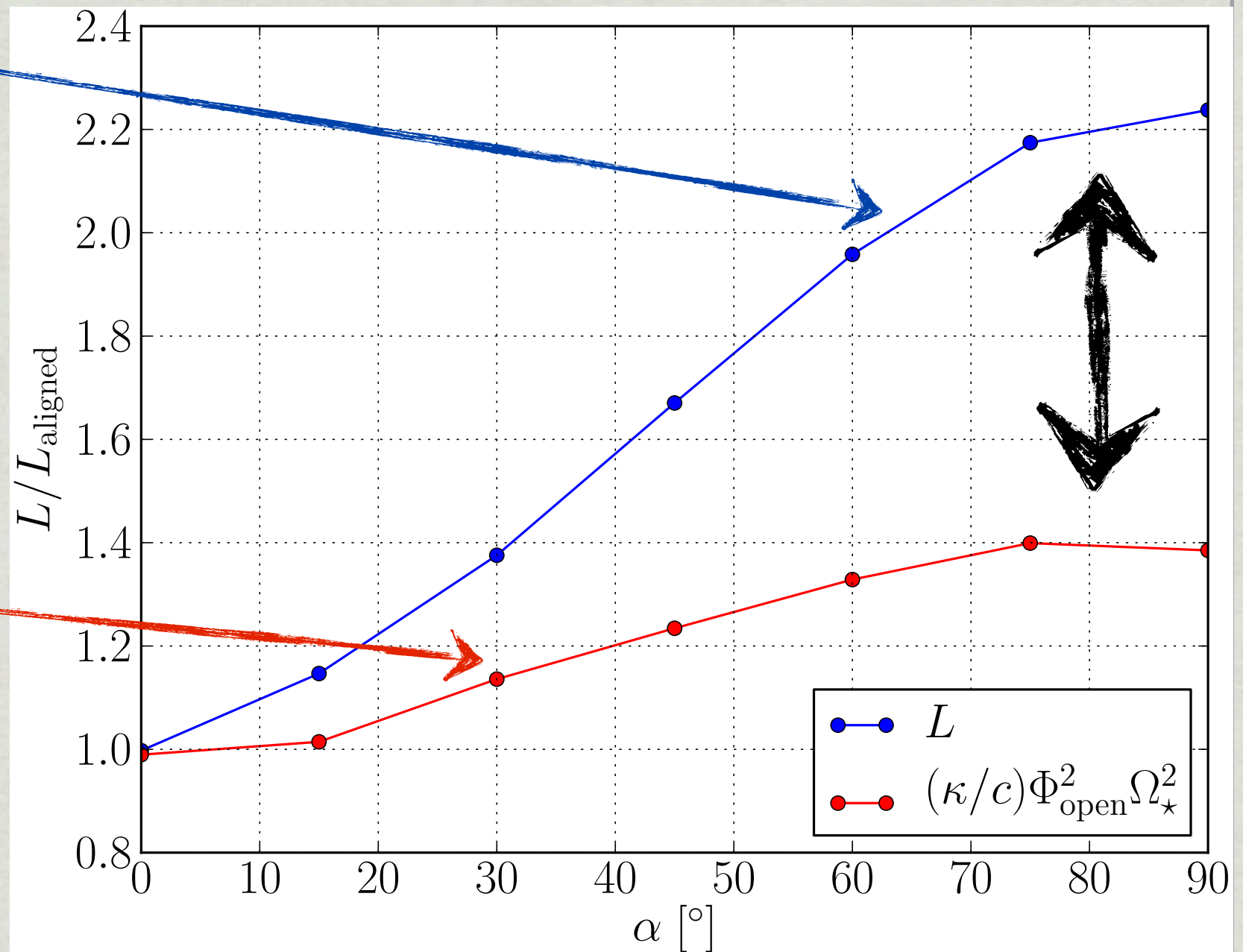
Field Non-uniformity Explains Enhanced Spindown of Oblique Pulsars

Enhanced
spindown due to
non-uniformity
of B-field?

(Spitkovsky'06, Petri'12, AT, Spitkovsky, Li'13)

Assumption of
uniform B-field
under-predicts
spindown

$$P_{\text{NS}} = \frac{\kappa}{c} \Phi_{\text{open}}^2 \Omega_{\star}^2$$



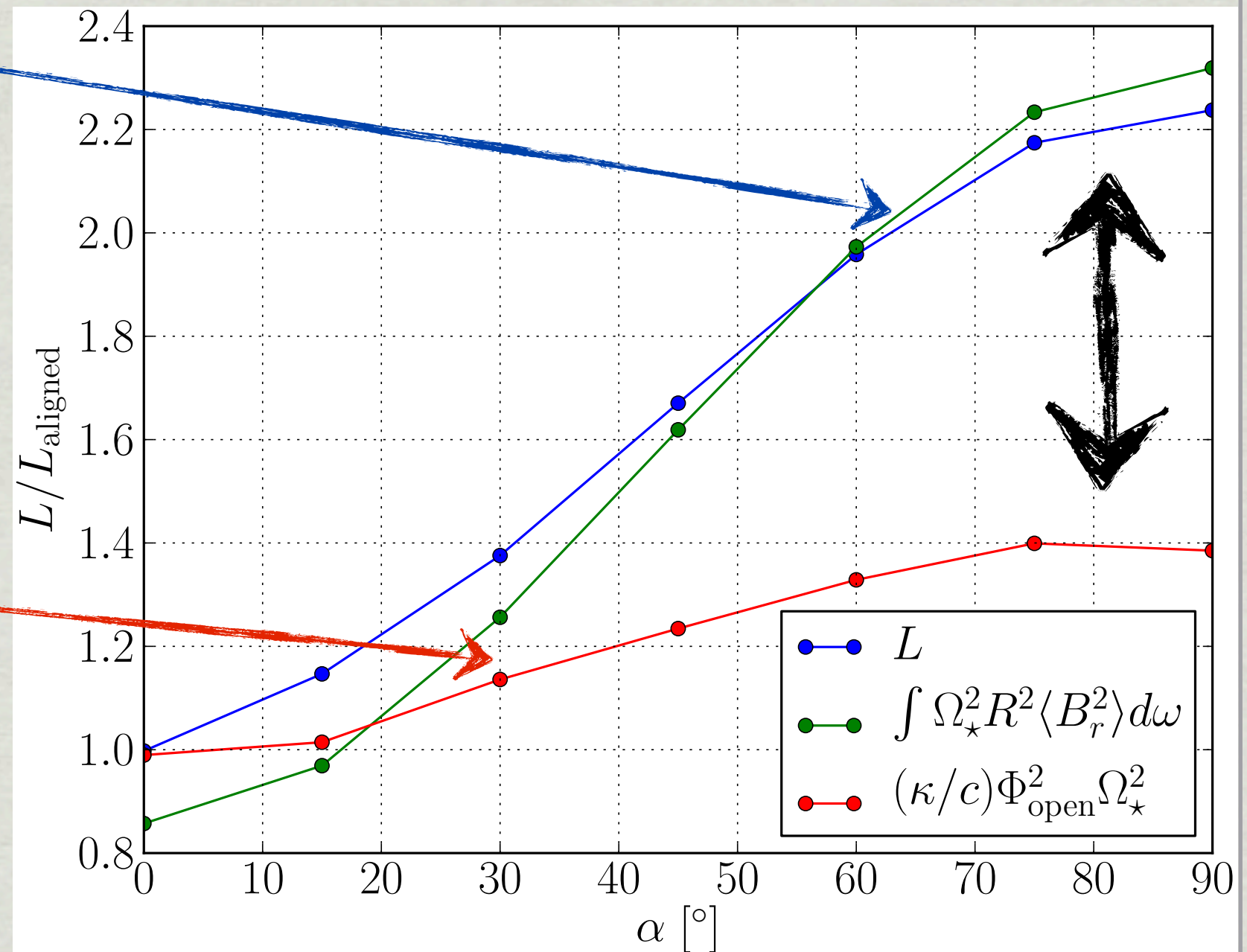
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Tchekhovskoy, Philippov, Spitkovsky in prep.

just B_r variation from inclined dipolar field gives $1 + \sin^2 \alpha$

Analytic fitting model of 3D pulsar wind

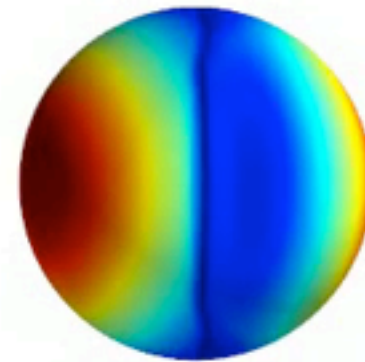
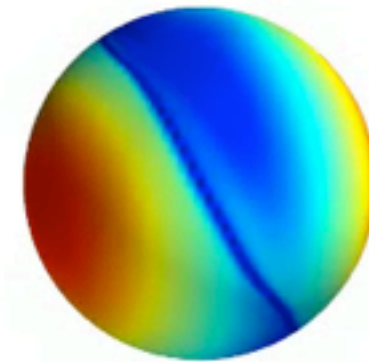
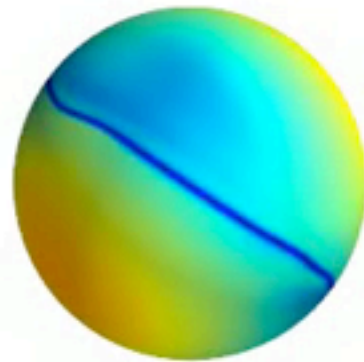
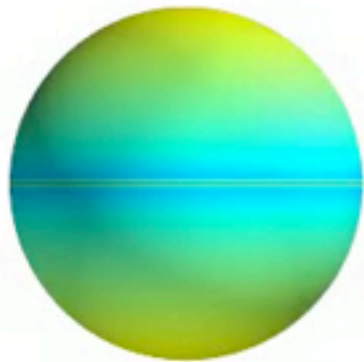
$\alpha = 0^\circ$

$\alpha = 30^\circ$

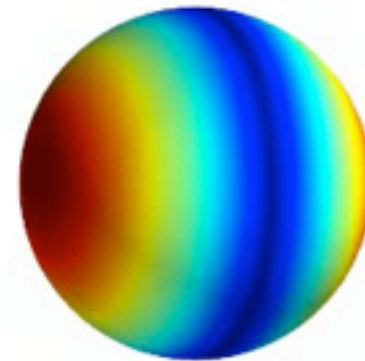
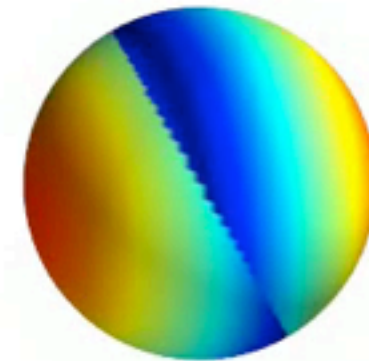
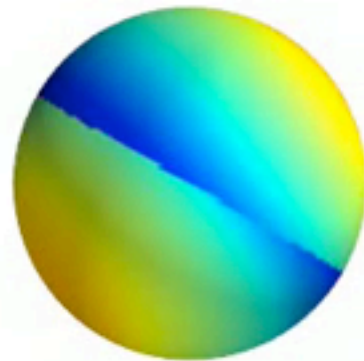
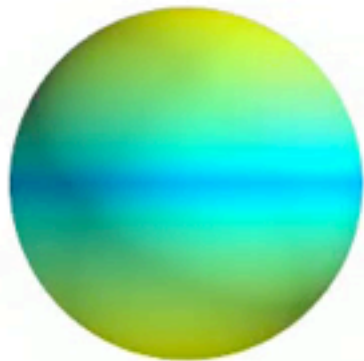
$\alpha = 60^\circ$

$\alpha = 90^\circ$

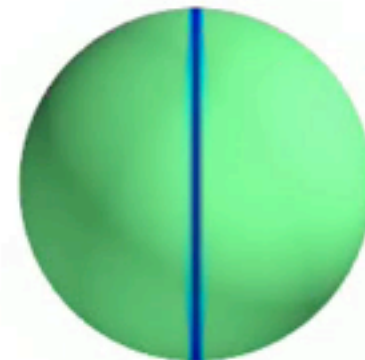
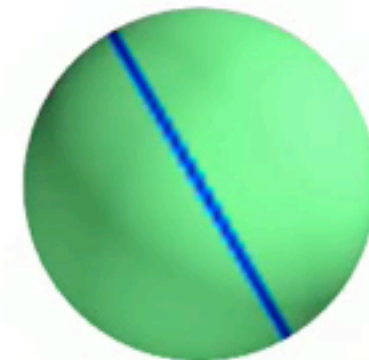
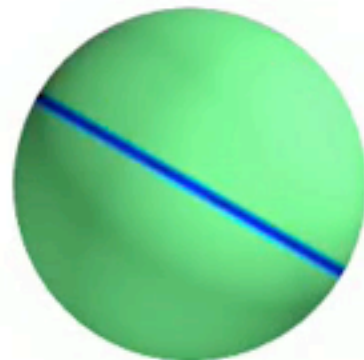
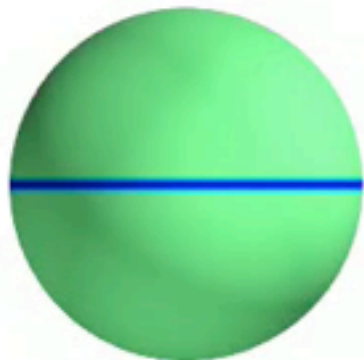
MHD
simulation



Analytic
Model



Oblique split-
monopole
(Bogovalov 1999)



$|B_r|$



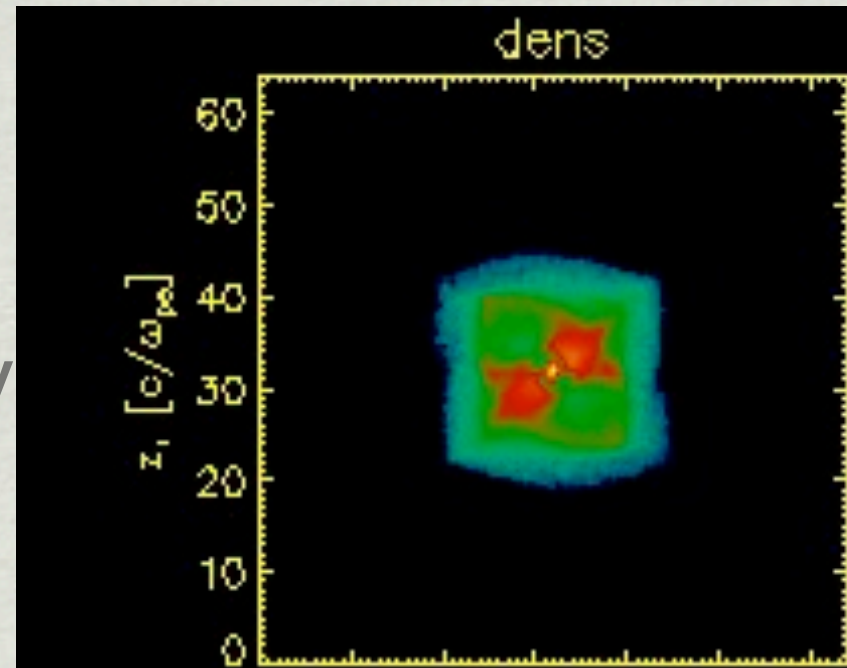
Fitting model for oblique pulsar wind is now available

More recent advances:

Vacuum field + test particles

- * Full particle modeling with PIC (particle-in-cell) in 3D (AS in prep).
- * Idea: supply enough neutral plasma and “let it figure it out.”

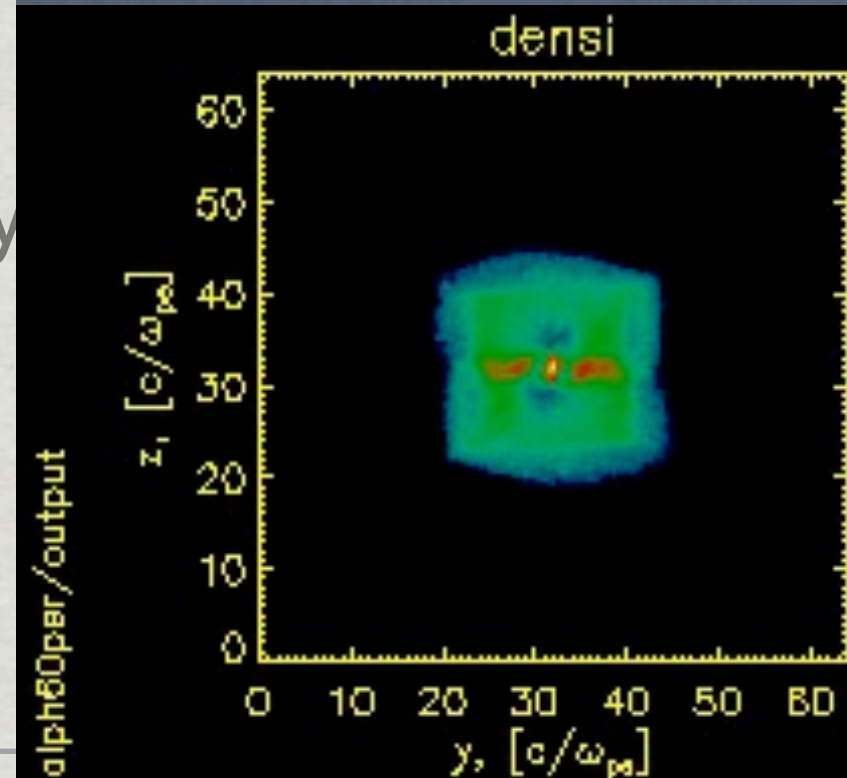
Full
density



<--- 7 R_{LC} --->

- * $E = E_{\text{plasma}} + E_{\text{Deutsch}}$;
 $B = B_{\text{plasma}} + B_{\text{Deutsch}}$

Ion
density



Step 1: Test electrons & ions become distributed in different locations -- no clear sheet

Recent advances:

Vacuum field + test particles

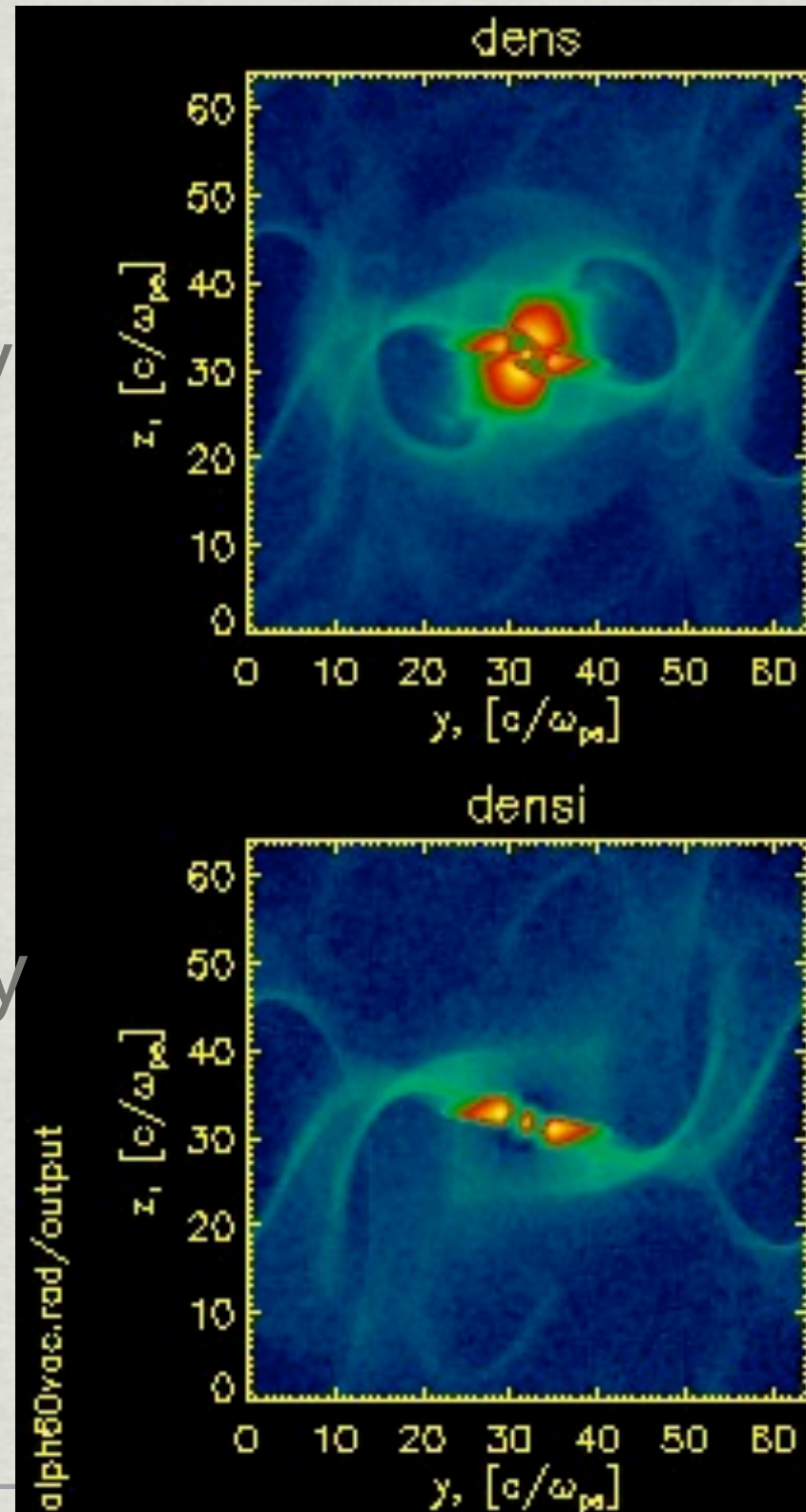
- * Full particle modeling with PIC (particle-in-cell) in 3D (AS in prep).
- * Idea: supply enough neutral plasma and “let it figure it out.”

Full
density

- * $E = E_{\text{plasma}} + E_{\text{Deutsch}};$
 $B = B_{\text{plasma}} + B_{\text{Deutsch}}$

Ion
density

Step 1: Test electrons & ions become distributed in different locations -- no clear sheet



Recent advances:

Vacuum field + plasma

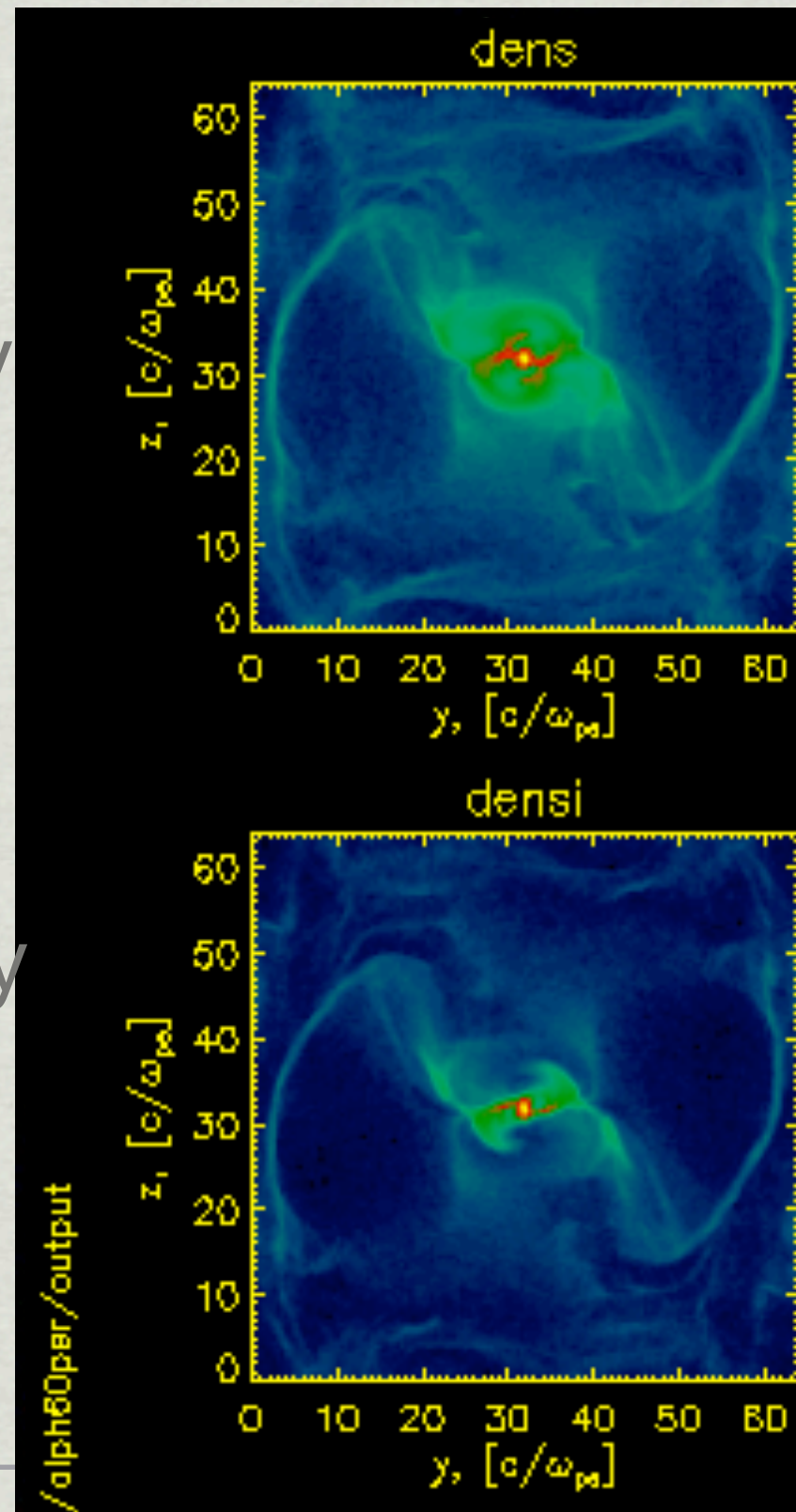
- * Full particle modeling with PIC (particle-in-cell) in 3D.

Full density

- * Turn on plasma back reaction:

Ion density

Electrons & ions now are together and form a current sheet. Approaches force-free.

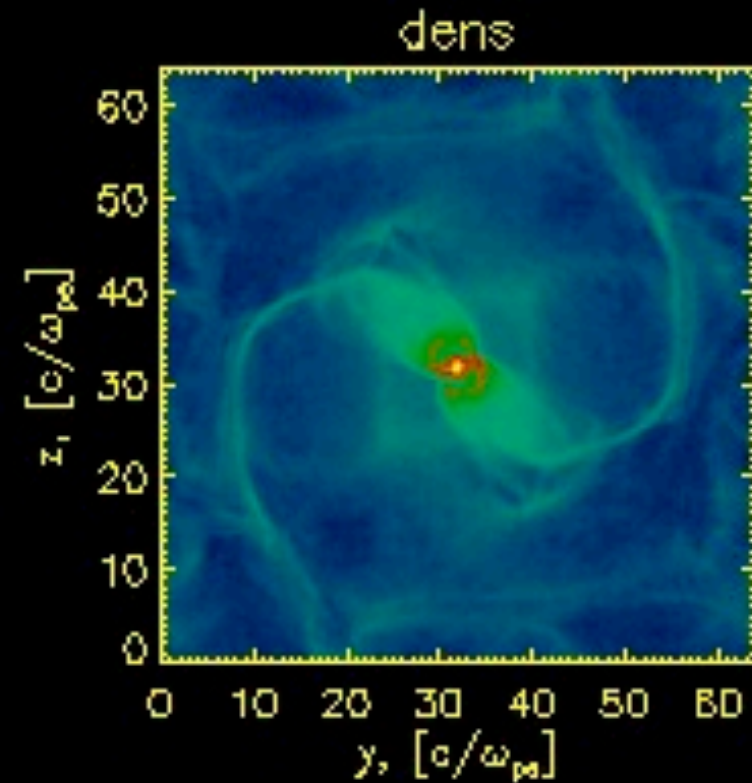


Recent advances:

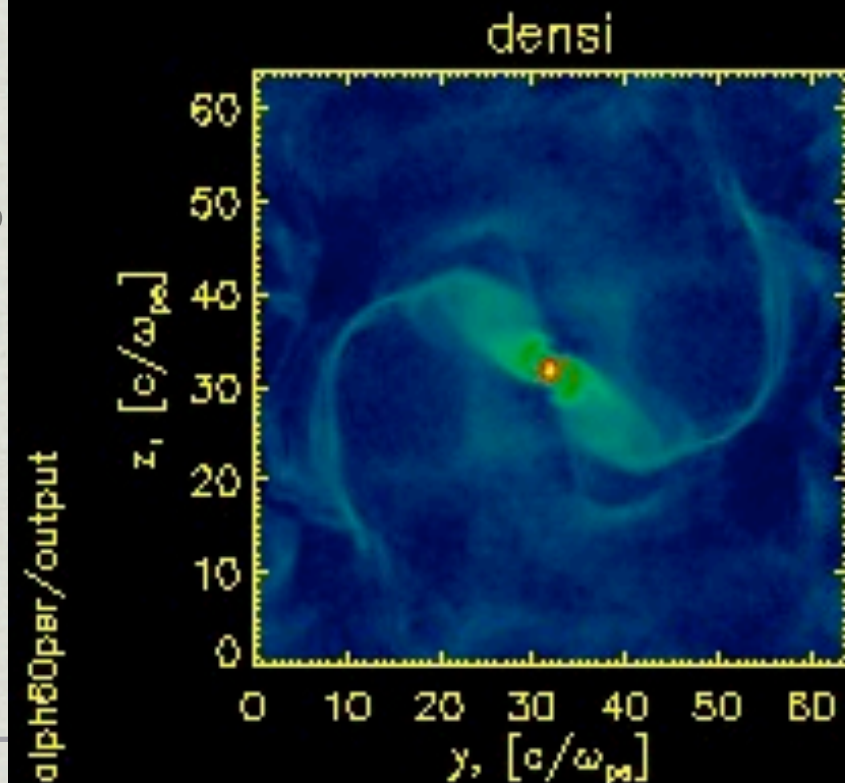
Vacuum field + plasma

- * In principle, such a solution has everything: acceleration, plasma velocity, reconnection.
- * In practice the resolution will always be limited.
- * Also, the aligned rotator part depends on good BCs on the star.

Full
dens



Ion
dens

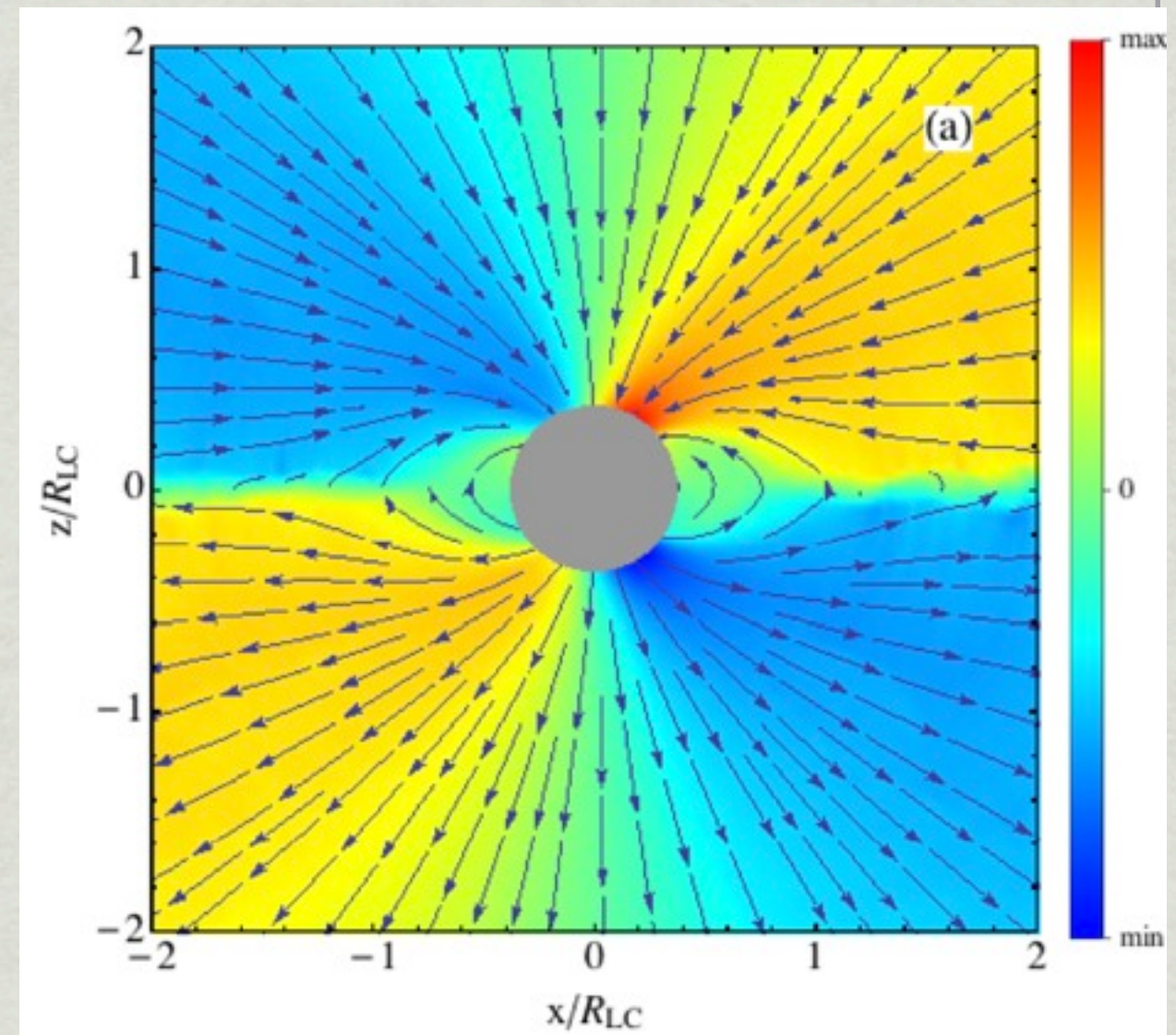


alpha50per/output

Ab-initio pulsars

Philippov + AS 2014

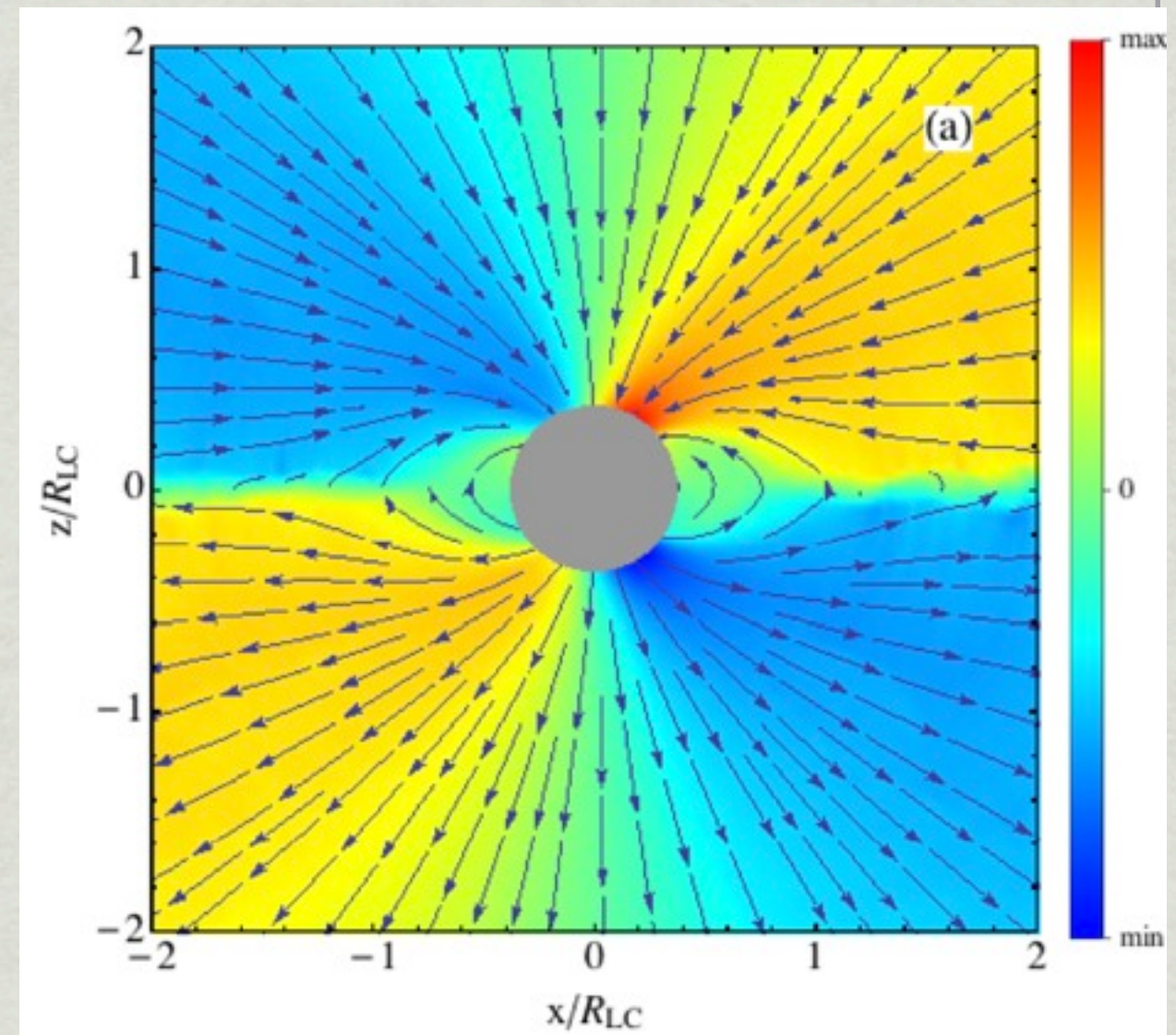
- * BC on the star: space-charge limited flow, particle escape, good spherical conductor (challenge on Cartesian grid) .
- * Use spherical polar coordinates (Cerutti, Belyaev) -- 2D limitation
- * We used “plasma sphere” BC (see Philippov’s talk)



Ab-initio pulsars

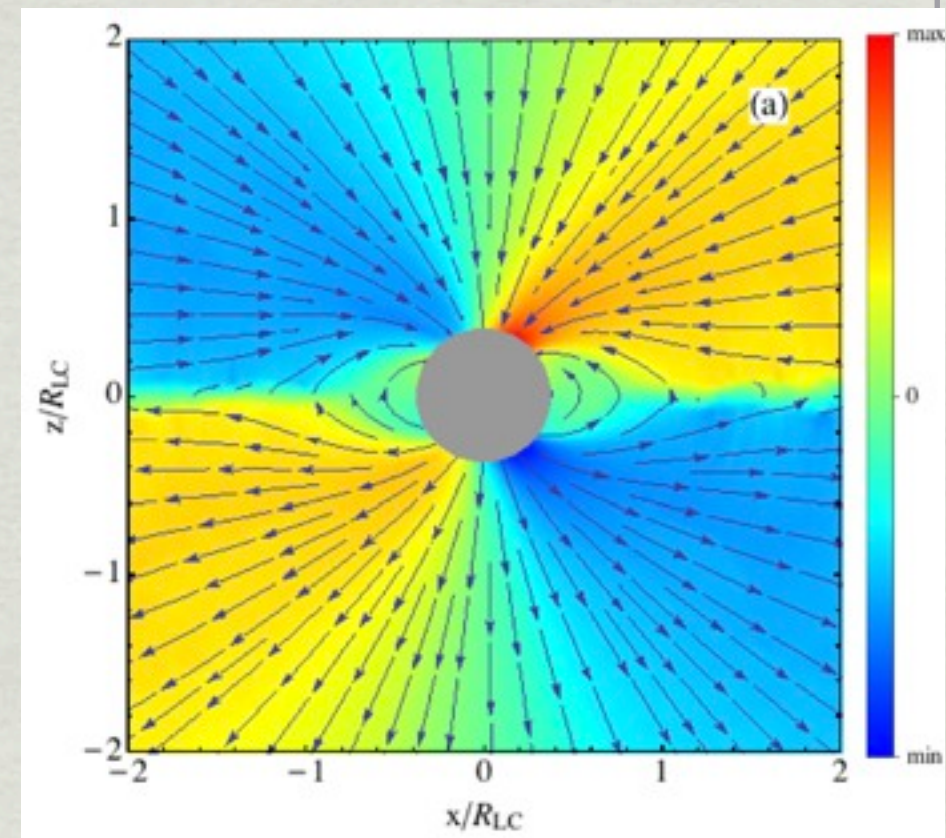
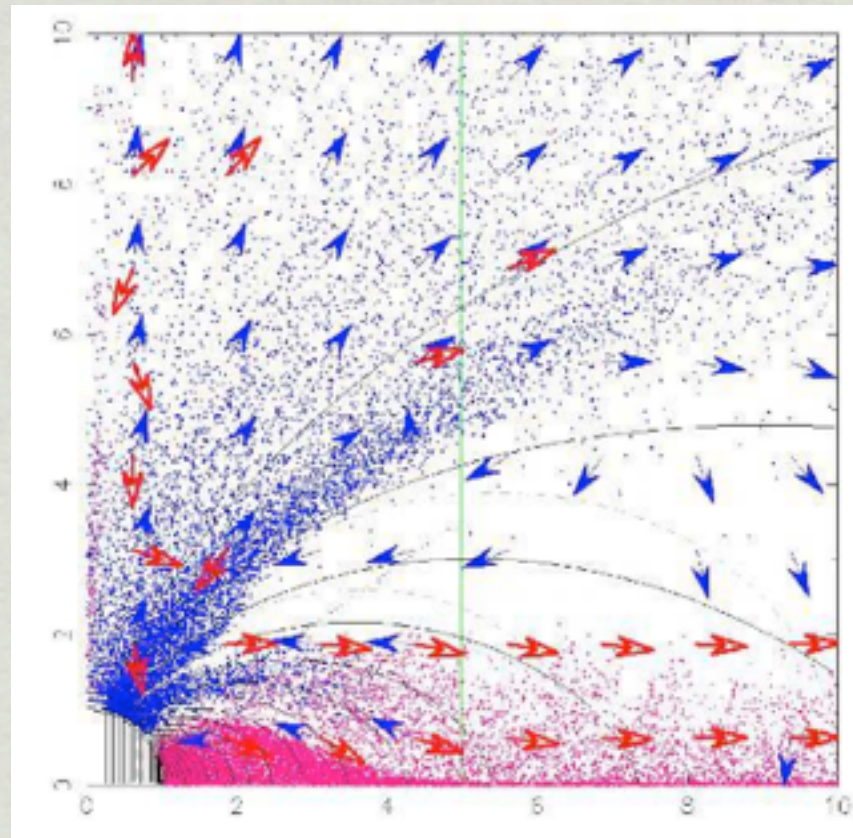
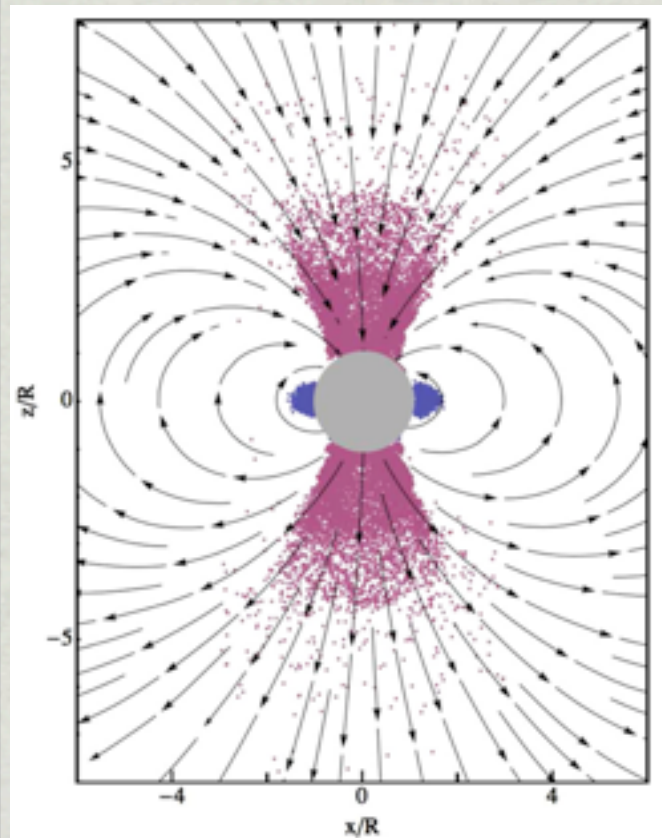
Philippov + AS 2014

- ✱ Abundant pair plasma with PIC reproduces force-free
- ✱ Small dissipation ($\sim 10\%$ in current sheet)
- ✱ Particle acceleration mainly in the sheet
- ✱ Drift-kink instability of the sheet



Ab-initio pulsars

- ✱ There may be other solutions depending on plasma supply; we are experimenting with pair formation prescriptions.



There is a class of solutions with $E > B$ and accelerated particles (e.g. Gruzinov's talk). They must be low-multiplicity states, that may not produce abundant pulsar wind as needed by observations.

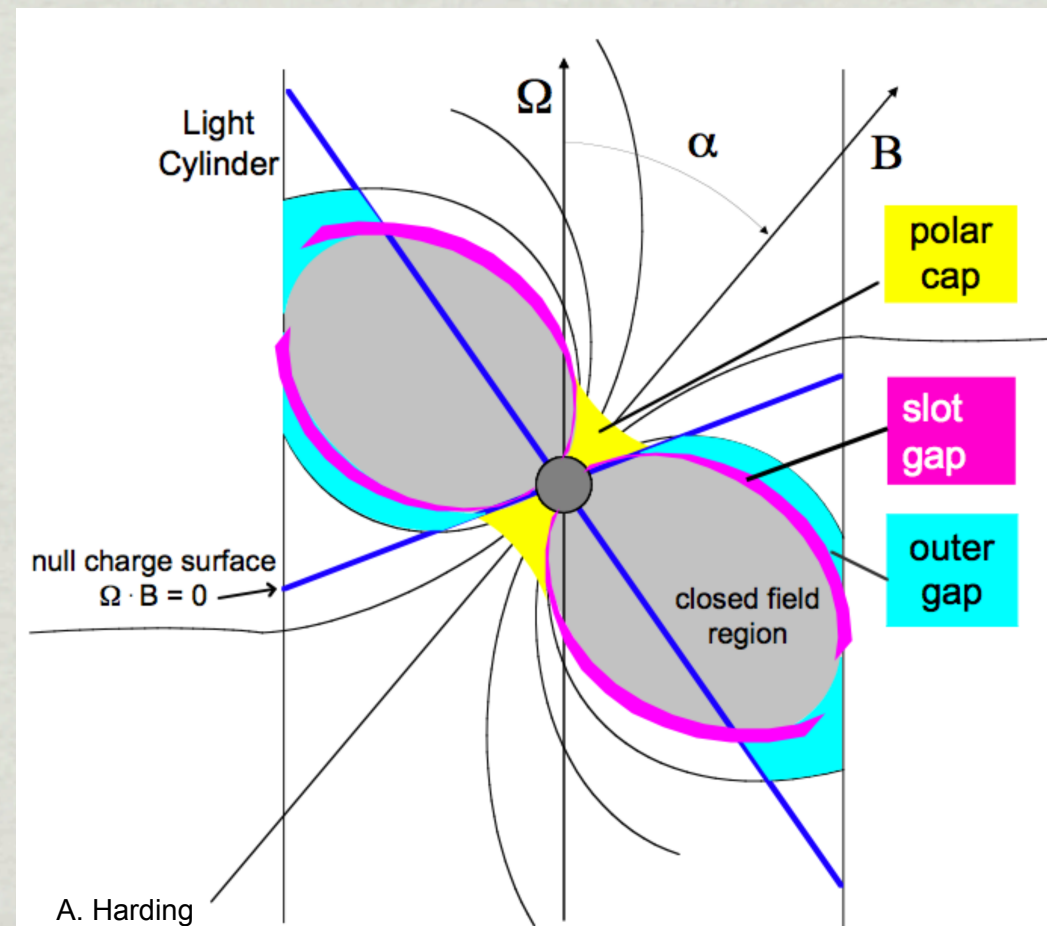
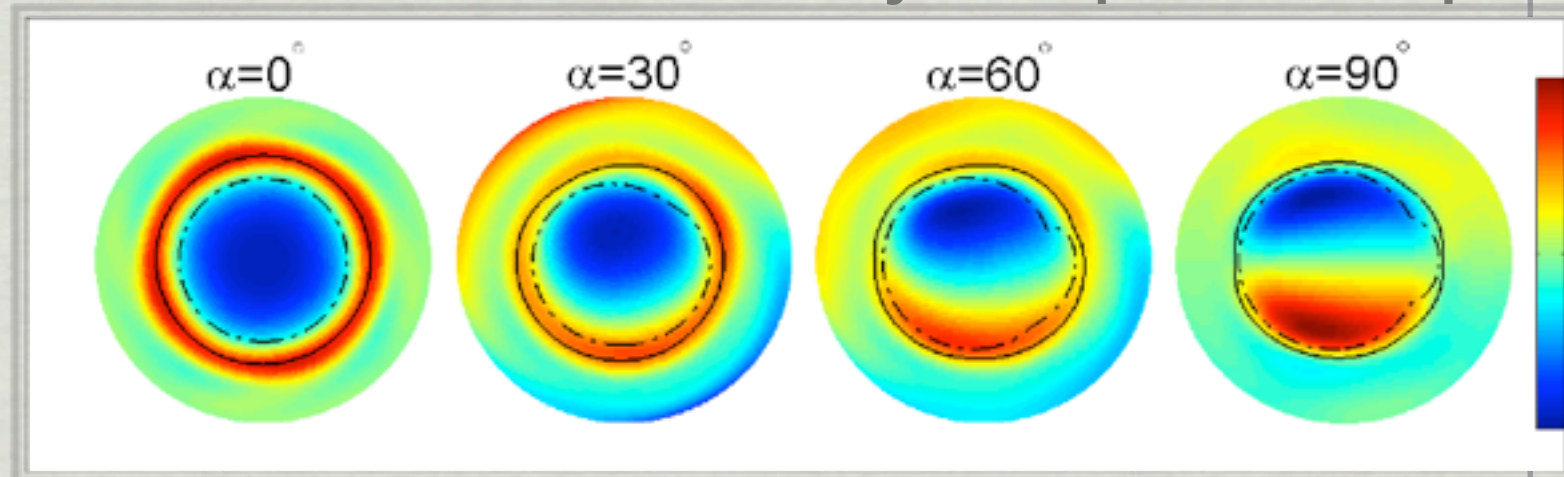
Where does emission come from?

- ✳ With force-free shape of the magnetosphere at hand, emission physics can be studied again.

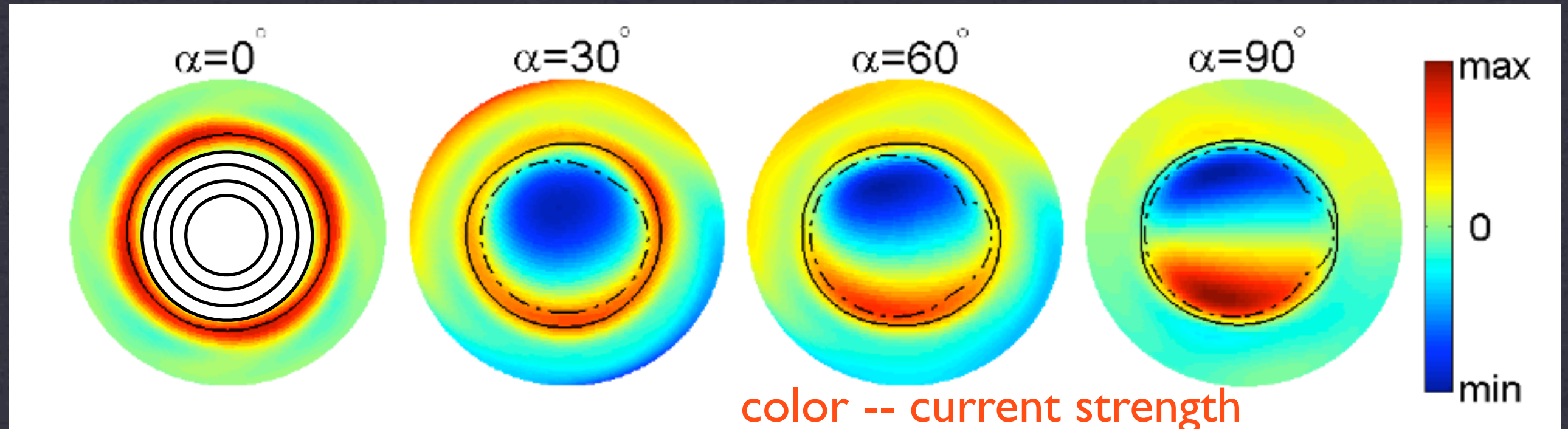
- ✳ Most progress so far is on the geometry of the emission so far (gamma-rays).

- ✳ Does the emission come from gap regions? Inside or outside LC?

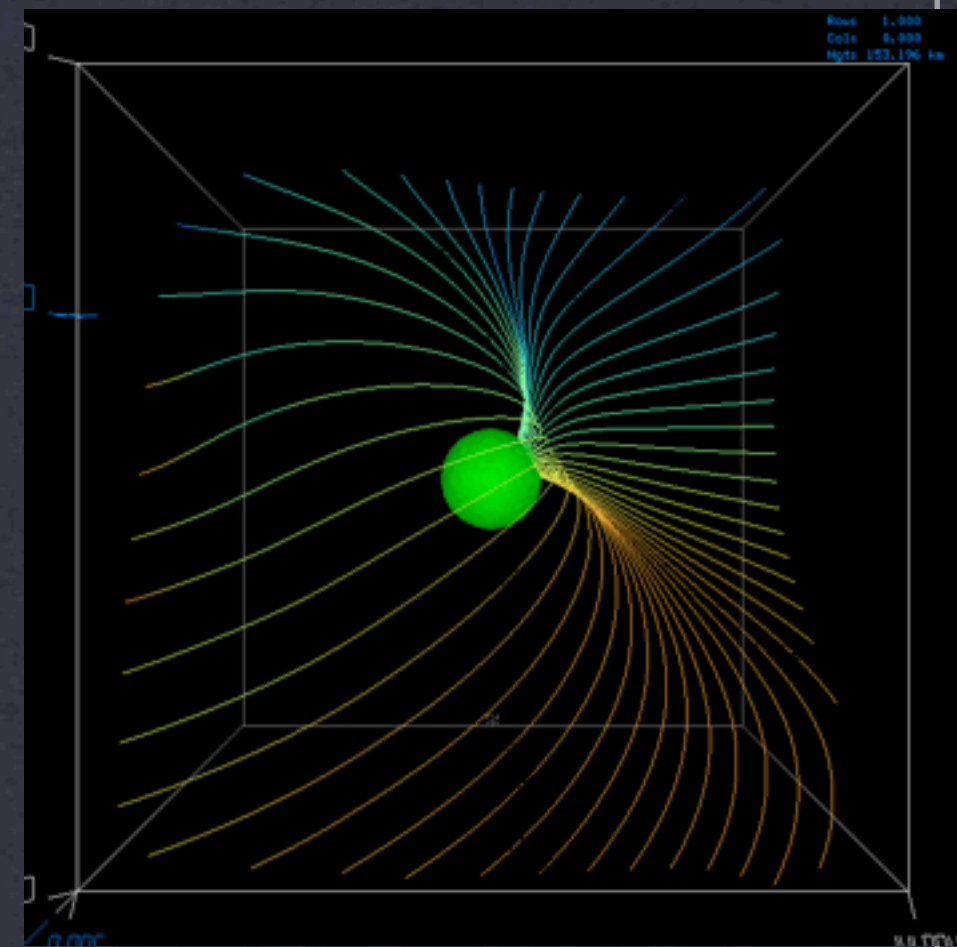
Current density on polar cap



Where does emission come from?



- Select flux tubes that map into rings on the polar caps. The rings are congruent to the edge of the polar cap.
- While ad-hoc, the point is to study the geometry of the possible emission zone.
- Emission is along field lines, with aberration and time delay added

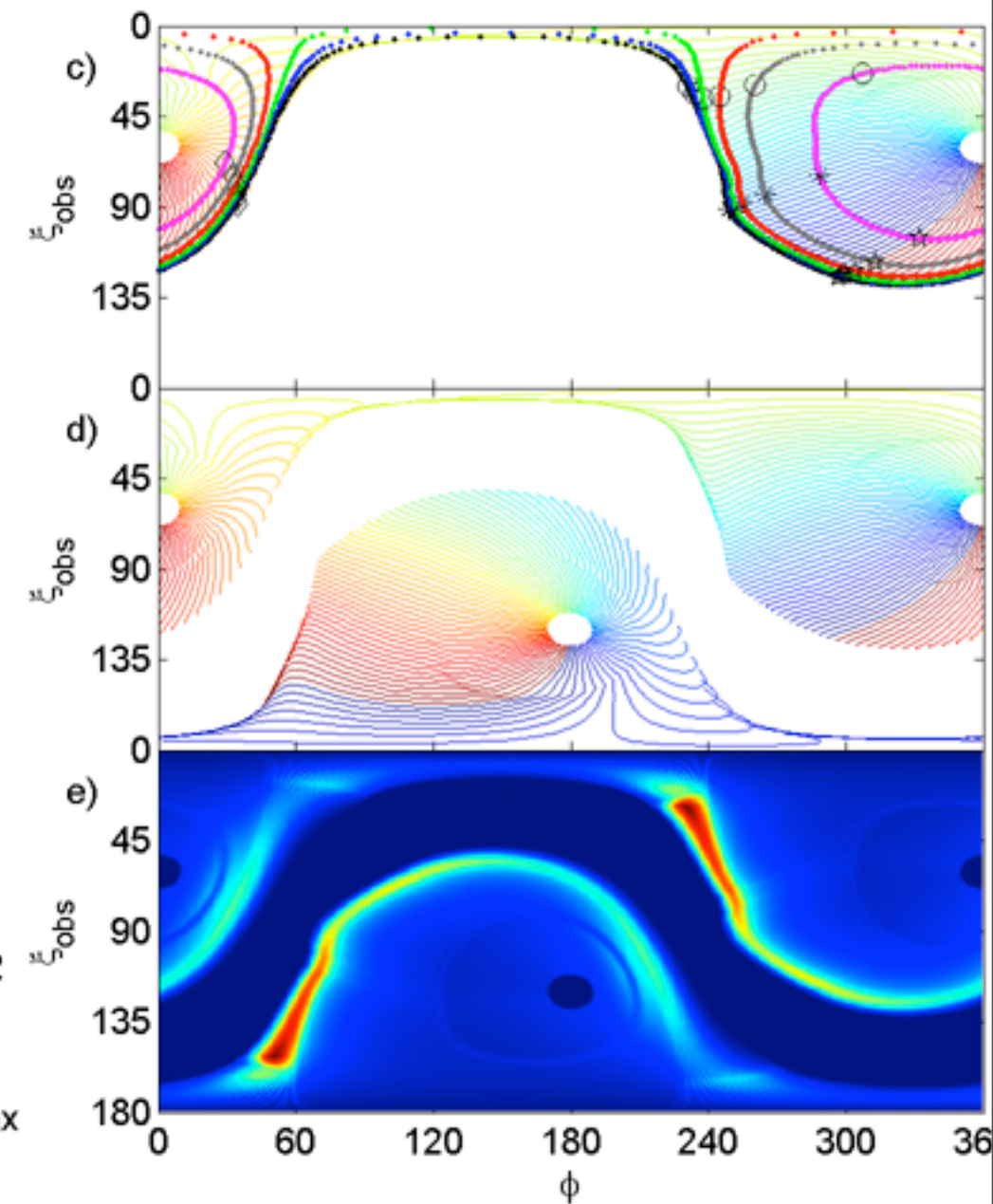
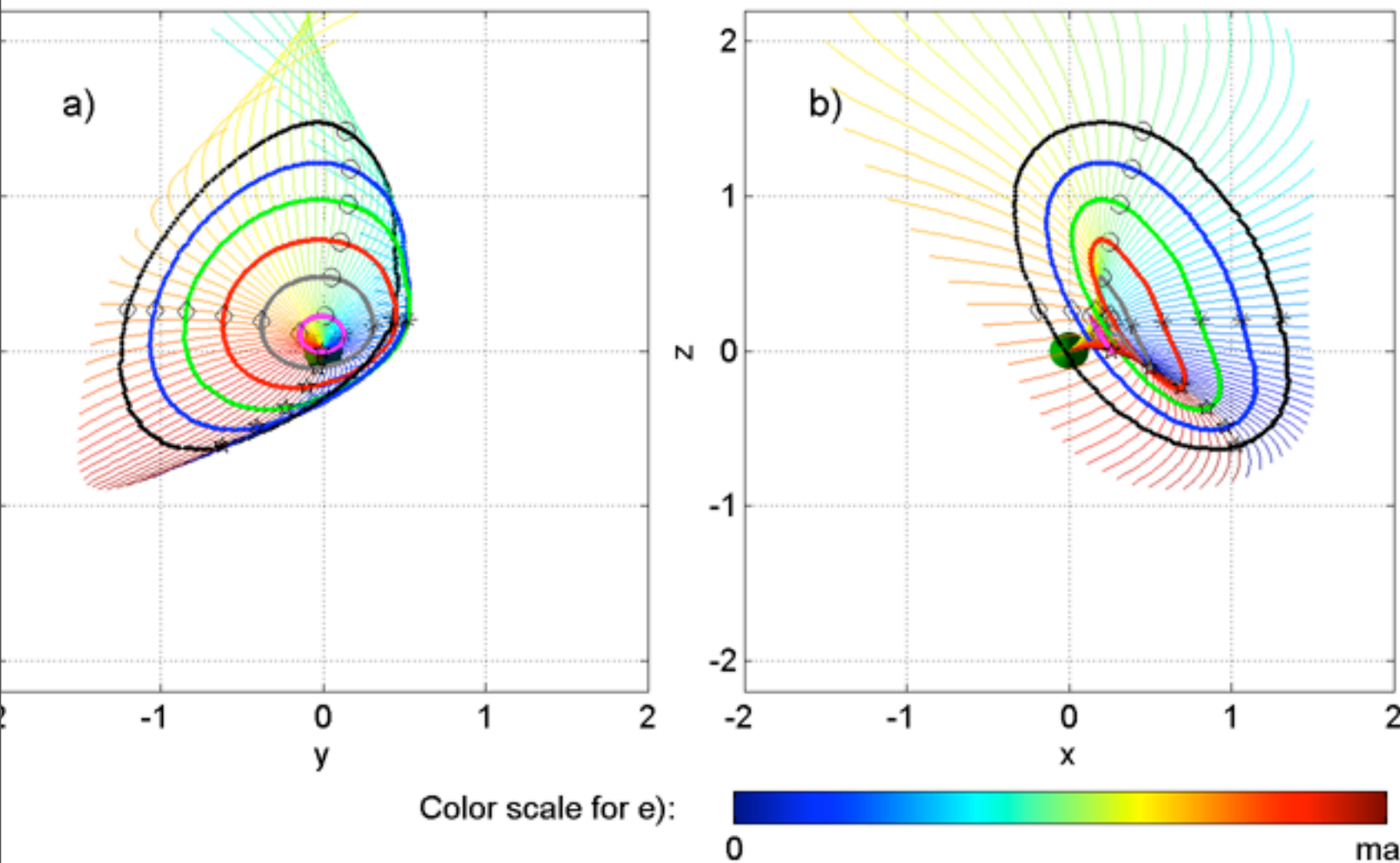


Emission from different flux tubes

Bai & A. S. 2010

AG-F60-80-90

a)-b): spatial plot; e): sky map intensity
c)-d): projection to sky map;



Emission from two poles merges on some flux tubes: what's special about them?

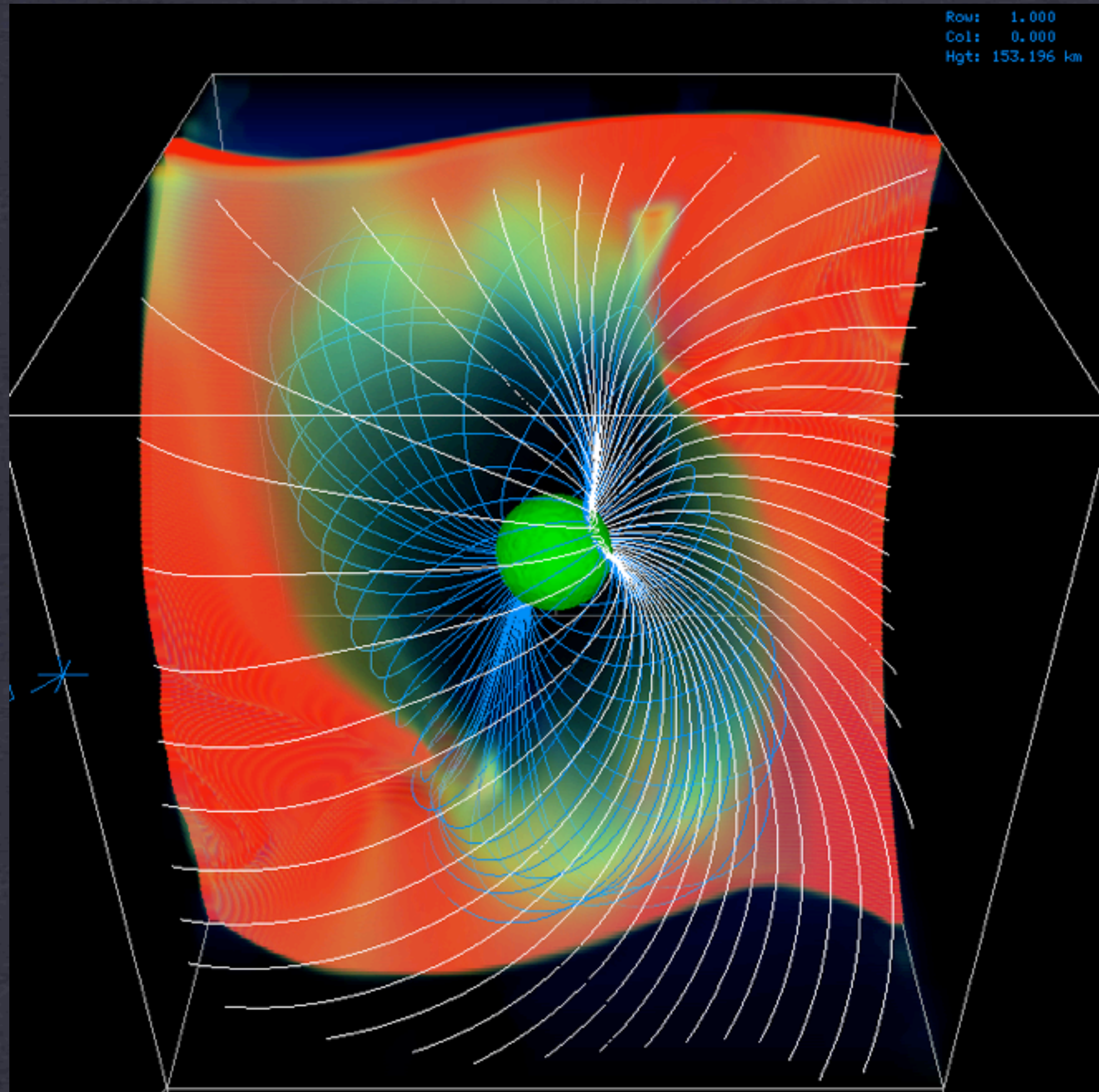
Association with the current sheet

Color -> current

Field lines that produce best force-free caustics seem to “hug” the current sheet at and beyond the LC.

Significant fraction of emission comes from beyond the light cylinder.

Best place to put a resistor in the circuit!



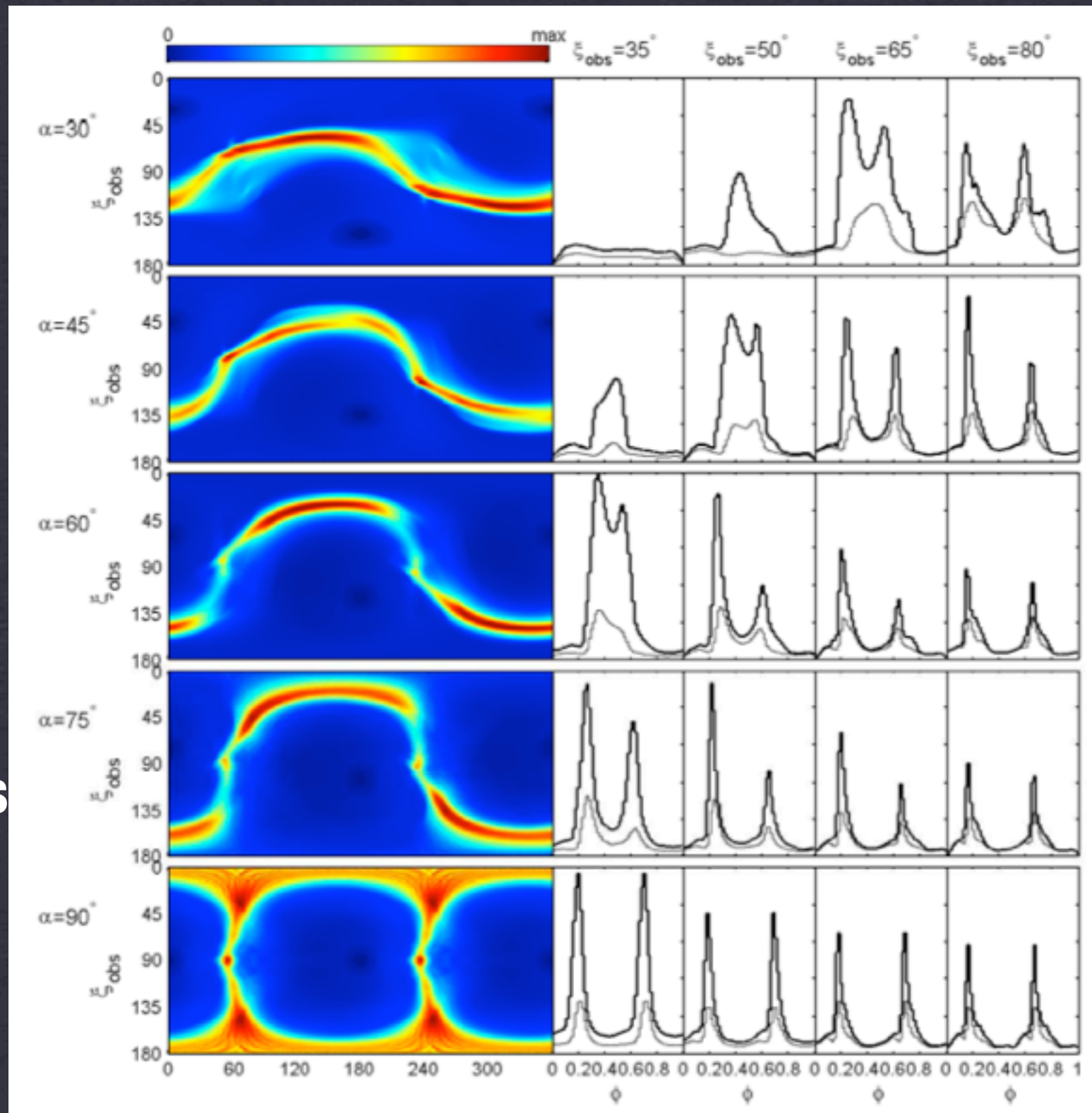
Force-free light curves

Viewing angle

**“Separatrix Layer”
model is a real
contender. It’s not
just outer gap vs
slot gap anymore!**

Inclination
angle

**Most of the
emission in FF
model accumulates
beyond $0.9 R_{lc}$**



Double peak profiles very common.

Bai & AS, 2010

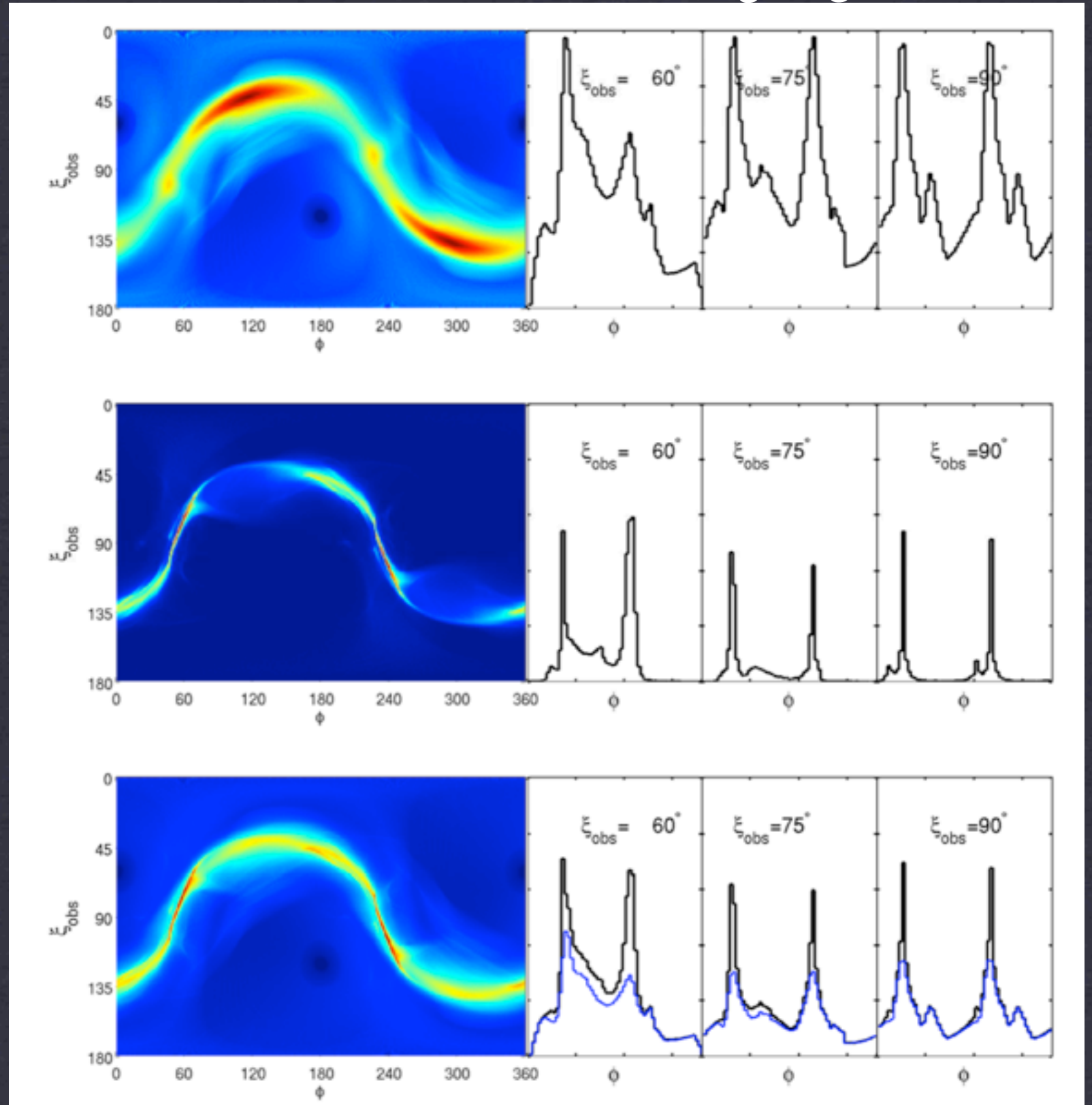
Resistive Force-free light curves

Beaming: along interpolated B field into the sheet.
Results in radial beaming.

Combine emission from current layer ($< R_{lc}$) for bridge emission with current sheet ($> R_{lc}$) for peaks

Inclination angle

Viewing angle



Li, AS, Tchekhovskoy 2014

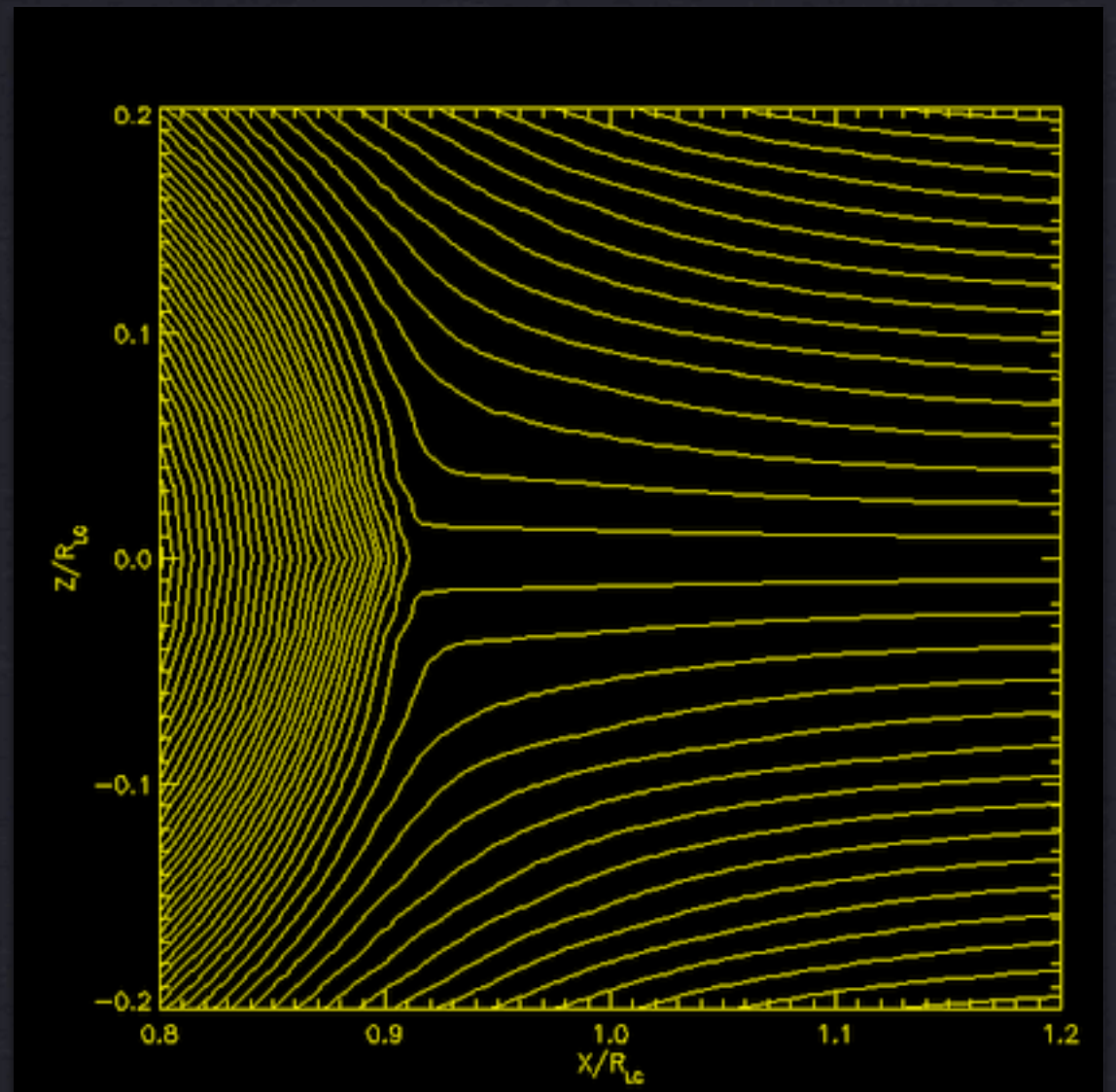
Source of emission

- ✧ Emission is geometrically associated with the current sheet
- ✧ What is the acceleration and radiation mechanism in current sheet?

Most likely culprit -- relativistic reconnection. This is different from conventional picture of accelerating gaps starved of plasma and curvature emission

- ✧ Boosted synchrotron from heated plasma can work

Reconnection controls magnetospheric shape!

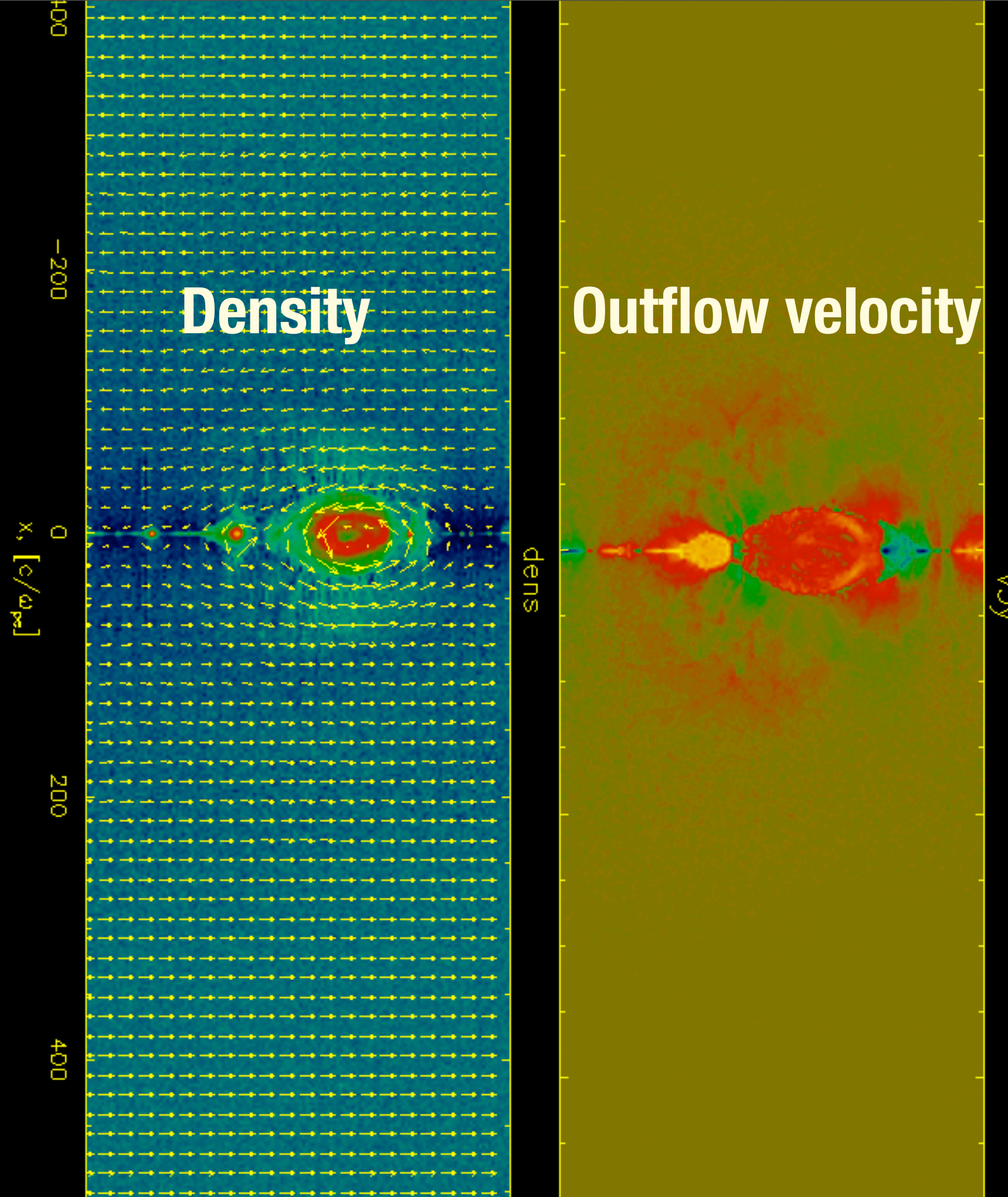


**Better ideas of
flow direction in the
current sheet needed.**

**In PIC simulations
get outflows near
 $\sqrt{\sigma}$.**

Minijets?

**Since beaming along
extrapolated B field in the
current sheet makes
double peaks, it's a
contender**



Why reconnection makes sense

- ✱ Conditions in the sheet can be obtained from:

Pressure balance

$$B_0^2/8\pi = 2nT$$

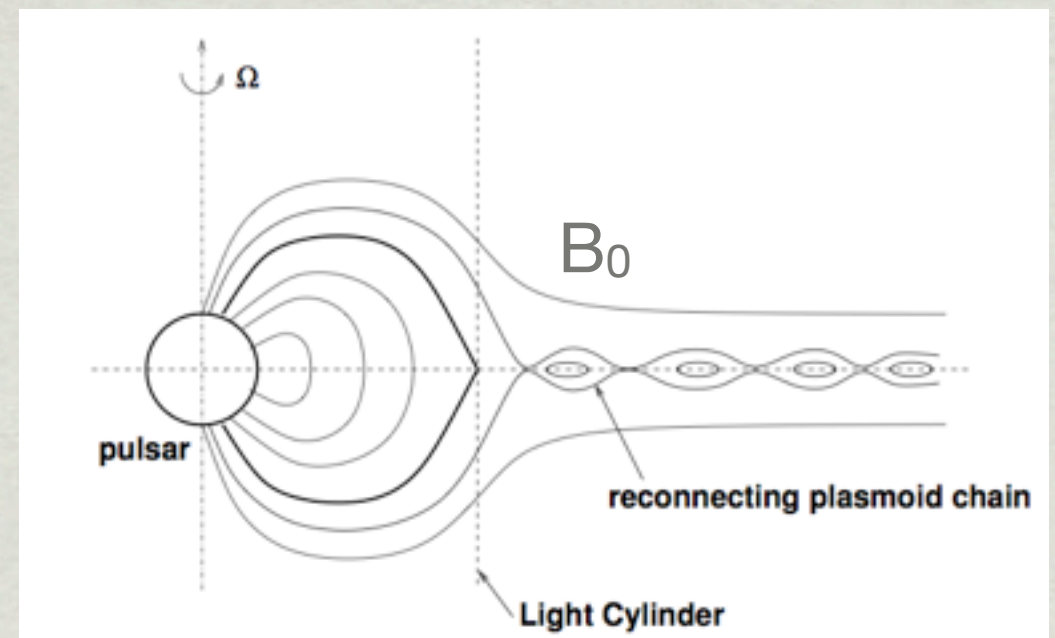
Strong synchr. cooling:

$$S_{in} = (c/4\pi) E_z B_0 \sim Q_{rad} \sim \delta (2n) P_{sync}(T)$$

Ampere's law:

$$j_z = 2ne v_{dr} = 2nec \beta_{dr} \sim (c/4\pi) B_0/\delta$$

Temperature at 10GeV comoving --> 160MeV synch radiation --> GeV pulsed emission in the lab boosted by bulk gamma of ~10. IC gives VHE.

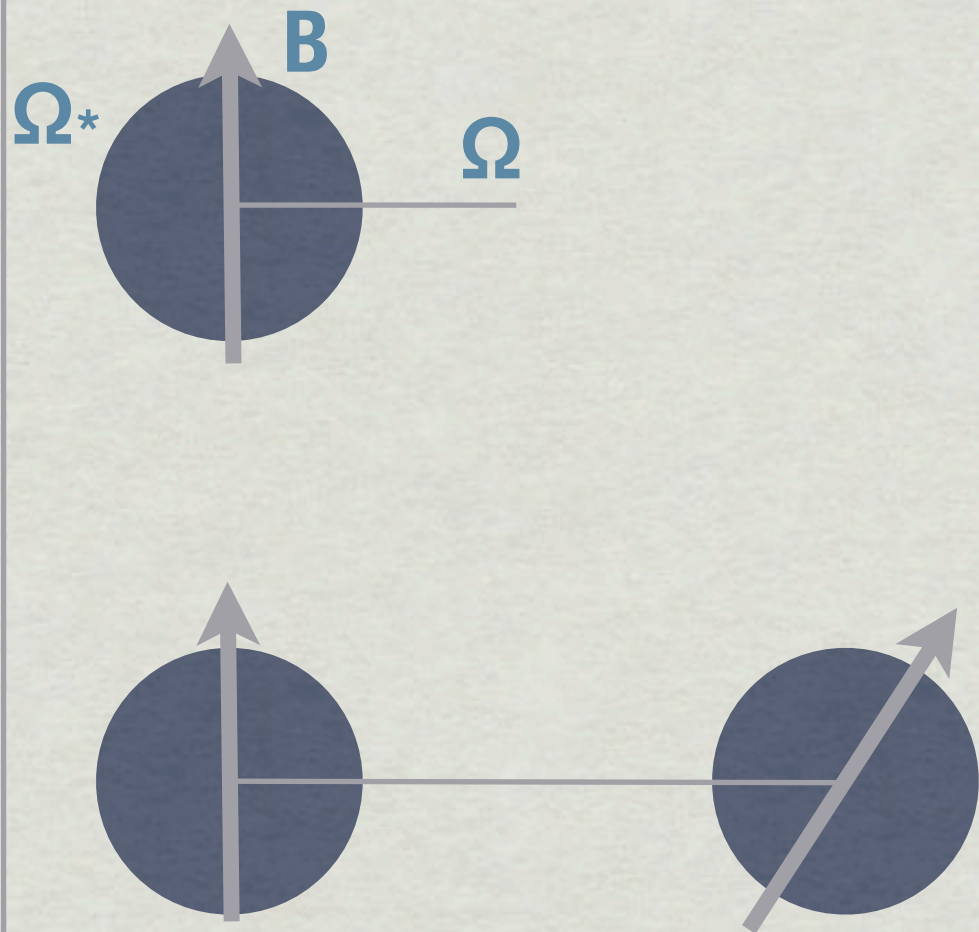


Uzdensky & AS 2014 (also, Lyubarsky 96, Petri 12, Arca 12)

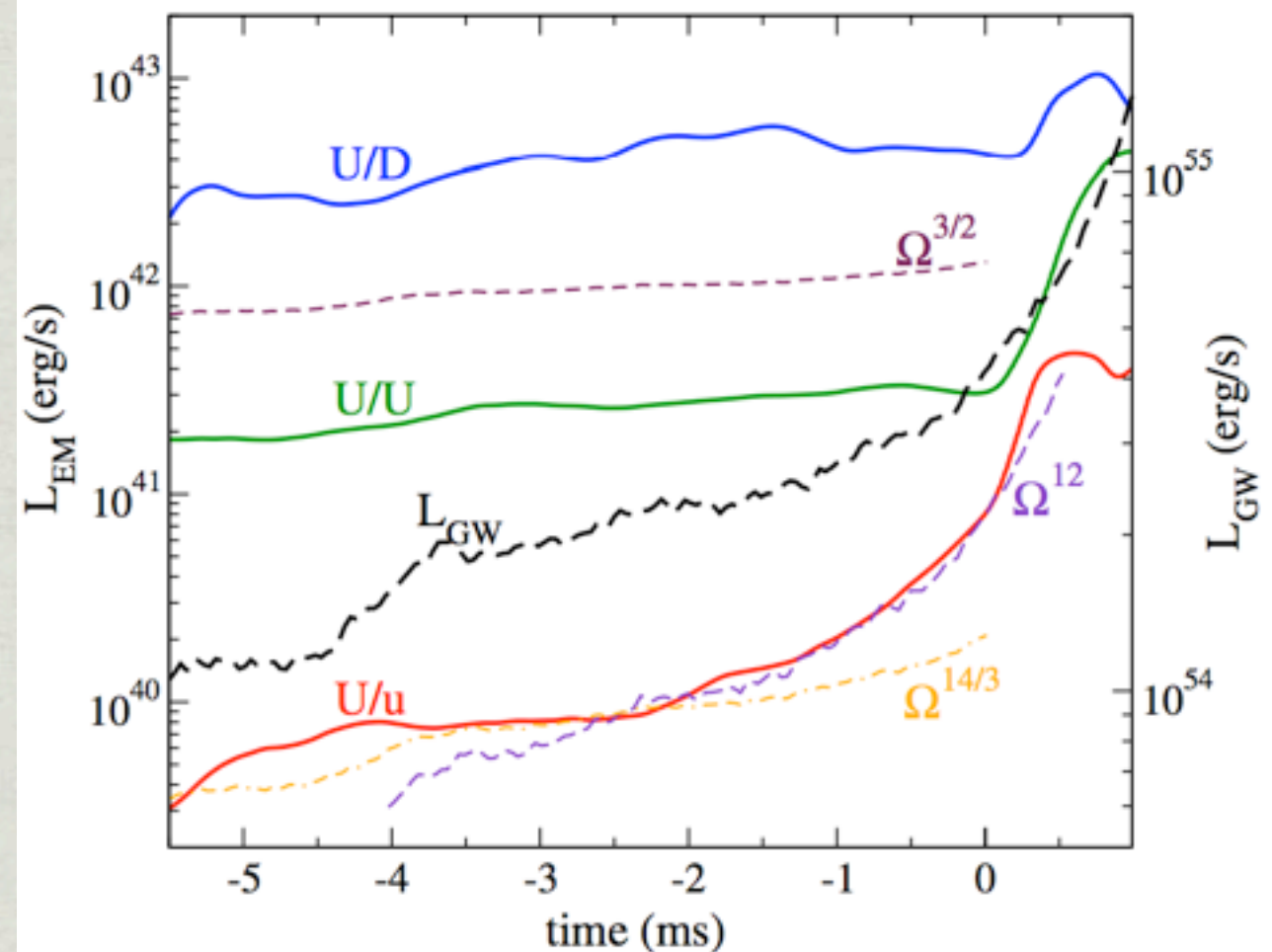
Temperature, density and thickness depend on B at LC. $\gamma_T = T/m_e c^2 \sim [\beta_{dr} \beta_{rec} 8\pi e/\sigma_T B_0]^{1/2} \sim (\beta_{dr} \beta_{rec})^{1/2} 4 \times 10^4$

Other EMF: binary inspiral

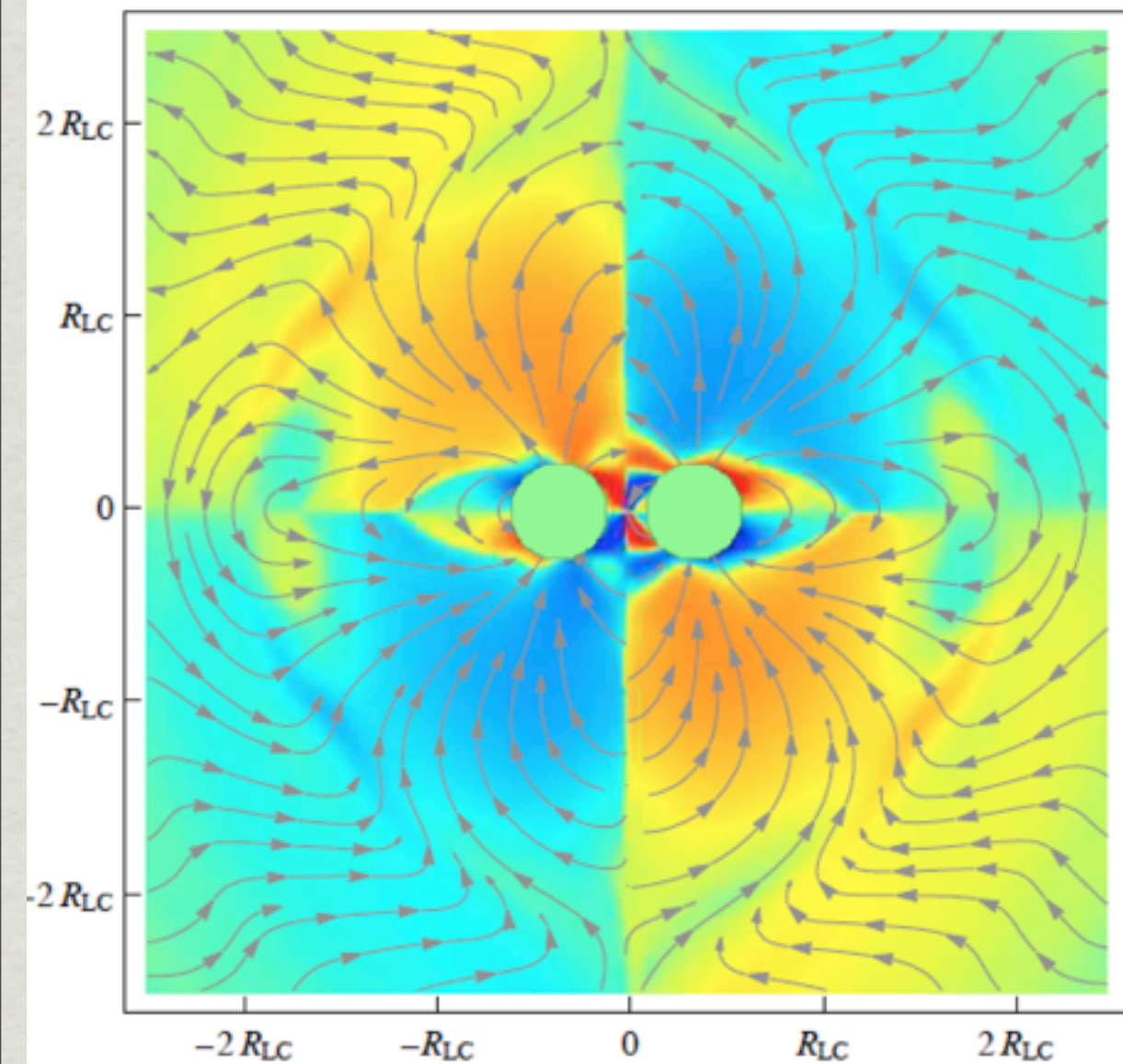
- ✱ Inspired by prospects for precursor emission for LIGO, we studied spin-down of single and binary NS in orbit (Li + AS 2014). Cf. work by Palenzuela et al 2013-14.
- ✱ How does the spin-down depend on orbital frequency and spin and orientation of NSs?



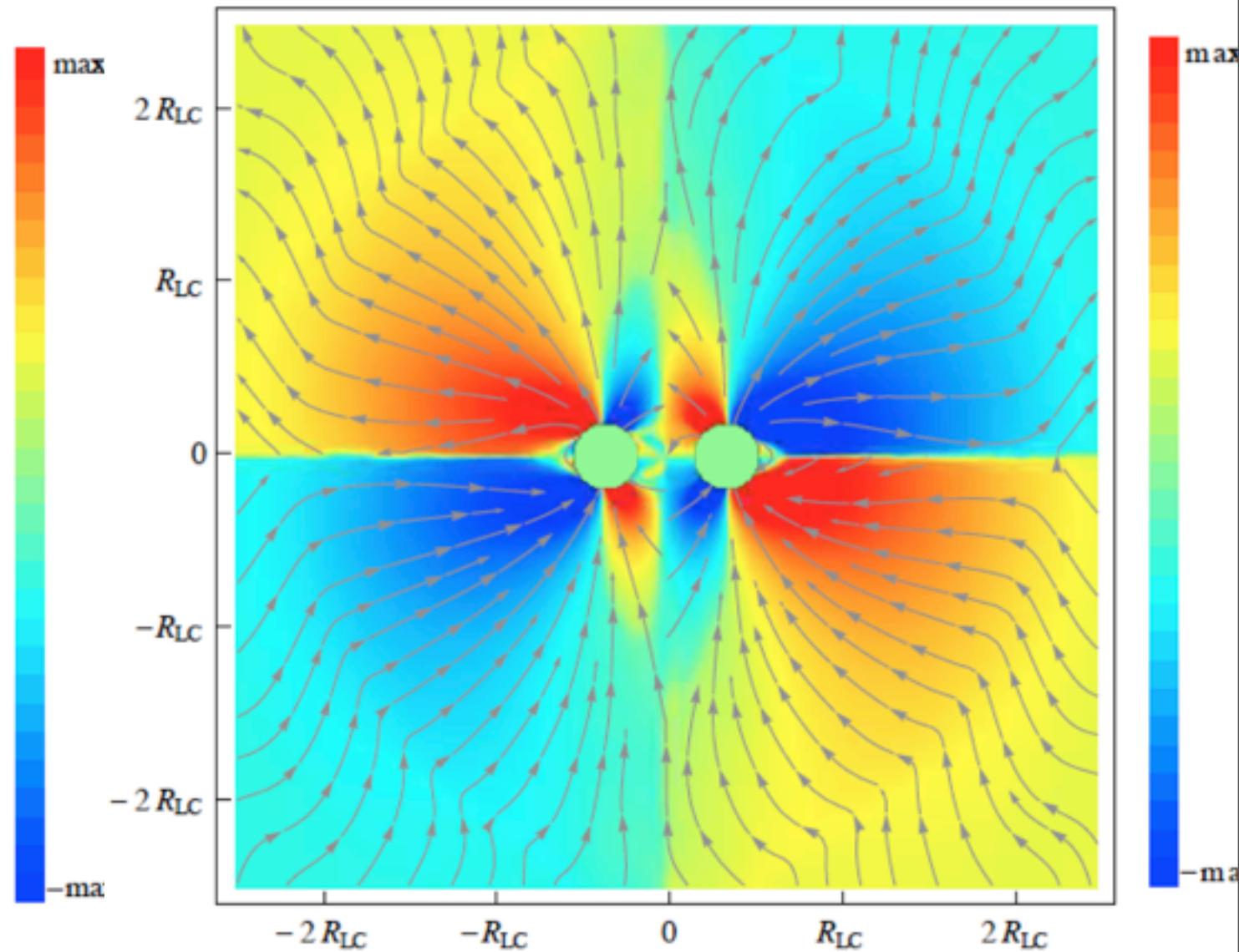
Palenzuela



Binaries: field patterns

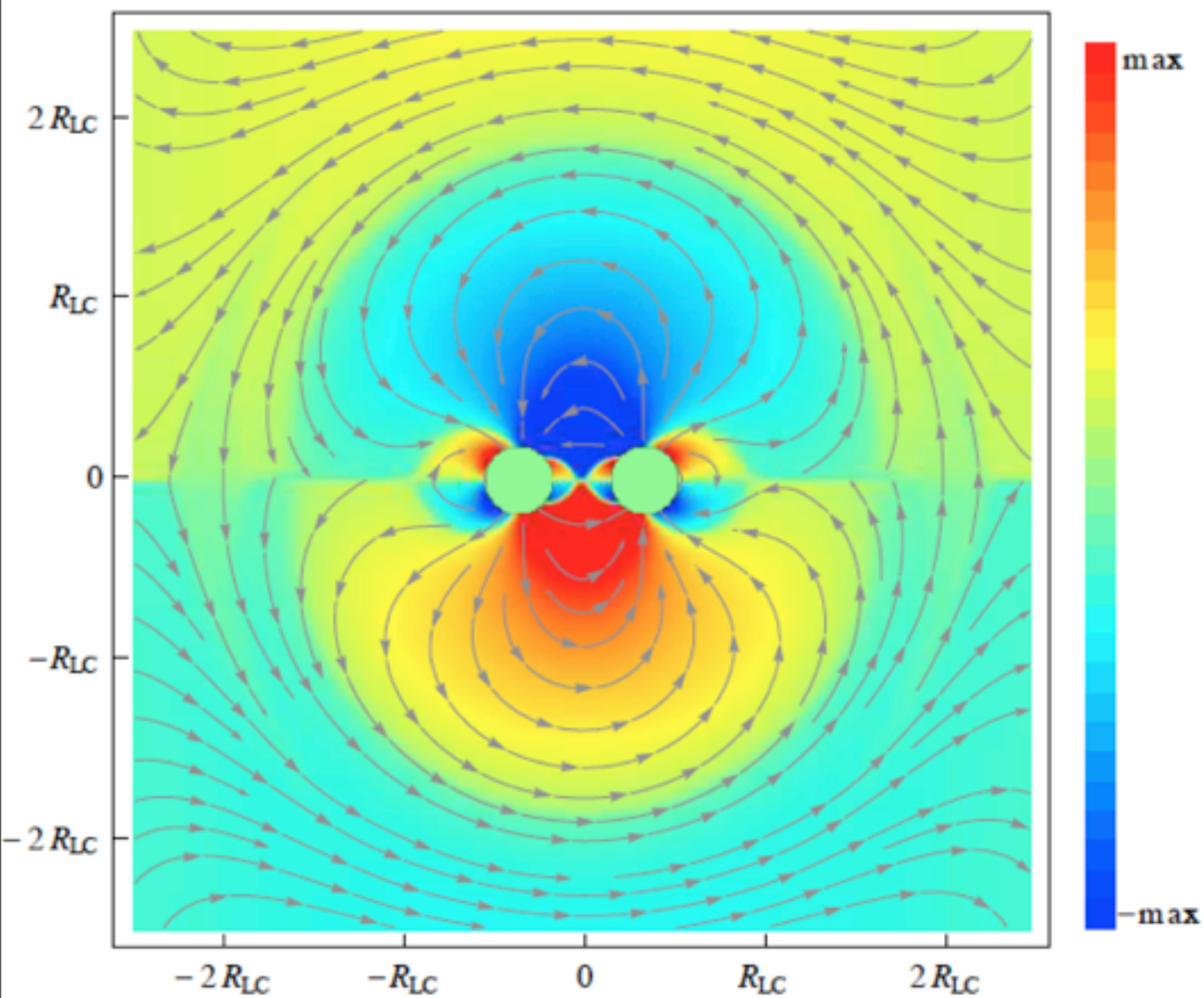


U/U, no NS spin

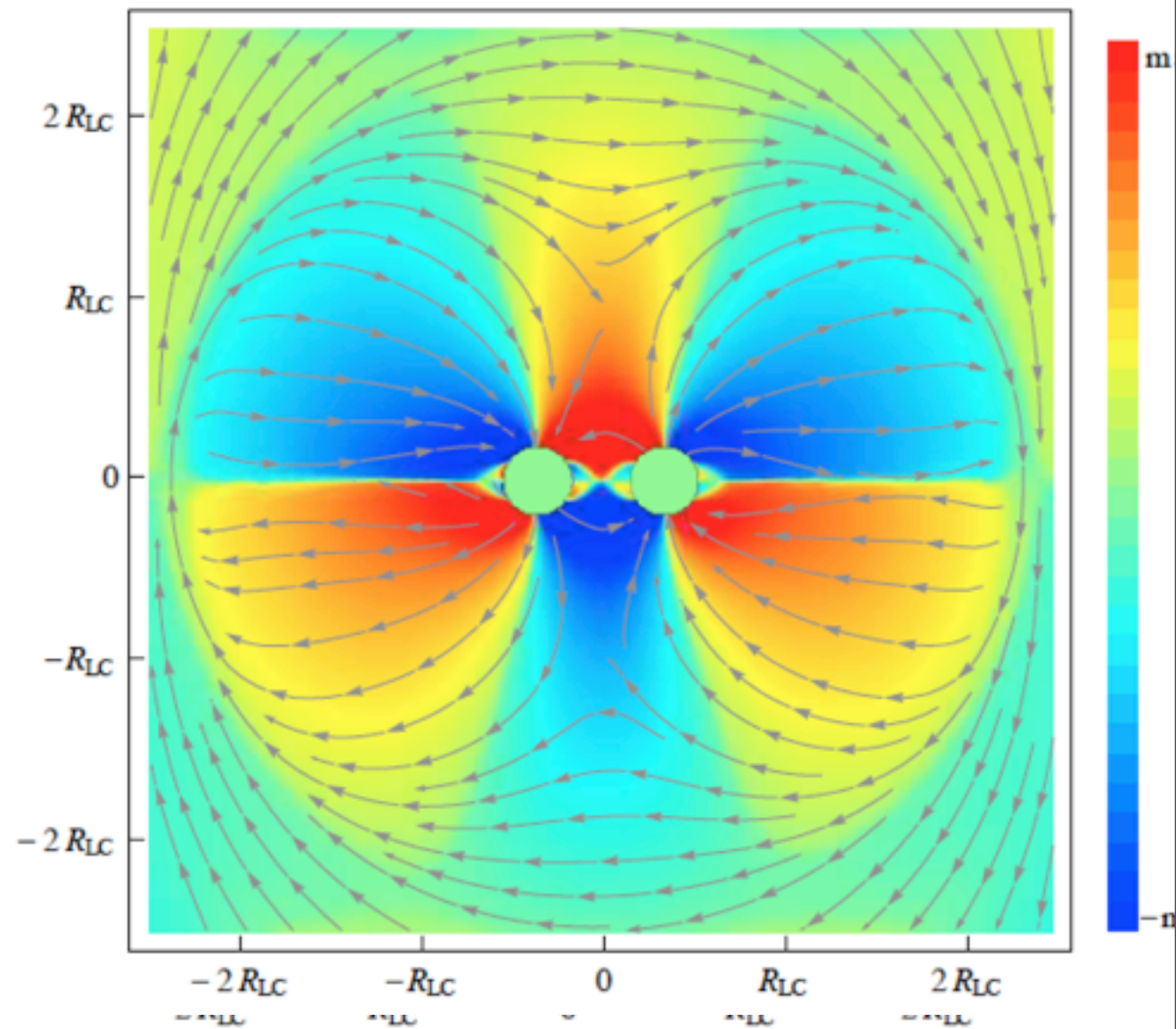


U/U, with spin

Binaries: field patterns

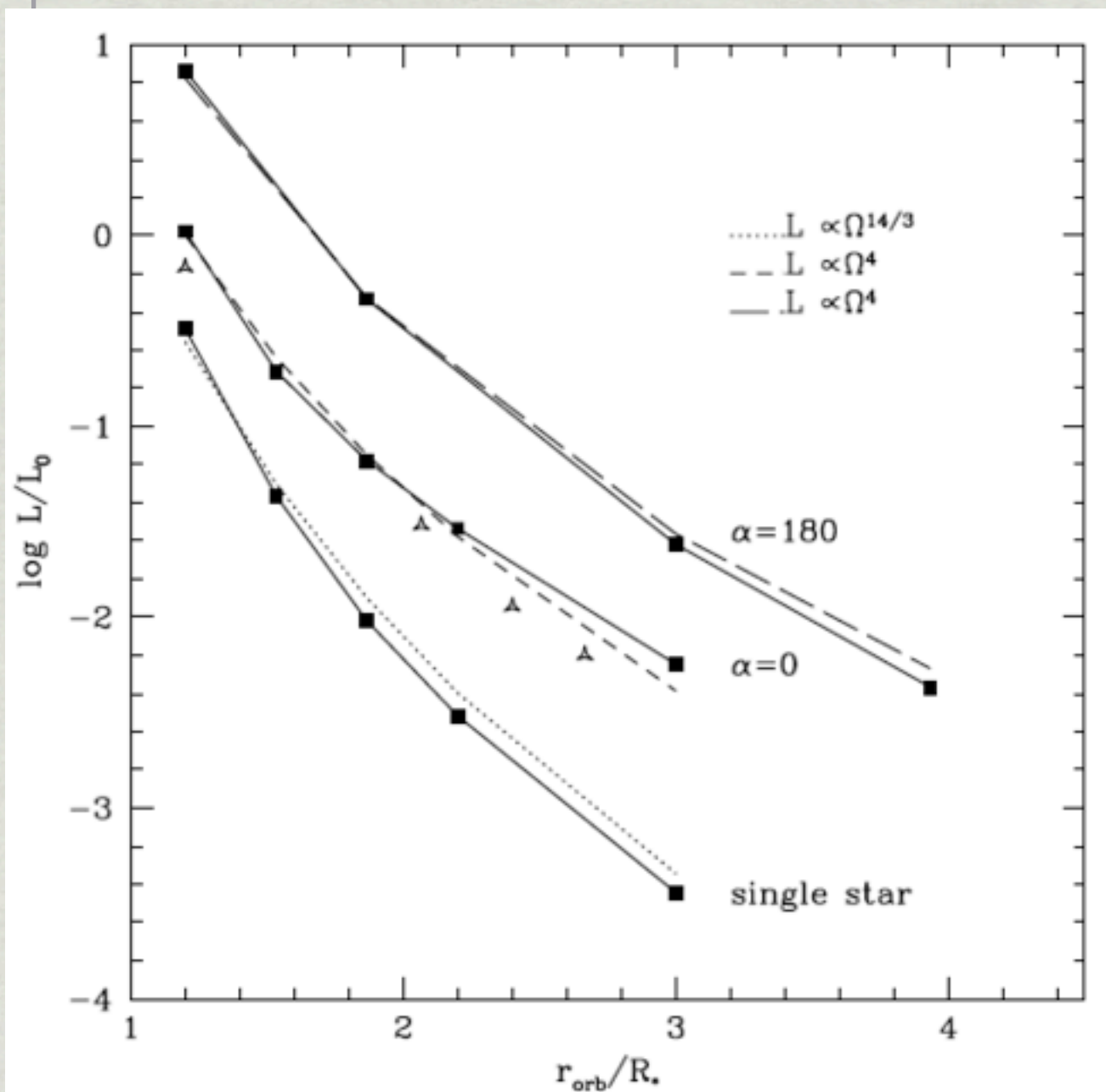


U/D, no NS spin

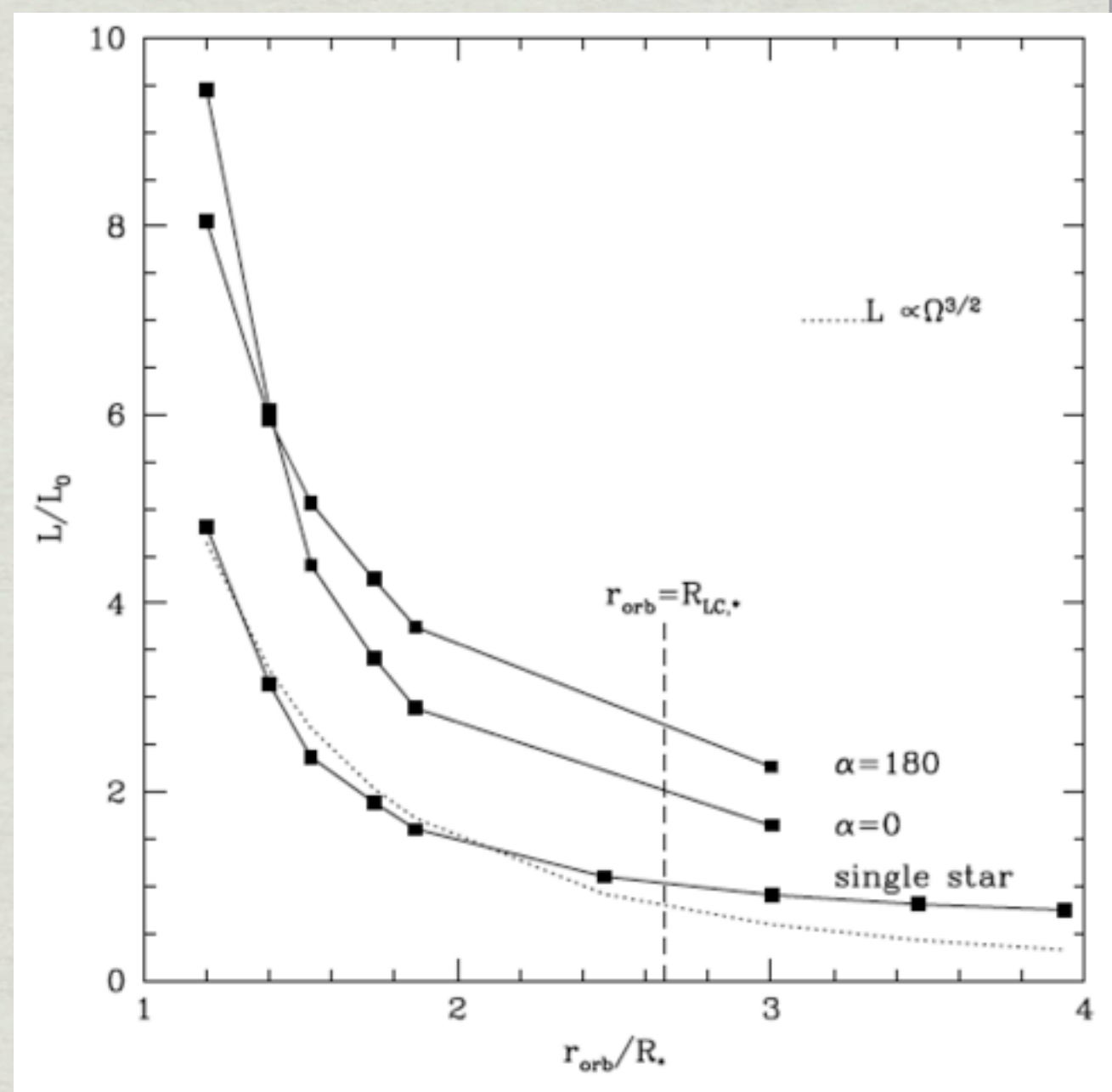


U/D, with spin

Binaries: EM luminosity

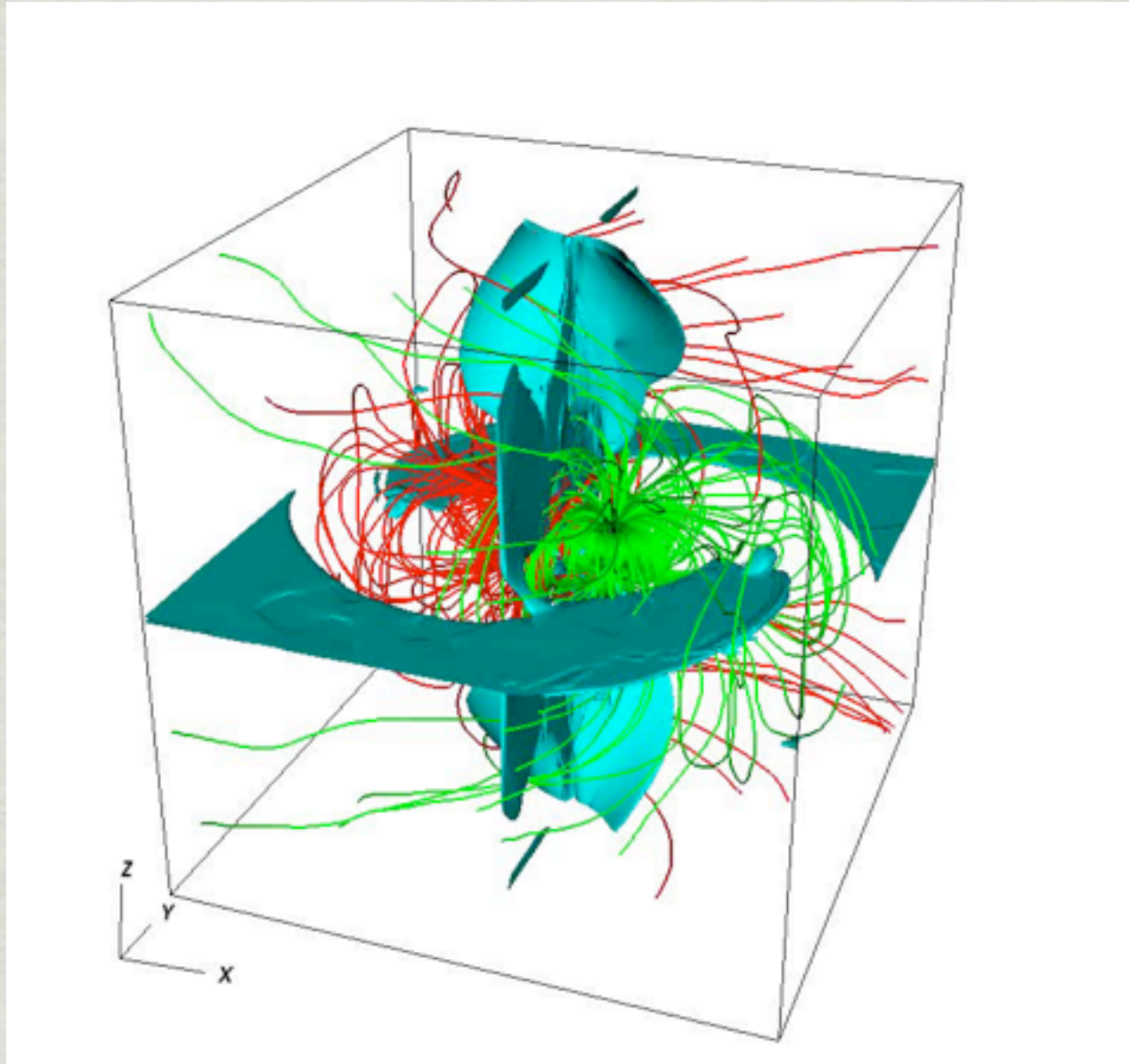


Non-spinning

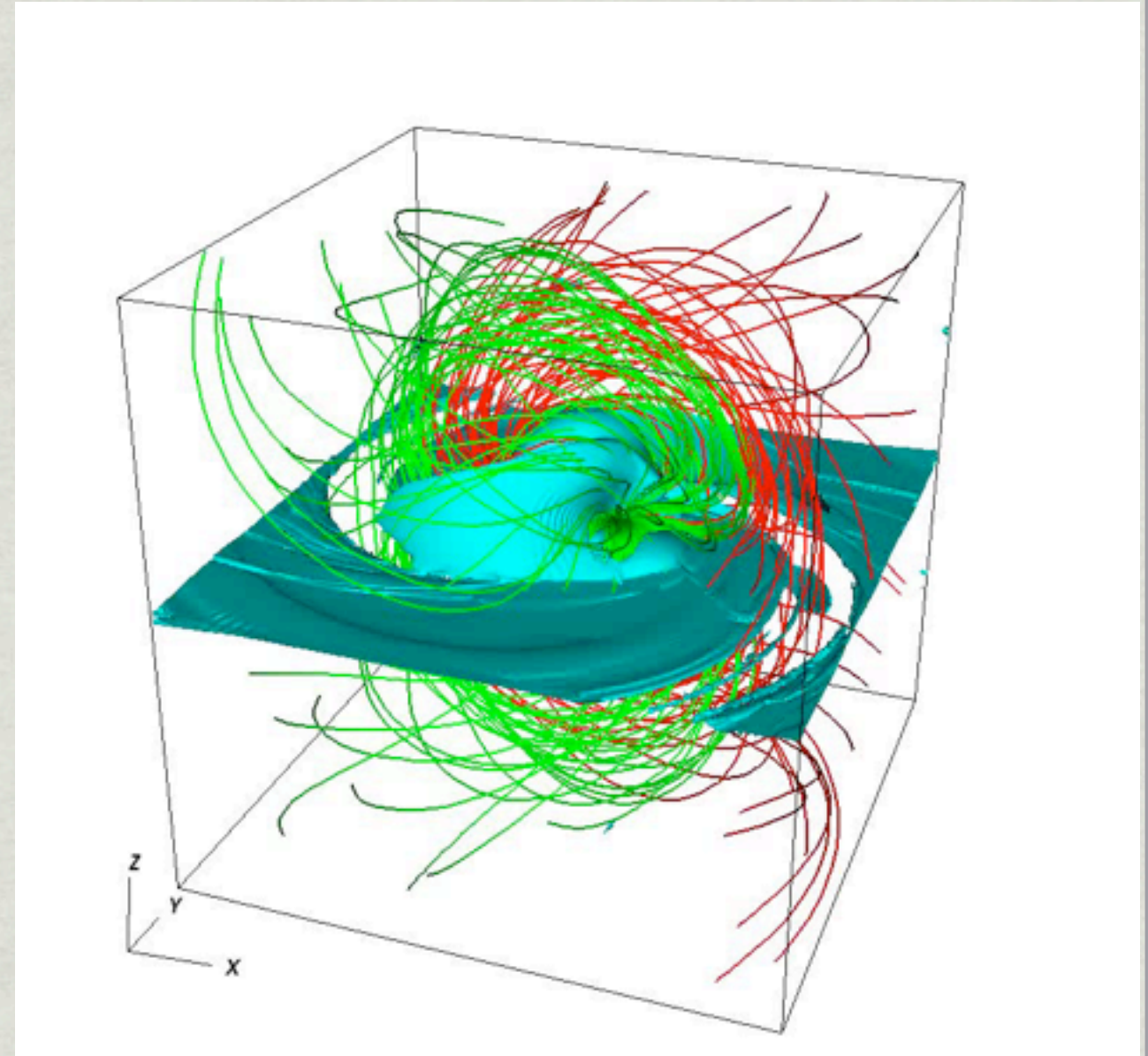


ms-period spin

Binaries: field patterns

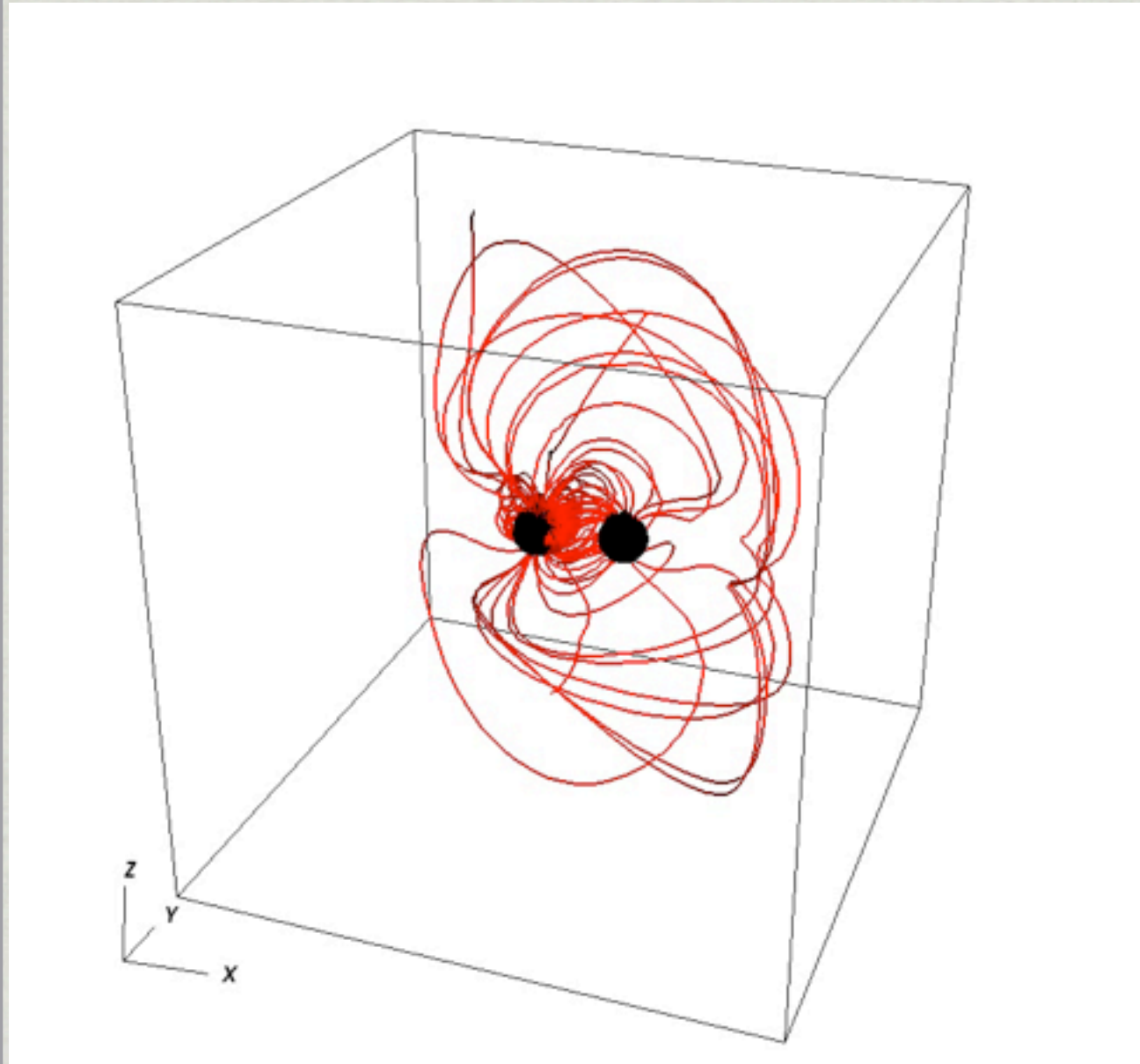


U/U, no spin

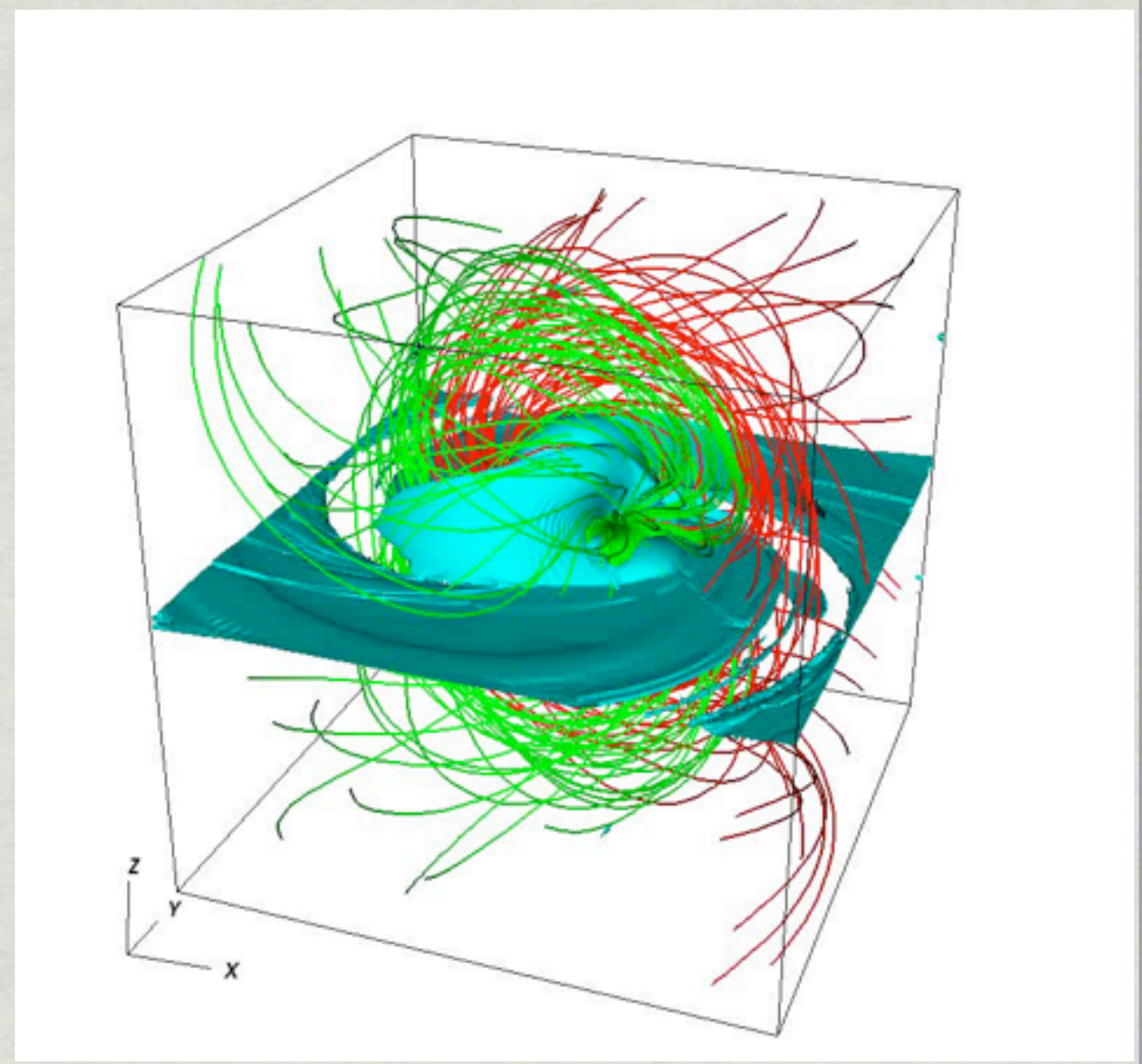


U/D, no spin

Binaries: field patterns



U/D, no NS spin



U/D, no NS spin

Analytic Guidance

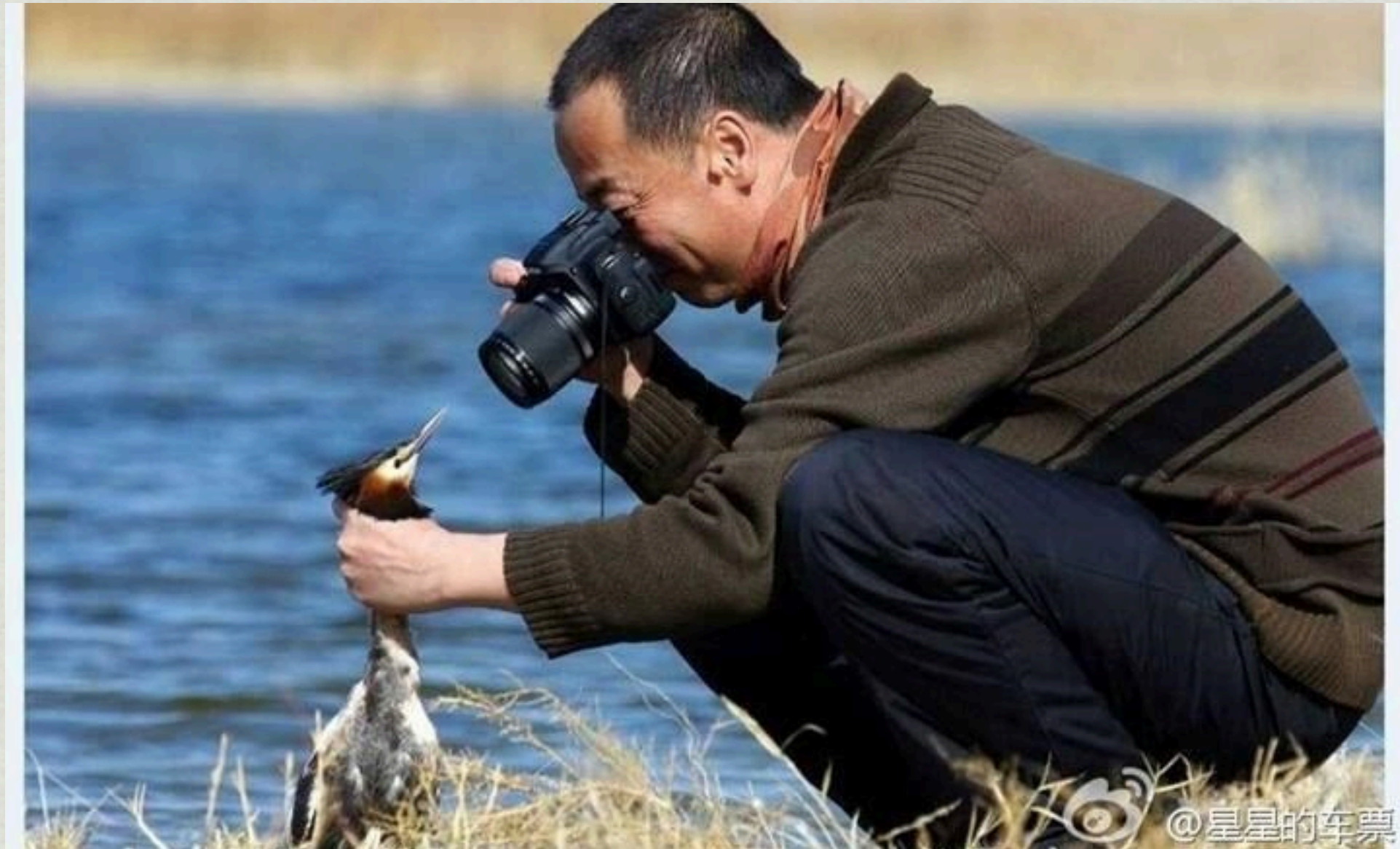
- In vacuum (e.g., Ioka & Taniguchi '00; Drell+ '65):
 - Electric dipole radiation $L \propto (\vec{v} \times \vec{\mu})'' \propto \Omega^{14/3}$
 - Cancels for aligned dipoles
 - Magnetic Quadrupole radiation $L \sim (r \mu)''' \propto \Omega^{14/3}$
 - Cancels for aligned dipoles
 - Potential drop $\Phi \approx (\Omega r_{\text{orb}}/c) B R_*$ yields $L \propto \Phi^2 \propto \Omega^{14/3}$
- In highly conducting plasma:
 - At LC $B_p \approx E_p \approx B_\phi \approx 1/R_{LC}^{1+\zeta}$ gives $L \propto \int \vec{E} \times \vec{B} d\vec{A} \propto \Omega^{2\zeta}$
 - $\zeta \rightarrow 0$ when $R_{LC,*} \rightarrow \infty$, $\zeta \rightarrow 2$ when $R_{LC,*} \rightarrow 0$
 - Aligned binaries:

$$L = 2 \frac{\mu^2 \Omega_*^4}{c^3} + \frac{\mu^2}{c^3} (1 + 6.5(\Omega_*/\Omega_0)^{1.7}) \Omega_0^{2.5\Omega_*/\Omega_0} \Omega^{4-2.5\Omega_*/\Omega_0},$$

Conclusions

- ✱ Magnetospheric shape is now known and confirmed in the limit of abundant plasma in 3D.
- ✱ Geometrically these models are being contrasted with gamma-ray observations (Separatrix Layer vs Gaps).
- ✱ More realistic models with 3D RMHD, cascade physics and full PIC are advancing -- expect cool results in the next few years. Benefit a lot from Moore's law.
- ✱ Reconnection may play an important and under-appreciated role in both emission and determining the magnetospheric shape.
- ✱ The origin of long time scale variability is likely not magnetospheric.

Conclusions



Conclusions

Pulsar
Magnetosphere

Particle
Simulations

