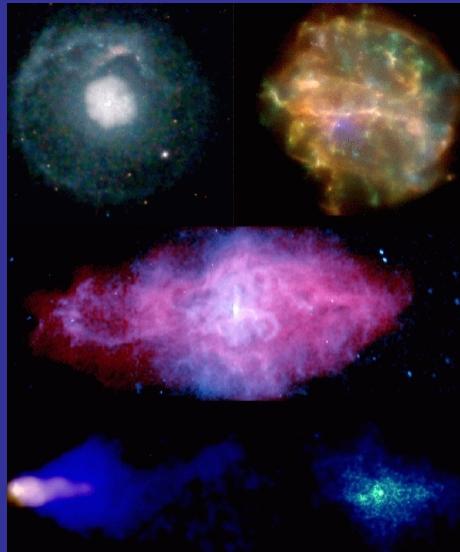


# Ionization Break-Out from Millisecond PWNe: Testing the Origin of Superluminous Supernovae

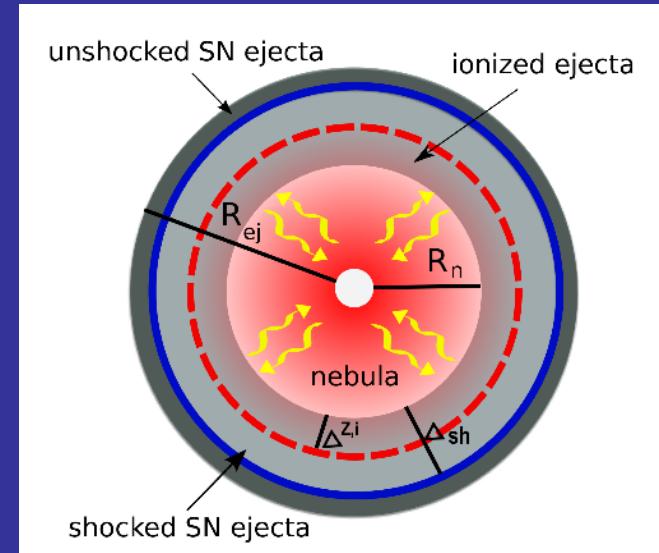


Brian Metzger  
Columbia University

In Collaboration with

Indrek Vurm, Romain Hascoet, Andrei Beloborodov (Columbia)

Andrew Levan (Warwick), Tony Piro (Caltech)



**BDM, Vurm, Hascoet & Beloborodov (2014), MNRAS, 437, 703**

Levan, Read, BDM, Wheatley, & Tanvir (2013) ApJ, 771, 136)

BDM & Piro (2014), MNRAS, 439, 3916

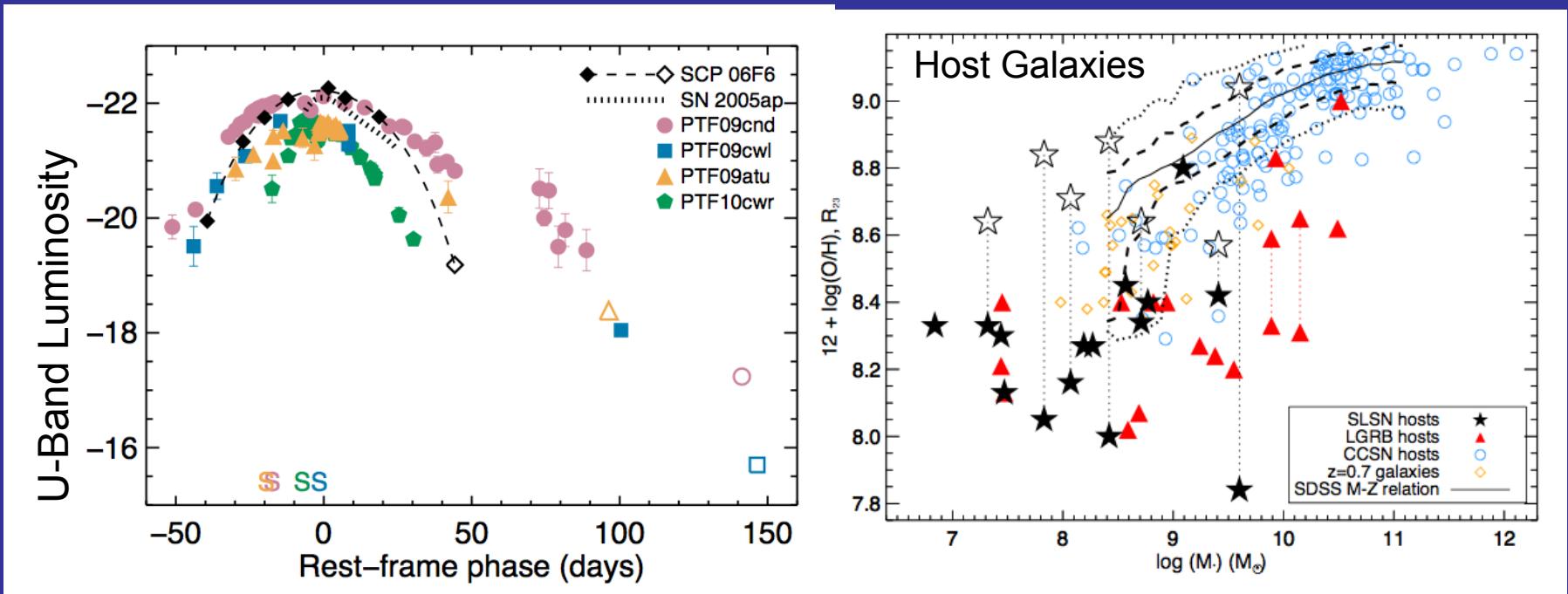
Relativistic Plasma Astrophysics, Purdue University - May 13, 2014

# H-Poor ‘Superluminous’ Supernovae

Quimby+07, Barbary+09, Pastorello+10, Chomiuk+11, Leloudas+12, Berger+12, Lunnan+13, Inserra+13

Quimby et al. 2011

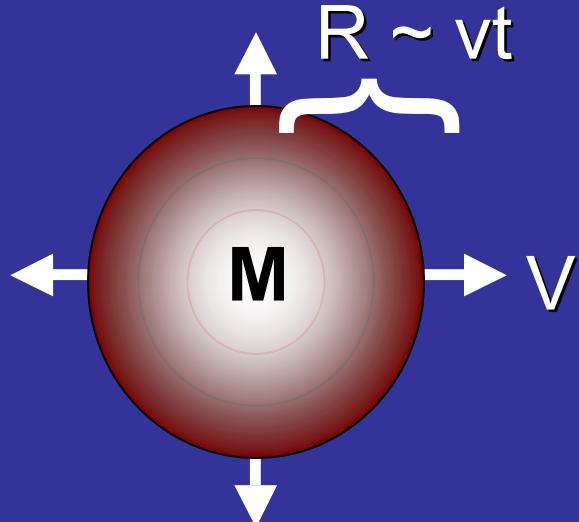
Lunnan et al. 2014



- $L_{\text{peak}} > 10^{44} \text{ erg s}^{-1}$ ,  $E_{\text{rad}} \sim 10^{50-51} \text{ ergs}$  ( $10-100 \times$  normal SNe)
- UV-rich spectrum with intermediate mass elements (CNO)
- Faint metal-poor host galaxies, similar to long GRBs  
(Quimby+11, Neill+11, Chomiuk+11, Chen+13; Lunnan+14; but see Berger+13, Chornock+13)
- Extremely Rare:  $\sim 10^{-4}$  of core collapse SN rate

# How Supernovae Shine

ejecta with mass  $M$ , velocity  $v$ , opacity  $\kappa$



$$R_0 \sim 10^{11.5} \text{ cm}$$

$$E_{th,0} \sim Mv^2 \sim 10^{51} \text{ ergs}$$

$$E_{th,pk} \sim E_{th,0}(R_0/R_{pk}) \sim 10^{-4} E_{th,0} \sim 10^{47} \text{ ergs}$$

$$\begin{aligned}L_{\text{pk}} &\sim 10^{47} \text{ ergs / t}_{\text{pk}} \\&\sim 10^{41} \text{ erg s}^{-1}\end{aligned}$$

$$R = v \cdot t$$

$$\rho \sim M/R^3$$

$$\tau \sim \kappa \rho R \quad t_{\text{diff}} \sim \frac{\tau R}{c} \propto t^{-1}$$

# PdV work

# radiation

$$\frac{dE_{\text{th}}}{dt} = -\frac{E_{\text{th}}}{t} - \frac{E_{\text{th}}}{t_{\text{diff}}} + ???$$

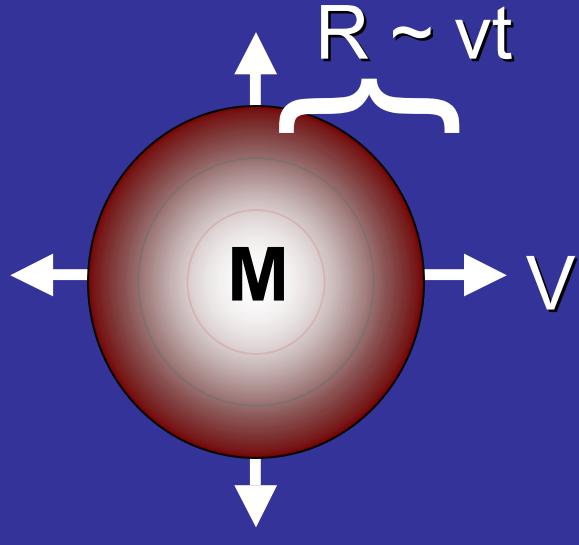
Luminosity peaks when  $t \sim t_{\text{diff}} \Rightarrow$

$$t_{\text{pk}} \sim \text{month} \left( \frac{v}{10^4 \text{ km s}^{-1}} \right)^{-1/2} \left( \frac{M}{3M_{\odot}} \right)^{1/2} \left( \frac{\kappa}{0.1 \text{ cm}^2 \text{g}^{-1}} \right)^{1/2}$$

$$R_{\text{pk}} \sim 3 \times 10^{15} \text{cm} \left( \frac{v}{10^4 \text{ km s}^{-1}} \right)^{1/2} \left( \frac{M}{3M_{\odot}} \right)^{1/2} \left( \frac{\kappa}{0.1 \text{ cm}^2 \text{g}^{-1}} \right)^{1/2}$$

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# PdV work

# **radiation**

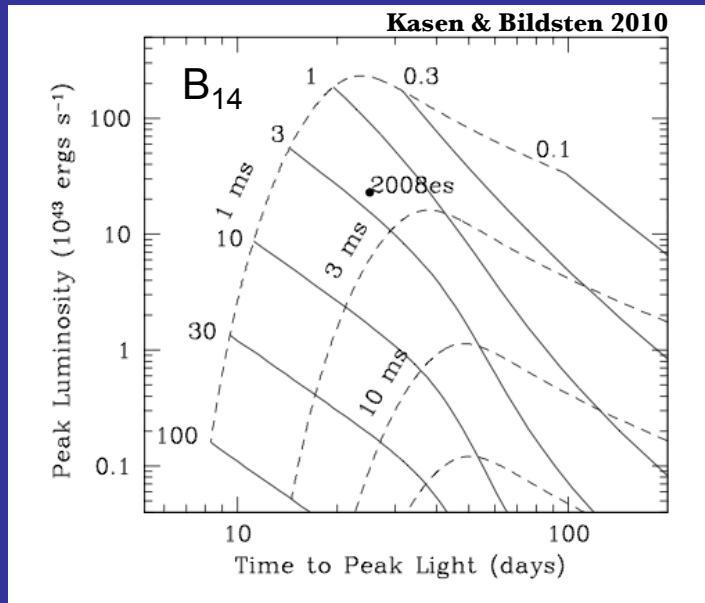
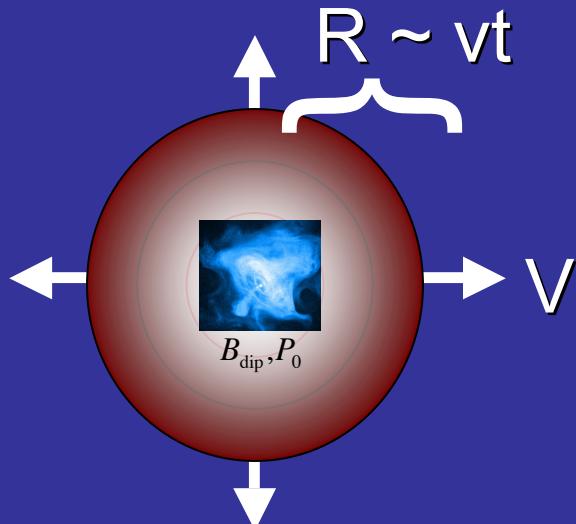
$$\frac{dE_{\text{th}}}{dt} = -\frac{E_{\text{th}}}{t} - \frac{E_{\text{th}}}{t_{\text{diff}}} + \dot{E}_{\text{Ni56}}$$

Luminosity peaks when  $t \sim t_{\text{diff}} \Rightarrow$

$$\dot{E}_{\text{Ni56}} \sim \frac{M_{\text{Ni56}} \epsilon}{\tau} e^{-t/\tau} \sim 10^{43} \left( \frac{M_{\text{Ni56}}}{0.2 M_{\odot}} \right) e^{-t/10d} \text{erg s}^{-1}$$

# Pulsar Powered Supernovae

(Ostriker & Gunn 1971; Kasen & Bildsten 2010; Woosley 2010)



**PdV work**      **radiation**

$$\frac{dE_{\text{th}}}{dt} = -\frac{E_{\text{th}}}{t} - \frac{E_{\text{th}}}{t_{\text{diff}}} + \dot{E}_{\text{pulsar}}$$

$$\dot{E}_{\text{pulsar}} = \frac{\mu^2 \Omega^4}{c^3} \approx 4 \times 10^{45} \left( \frac{P}{2 \text{ ms}} \right)^{-4} \left( \frac{B_{\text{dip}}}{10^{13.5} \text{ G}} \right)^2 \text{ erg s}^{-1}$$

$$\tau_{\text{sd}} = \frac{E_{\text{rot}}}{L_{\text{sd}}} \approx \text{month} \left( \frac{P_0}{2 \text{ ms}} \right)^2 \left( \frac{B_{\text{dip}}}{10^{13.5} \text{ G}} \right)^{-2}$$

$$t_{\text{pk}} \sim \text{month} \left( \frac{v}{10^4 \text{ km s}^{-1}} \right)^{-1/2} \left( \frac{M}{3M_{\odot}} \right)^{1/2} \left( \frac{\kappa}{0.1 \text{ cm}^2 \text{ g}^{-1}} \right)^{1/2}$$

Increase SN luminosity if  $\tau_{\text{sd}} \sim t_{\text{pk}} \Rightarrow B_{\text{dip}} \sim 10^{13-14} \text{ G}, P \sim 1\text{-few ms}$

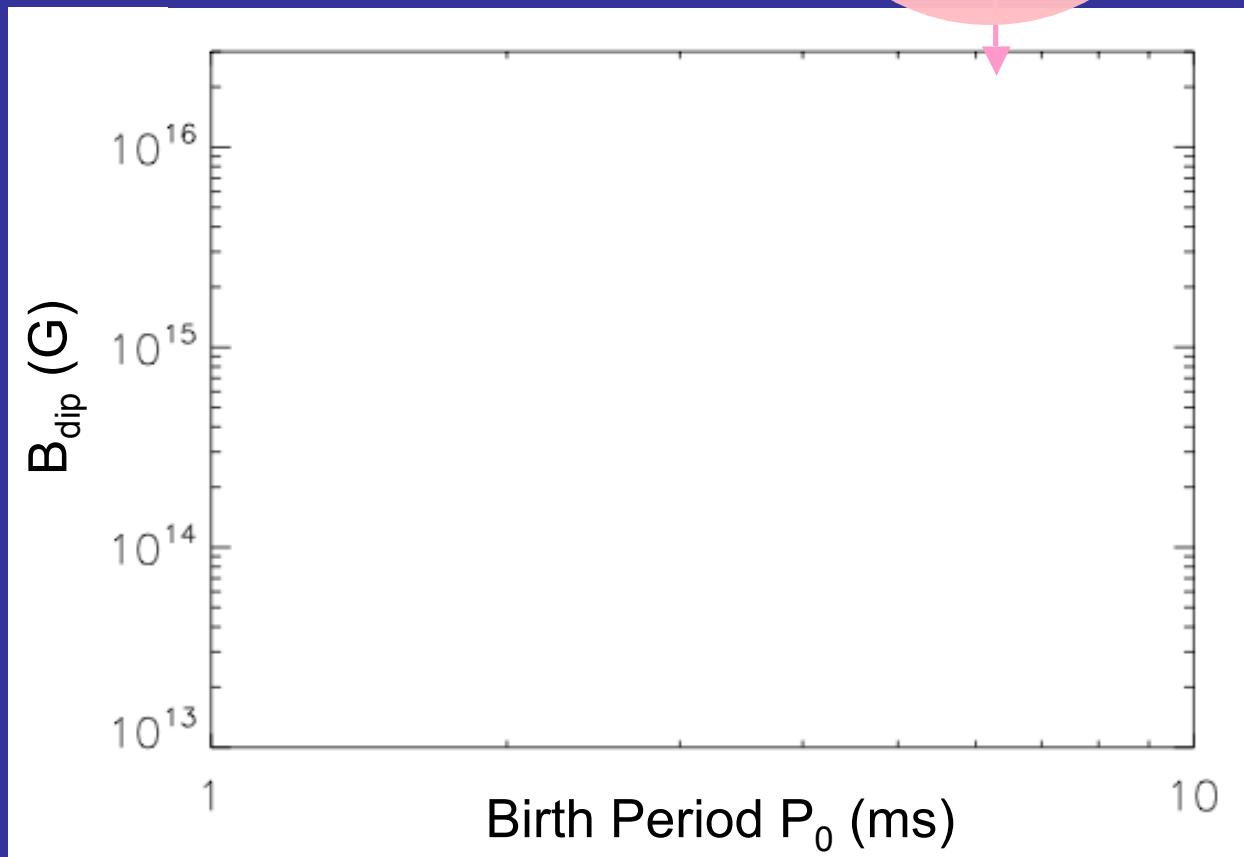
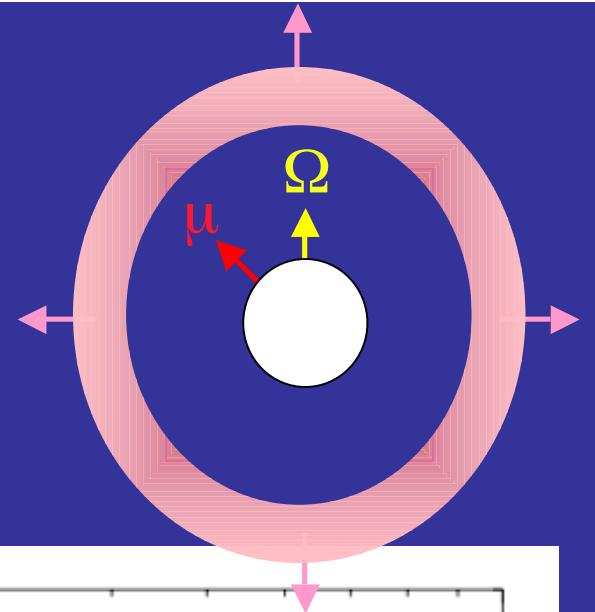
# Signatures of Magnetar Birth

spin-down ·  
luminosity ·

$$L_{\text{sd}} = \frac{\mu^2 \Omega^4}{c^3} \approx 6 \times 10^{49} \left( \frac{P}{1 \text{ ms}} \right)^{-4} \left( \frac{B_{\text{dip}}}{10^{15} \text{ G}} \right)^2 \text{ erg s}^{-1}$$

spin-down time : ·

$$\tau_{\text{sd}} = \frac{E_{\text{rot}}}{L_{\text{sd}}} \approx 10 \left( \frac{P_0}{1 \text{ ms}} \right)^2 \left( \frac{B_{\text{dip}}}{10^{15} \text{ G}} \right)^{-2} \text{ min}$$



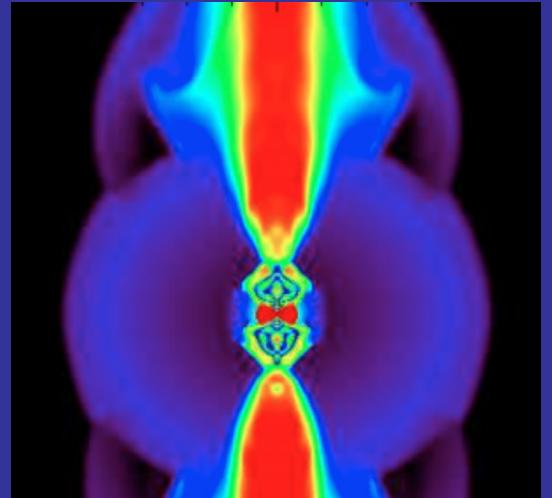
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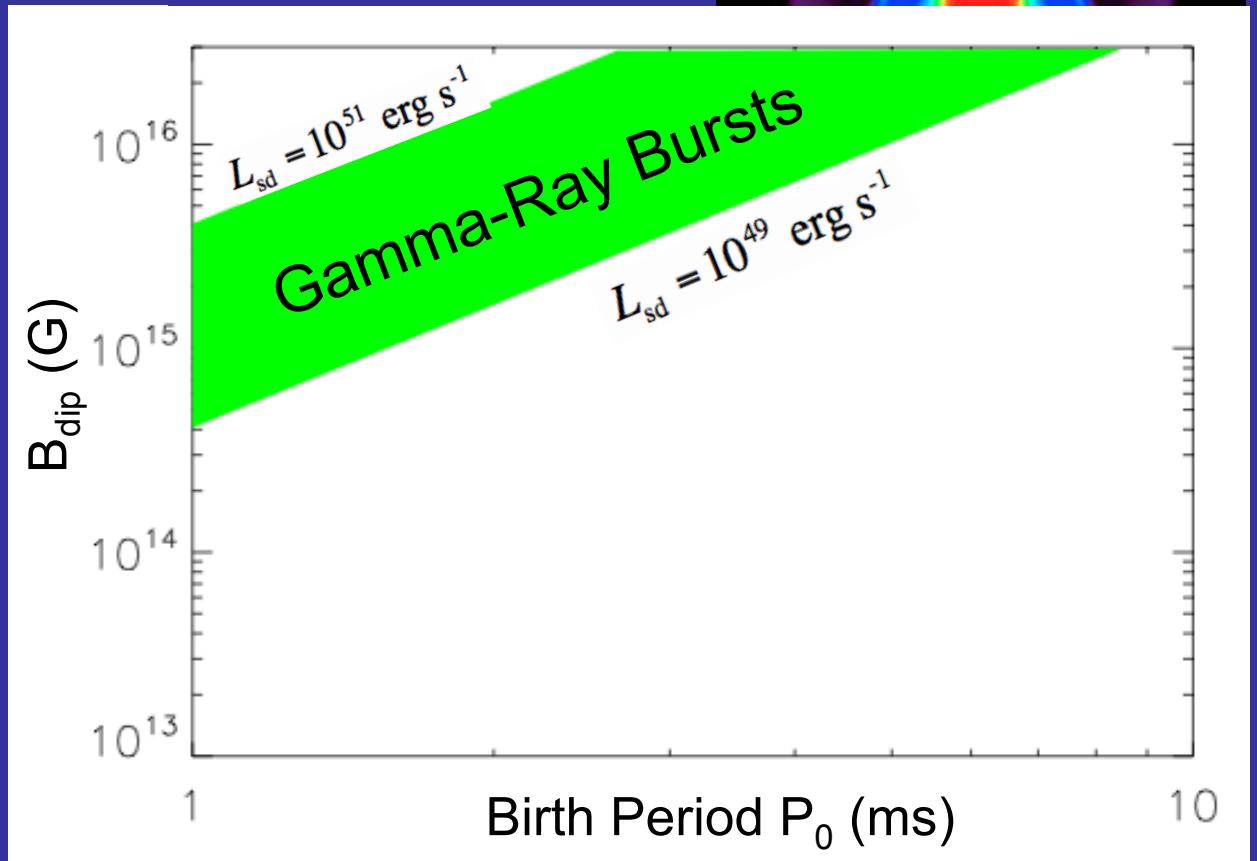
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## Gamma-Ray Burst

- Jet punches successfully through star (???)
- $L_{\text{sd}} \sim L_{\gamma} \sim 10^{49-51} \text{ erg s}^{-1}$
- $\tau_{\text{sd}} \sim \text{minutes-hours}$



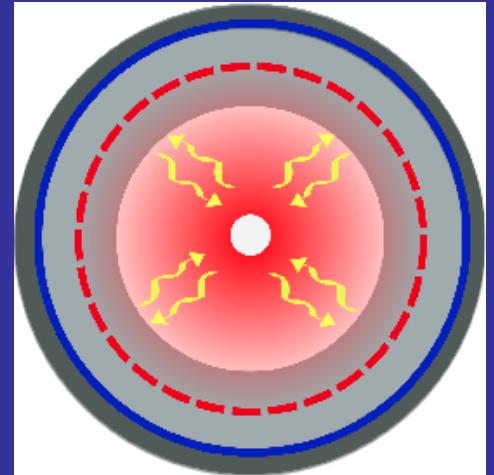
# Signatures of Magnetar Birth

spin-down luminosity :

$$L_{\text{sd}} = \frac{\mu^2 \Omega^4}{c^3} \approx 6 \times 10^{49} \left( \frac{P}{1 \text{ ms}} \right)^{-4} \left( \frac{B_{\text{dip}}}{10^{15} \text{ G}} \right)^2 \text{ erg s}^{-1}$$

spin-down time :

$$\tau_{\text{sd}} = \frac{E_{\text{rot}}}{L_{\text{sd}}} \approx 10 \left( \frac{P_0}{1 \text{ ms}} \right)^2 \left( \frac{B_{\text{dip}}}{10^{15} \text{ G}} \right)^{-2} \text{ min}$$

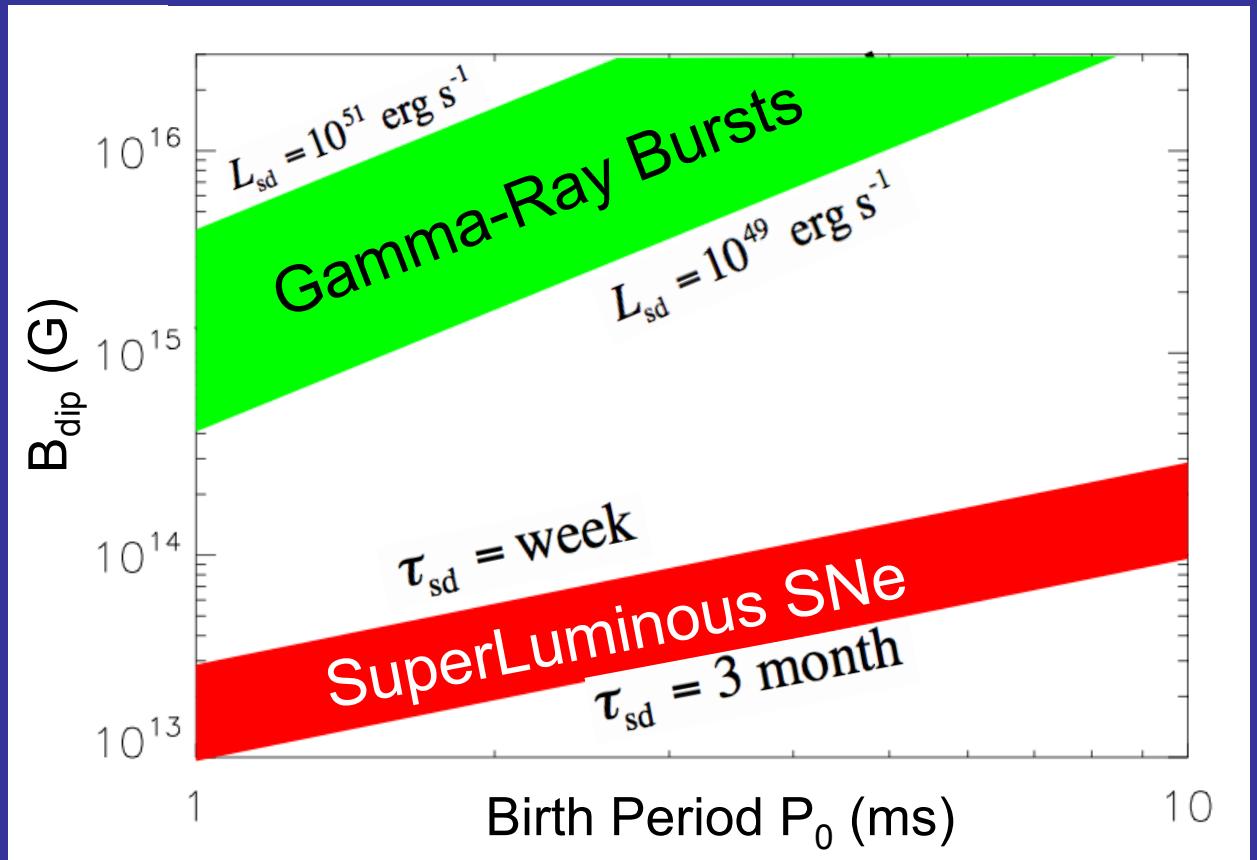


## Gamma-Ray Burst

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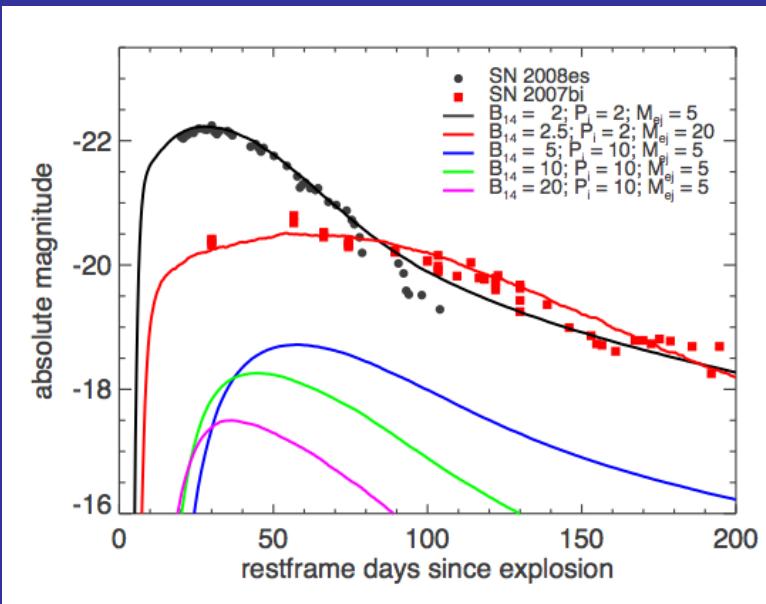
## Super-Luminous SN

- Jet stifled, but optical SN powered diffusively
- $L_{\text{sd}} \sim L_{\text{SN}} \sim 10^{43-45} \text{ erg s}^{-1}$
- $\tau_{\text{sd}} \sim \text{week - months}$

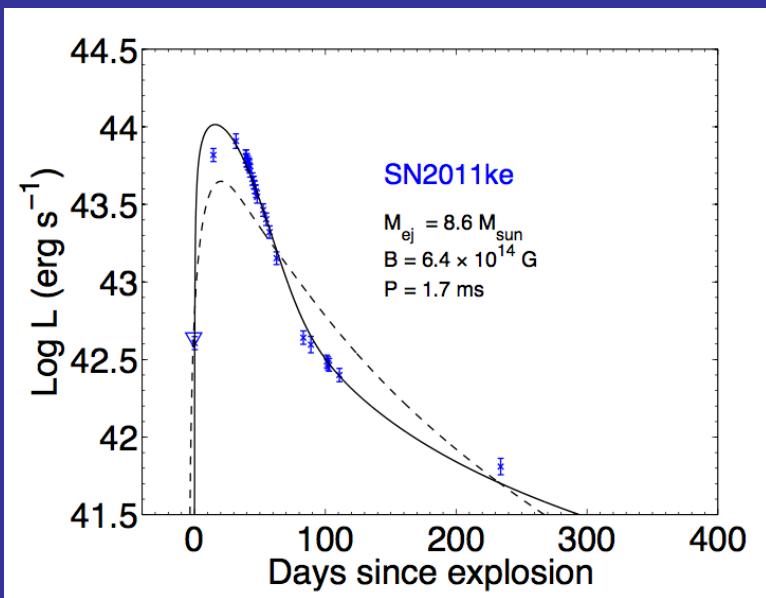


# “Magnetar” Model for Superluminous SNe

Kasen & Bildsten 2010



Inserra et al. 2013



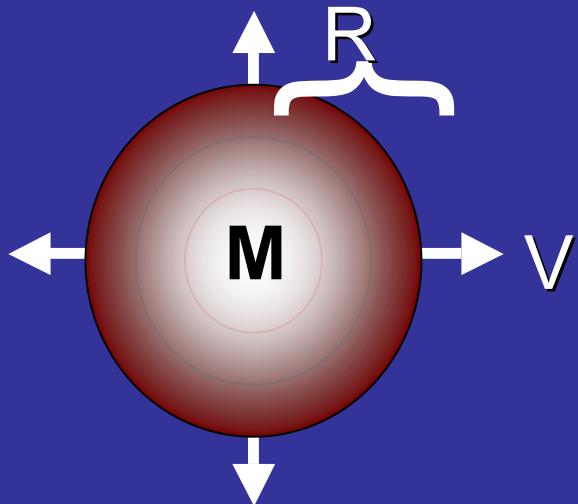
## PROS

- “Explains” similar host galaxies to long GRBs (both require rapid rotating progenitor)
- Can reproduce diversity of rise times and peak luminosities

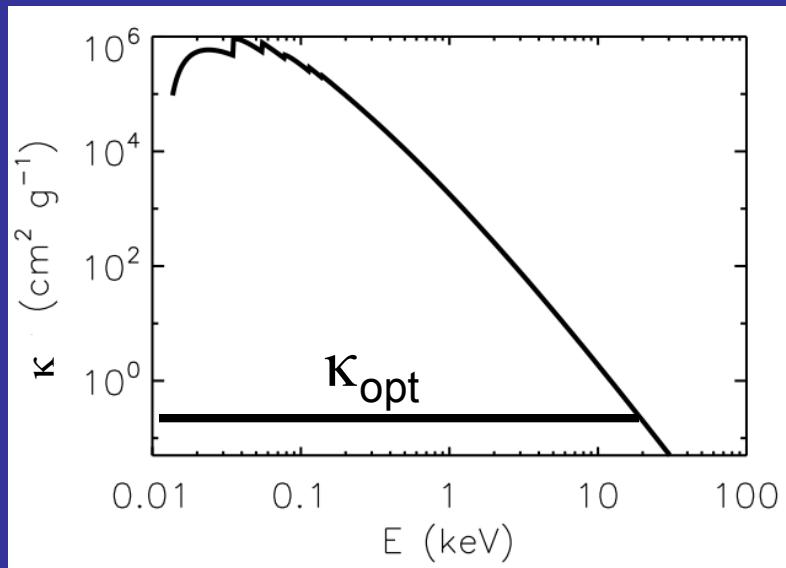
## CONS

- Can reproduce diversity of rise times and peak luminosities (hard to test)
  - How to distinguish from other hidden energy sources (optically-thick CSM interaction)
  - Assumes pulsar wind thermalizes with 100% efficiency.
- Reality:** Poynting flux  $\Rightarrow e^{+/-}$  pairs  $\Rightarrow$  non-thermal radiation  $\Rightarrow$  thermal radiation

# Can we see inside nebula directly?



UV/X-ray opacity of neutral oxygen



Light escapes when  $t > t_{\text{diff}} \Rightarrow$

$$t > \text{month} \left( \frac{v}{10^4 \text{ km s}^{-1}} \right)^{-1/2} \left( \frac{M}{3M_\odot} \right)^{1/2} \left( \frac{\kappa}{0.1 \text{ cm}^2 \text{g}^{-1}} \right)^{1/2}$$

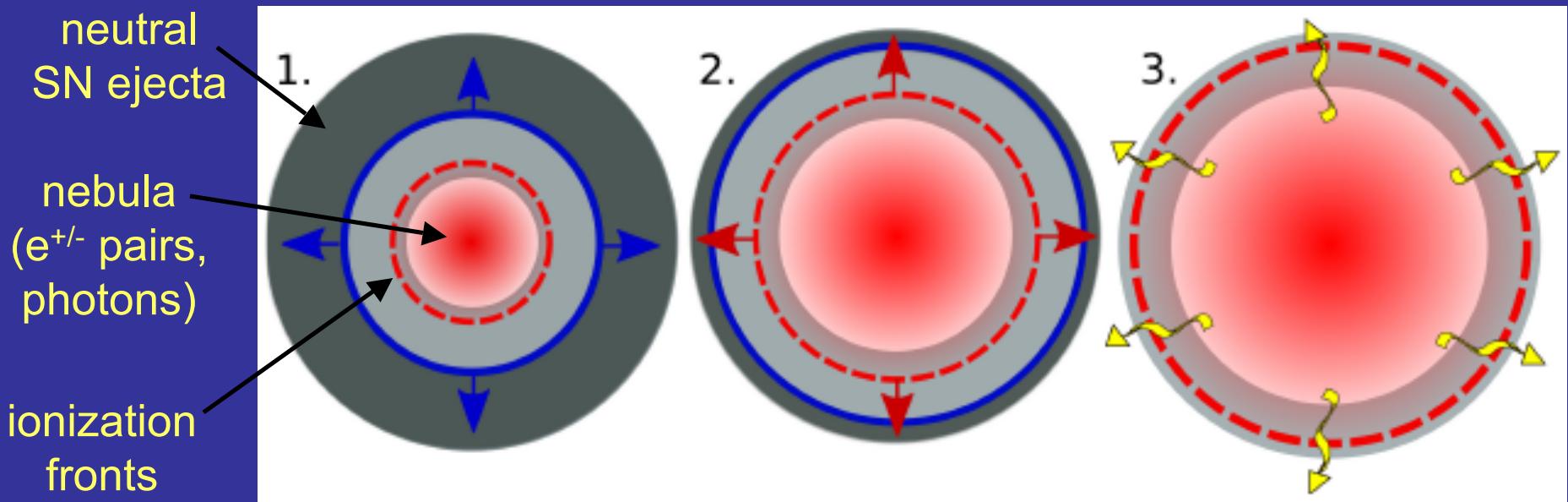
Optical ( $\kappa \sim 0.1 \text{ cm}^2 \text{g}^{-1}$ )  $\Rightarrow t \sim \text{month}$

Hard X-ray  $\sim 10 \text{ keV}$  ( $\kappa \sim 30 \text{ cm}^2 \text{g}^{-1}$ )  
 $\Rightarrow t \sim \text{few years}$

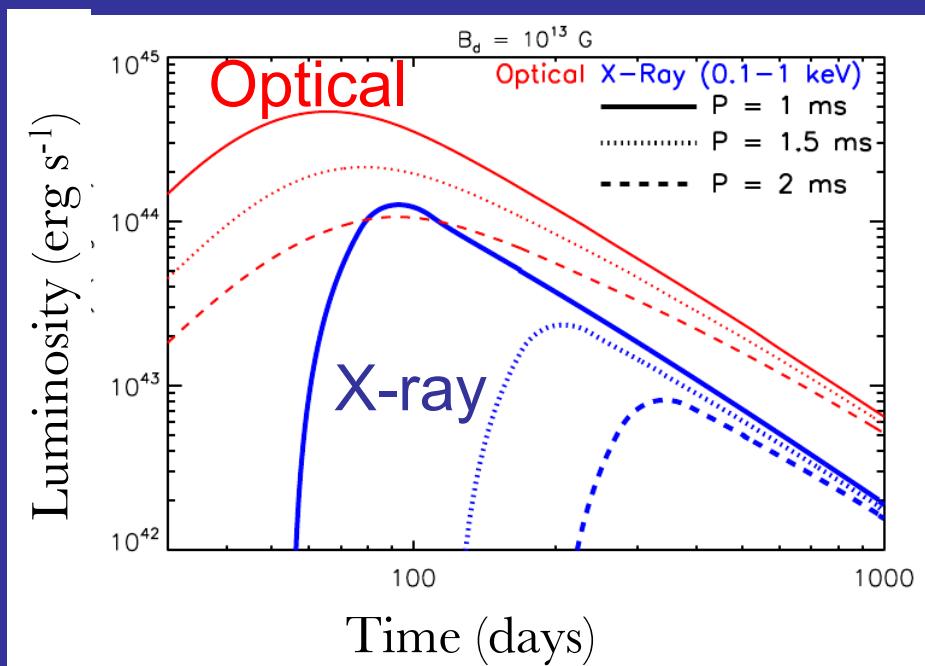
Soft X-ray  $\sim 0.1 \text{ keV}$  ( $\kappa \sim 10^5 \text{ cm}^2 \text{g}^{-1}$ )  
 $\Rightarrow t \sim 100 \text{ years}$

# X-ray Ionization Break-Out

(BDM, Vurm, Hascoet & Beloborodov 2013)

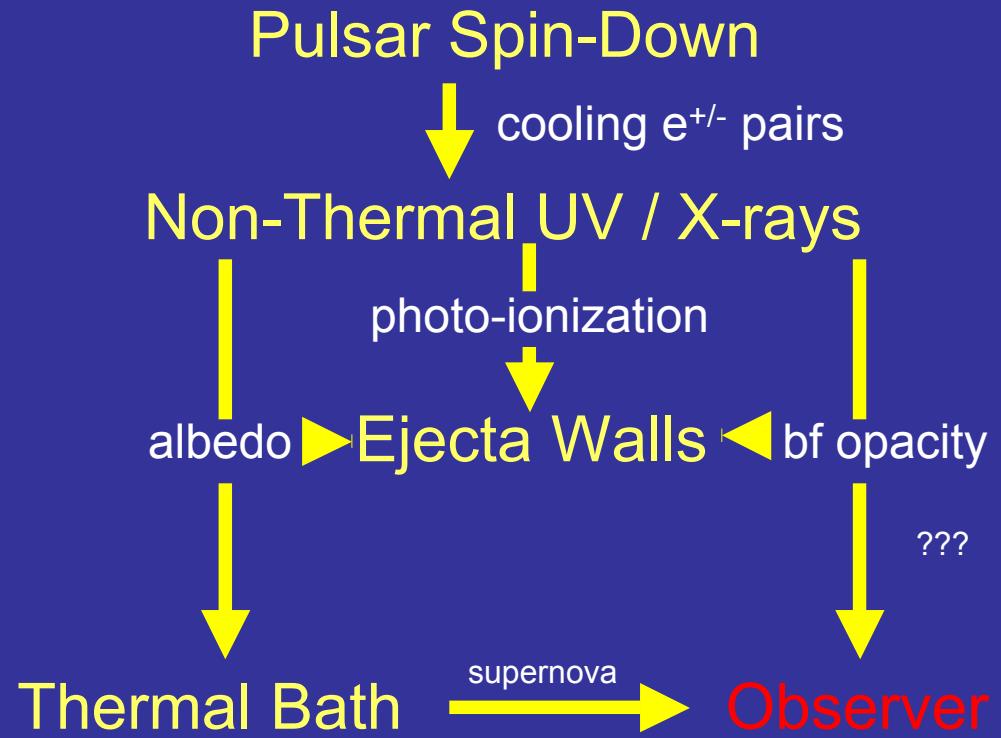
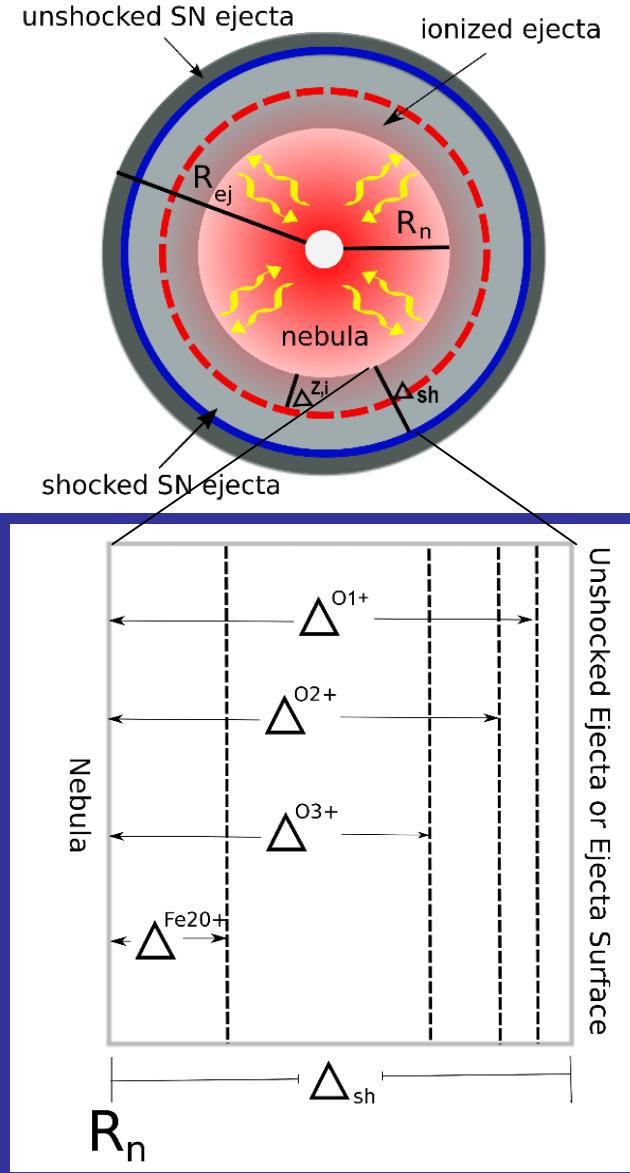


1. Pulsar inflates nebular cavity.
2. Nebula X-rays ionize inner exposed surface of ejecta
3. Ionization front reaches outer surface  $\Rightarrow$   
X-rays escape to observer.



# Model for Evolution of Millisecond Pulsar Wind Nebulae

(BDM, Vurm, Hascoet & Beloborodov 2013)



Pairs injected into nebula at rate

$$\dot{N}_{\pm} = \mu_{\pm} \dot{N}_{\text{GJ}} \simeq 5 \times 10^{37} \mu_{\pm} B_{13} P_{-3}^{-2} (1 + t/t_{\text{sd}})^{-1} \text{s}^{-1}$$

...with energy per particle

$$\epsilon_{\pm} = \frac{L_{\text{sd}}}{\dot{N}_{\pm}} \simeq 1.1 \times 10^8 \text{ erg} \mu_{\pm}^{-1} B_{13} P_{-3}^{-2} (1 + t/t_{\text{sd}})^{-1}$$

...and particle lorentz factor

$$\gamma_{\pm} = \frac{\epsilon_{\pm}}{m_e c^2} \simeq 1.3 \times 10^{14} \mu_{\pm}^{-1} B_{13} P_{-3}^{-2} (1 + t/t_{\text{sd}})^{-1}$$

...which will actually be limited by cooling to a much lower value:

...producing photons of energy:

sufficient to create pairs, which given the high compactness of the nebula,

$$\gamma_{\pm}^{\text{max}} = \sqrt{\frac{3e}{\sigma_T B_n}} \simeq 5 \times 10^6 \epsilon_{\text{B},-2}^{-1/4} L_{45}^{-1/4} v_9^{1/2} M_3^{1/4} \left( \frac{t}{t_{d,0}} \right)^{1/2}$$

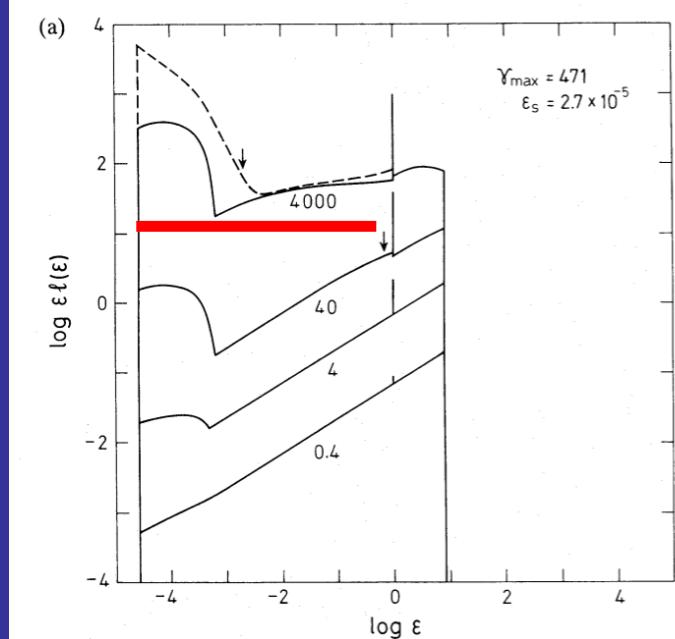
$$\hbar e B_n (\gamma_{\pm}^{\text{max}})^2 / m_e c \sim 100 \text{ MeV}$$

$$\ell \equiv \frac{E_{\text{nth}} \sigma_T R_n}{V_n m_e c^2} \approx 2.1 (1 - \bar{\mathcal{A}}_{\nu})^{-1} L_{45} v_9^{-1/2} M_3^{-1/2} \left( \frac{t}{t_{d,0}} \right)^{-1}$$

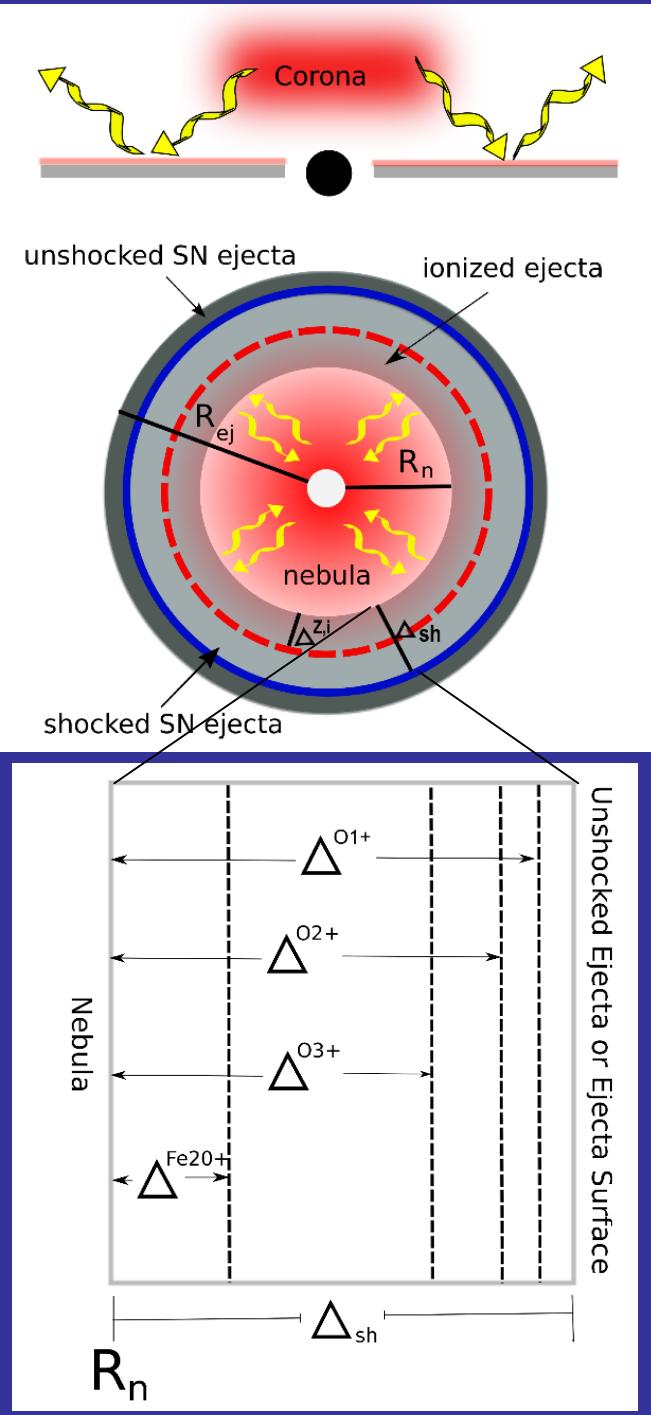
“saturated” pair cascades  $\Rightarrow$   
flat photon spectrum

$$\dot{E}_{\text{sd},\nu} \simeq \frac{L_{\text{sd}}}{14\nu} \quad (3kT_{\text{th}} \sim 1 \text{ eV} \lesssim h\nu \lesssim 1 \text{ MeV})$$

e.g. Zdziarski & Lightman 1985

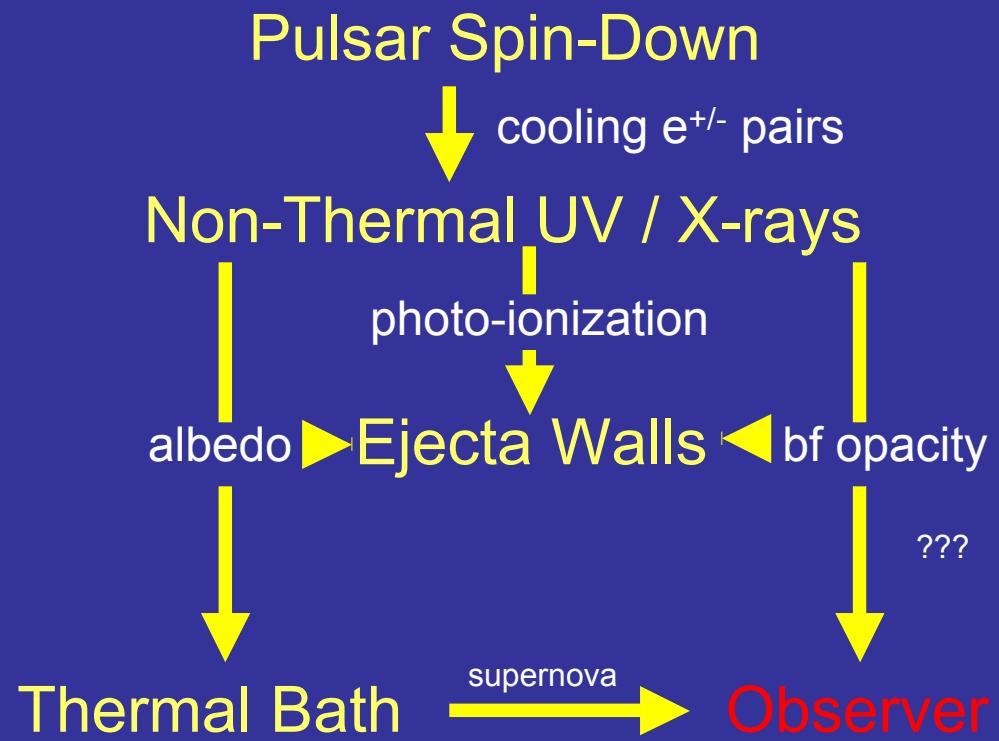


Svensson 1987



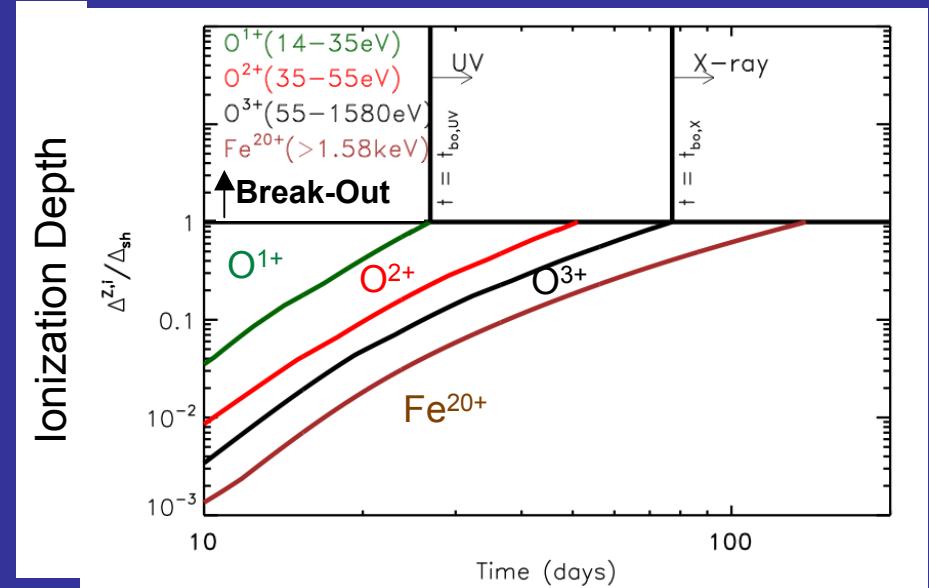
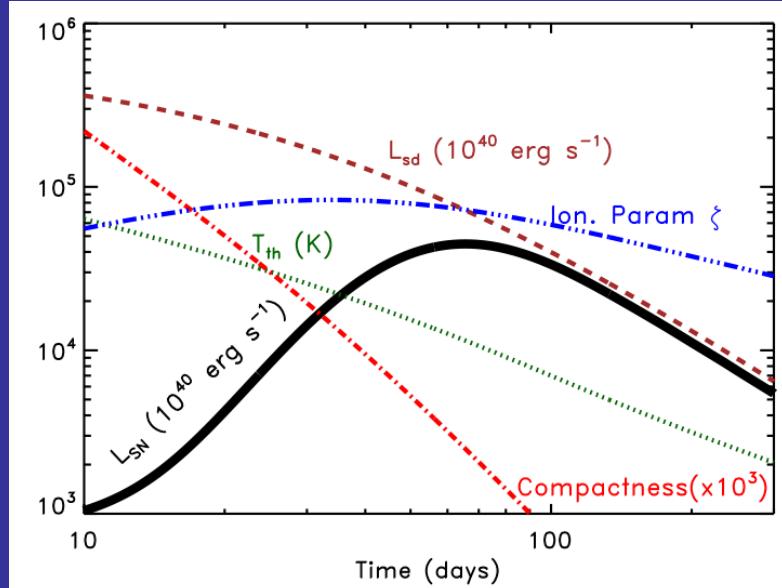
# Model for Evolution of Millisecond Pulsar Wind Nebulae

(BDM, Vurm, Hascoet & Beloborodov 2013)

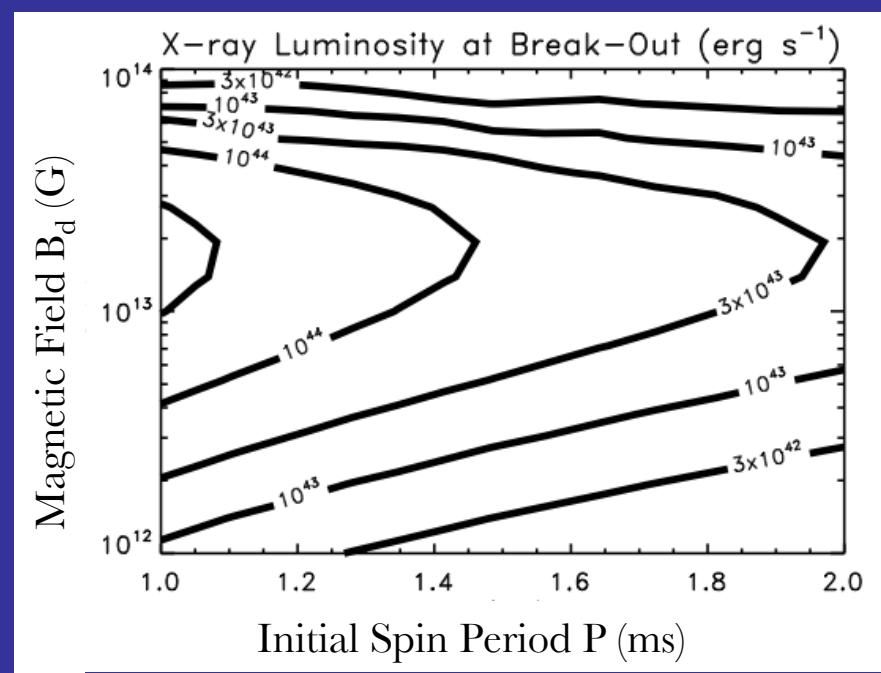
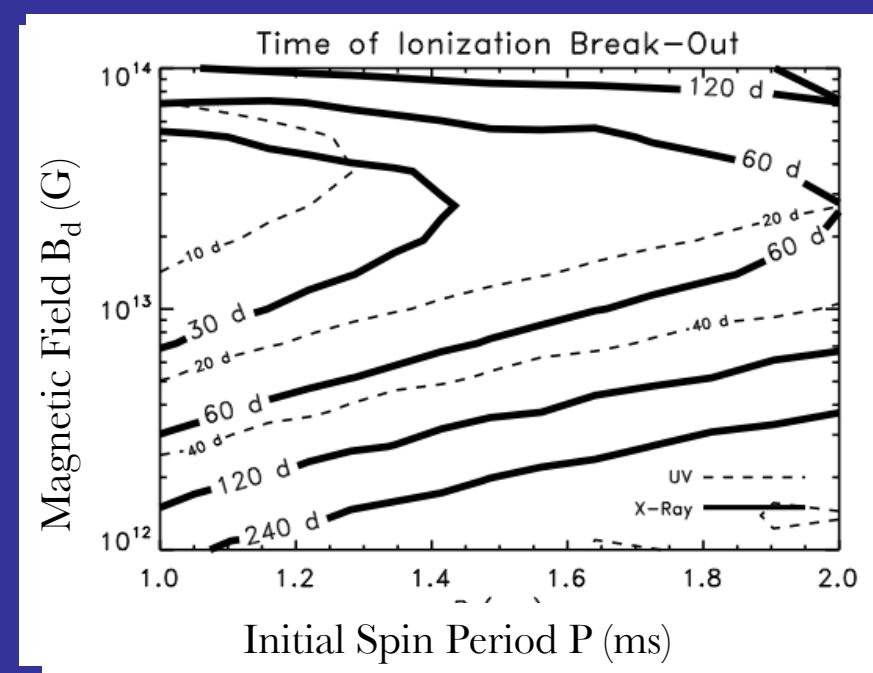
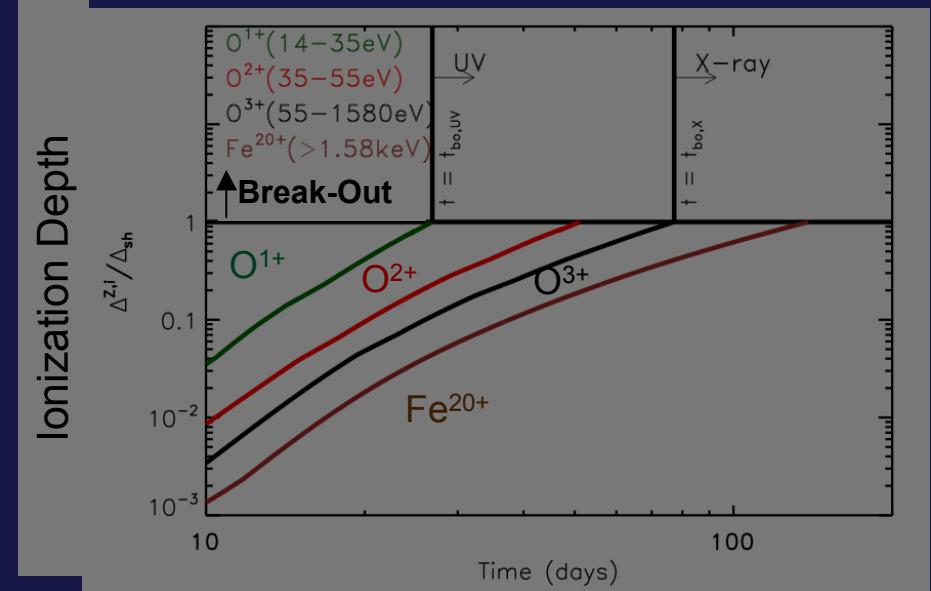
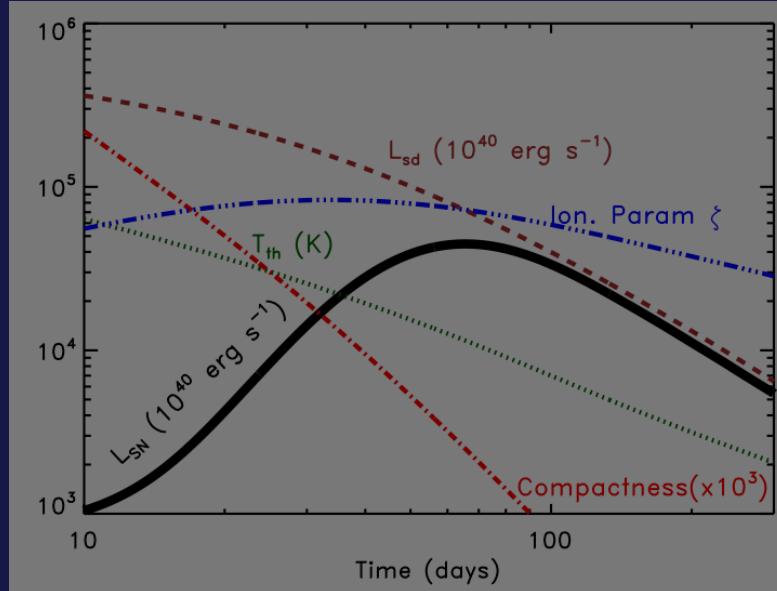


analogy to AGN accretion disks

Example:  $B = 10^{13}$  G,  $P = 1$  ms,  $M_{ej} = 3 M_\odot$



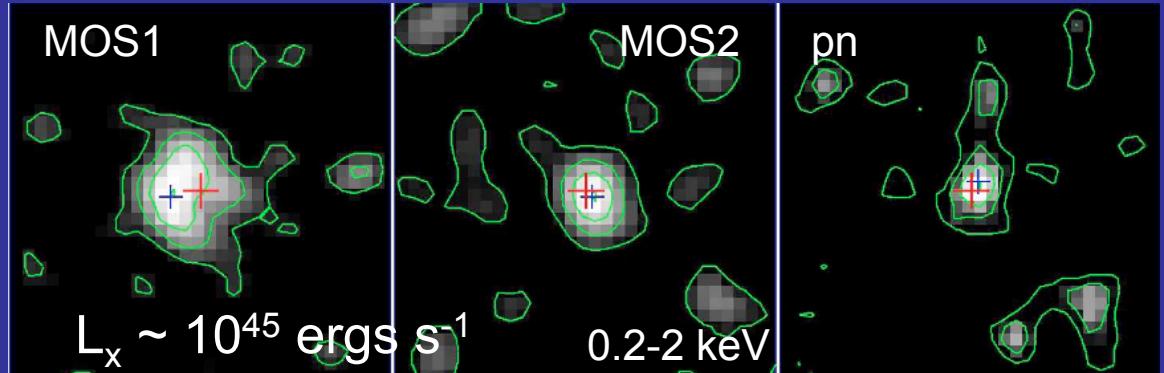
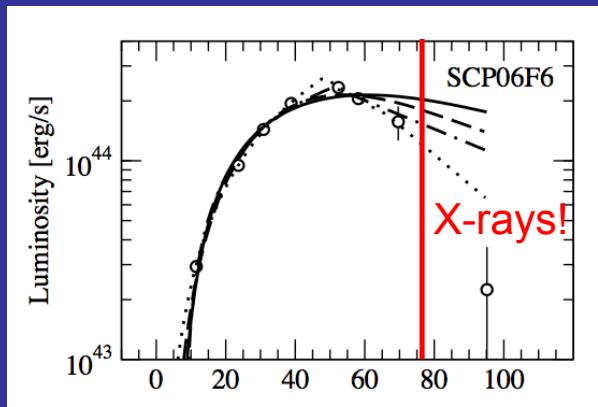
Example:  $B = 10^{13}$  G,  $P = 1$  ms,  $M_{ej} = 3 M_\odot$



# Superluminous X-rays from a Superluminous SN

(Levan, Read, BDM, Wheatley, Tanvir 2013)

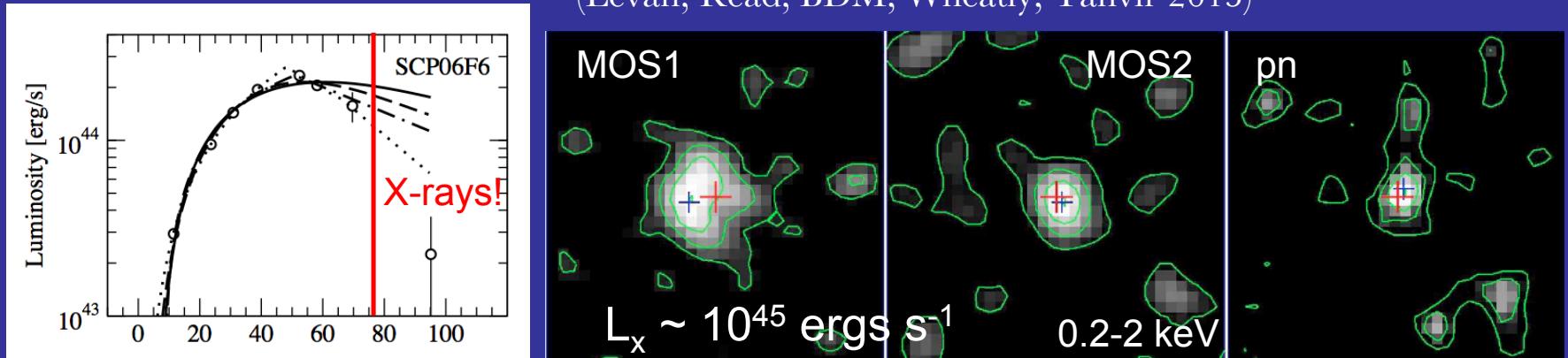
Chatzopoulos et al. 2013  
(cf. Barbary et al. 2009)



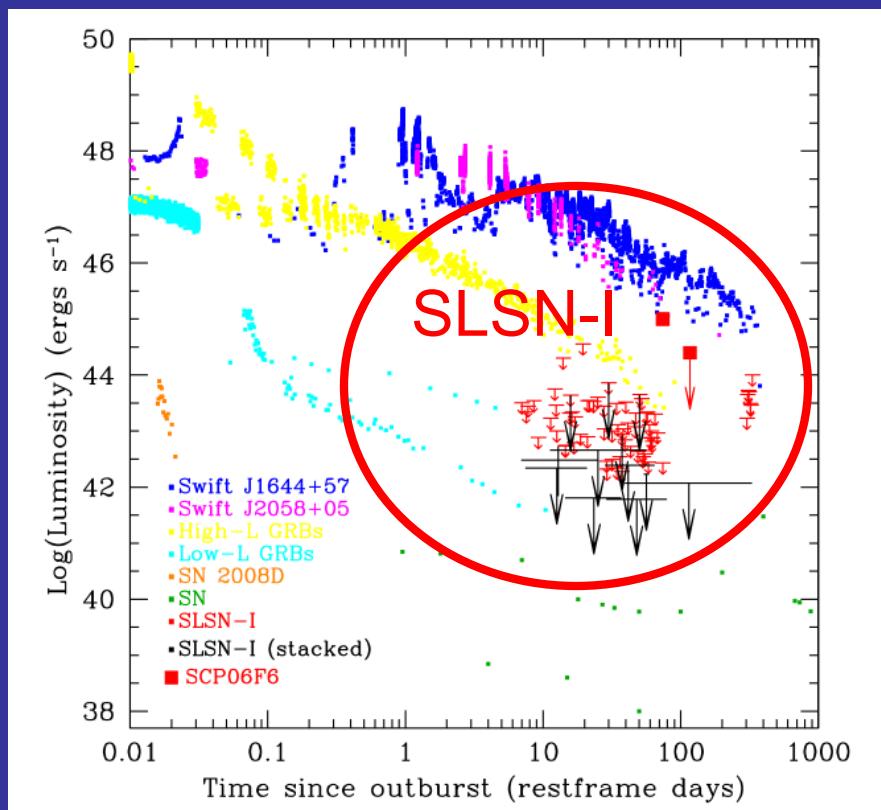
# Superluminous X-rays from a Superluminous SN

(Levan, Read, BDM, Wheatley, Tanvir 2013)

Chatzopoulos et al. 2013  
(cf. Barbary et al. 2009)



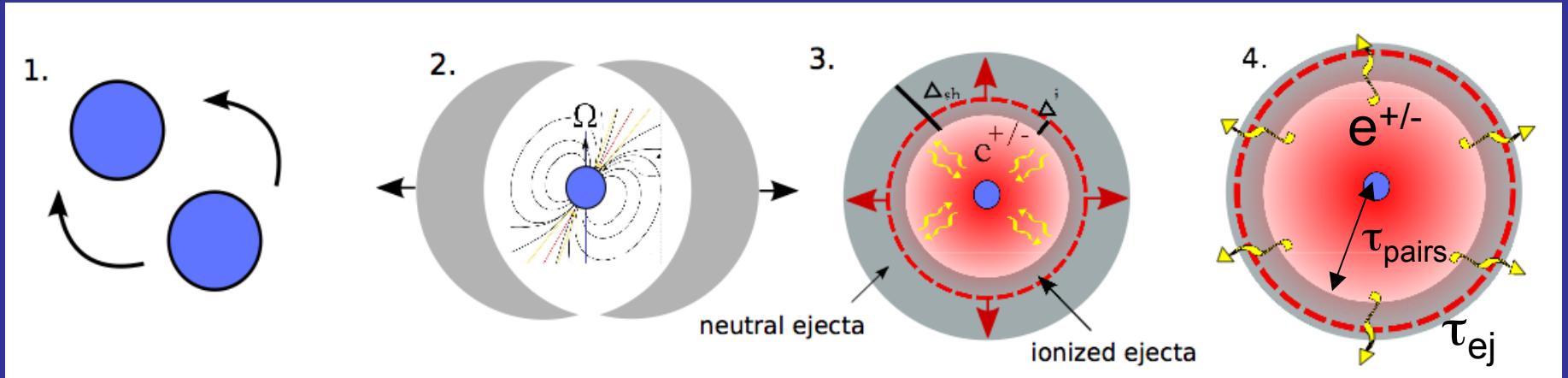
Levan et al. 2013



- No X-ray detections from other SLSNe
- Limits  $L_x < 10^{42}-10^{44}$  erg s $^{-1}$  on timescales  $< 70$  days (usually too early!)
- Future: X-ray follow-up *after* optical peak confirm/constrain pulsar model for SLSNe

# Application to Binary Neutron Star Mergers

BDM & Piro 2014



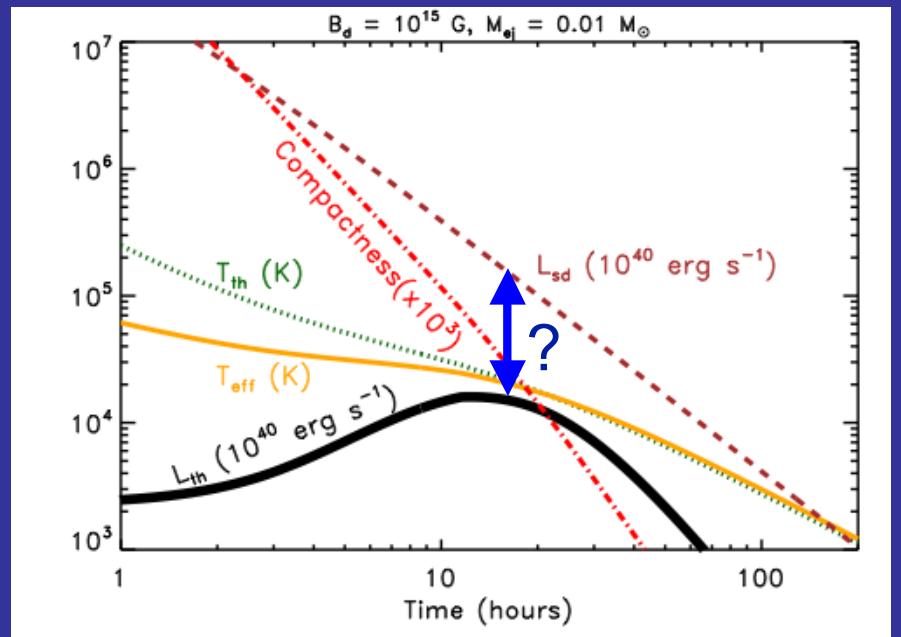
peak optical luminosity

$$L_{\text{opt,peak}} \approx \chi^{-3/2} L_{\text{sd}} \Big|_{t_{\text{peak}}}$$

$$\chi \equiv \frac{\tau_{\text{pairs}}}{\tau_{\text{ej}}} \Big|_{t_{\text{peak}}} \approx 2.7 \left( \frac{B_d}{10^{15} \text{ G}} \right)^{-2/3} \left( \frac{M_{\text{ej}}}{10^{-2} M_\odot} \right)^{-2/3} \left( \frac{v_{\text{ej}}}{c} \right)^{5/6}$$

$\chi > 1 \Rightarrow$  X-rays lose energy to PdV work before reaching walls (thermalization)

example ( $M_{\text{ej}} = 10^{-2} M_\odot$ ;  $B_d = 10^{15} \text{ G}$ )



# Summary

- Hydrogen-Poor SLSNe may be powered by spin-down of nascent millisecond pulsar
- However, previous models [1] challenging to test (degenerate) and [2] assume pulsar luminosity thermalizes
- We have developed a model for the evolution of millisecond PWNe that couples X-ray to thermal radiation via interaction with ejecta walls (self-consistent)
- Pulsar wind  $\Rightarrow e^{+/-}$  pairs  $\Rightarrow$  X-rays  $\Rightarrow$  thermal (optical) photons  $\Rightarrow$  observer (optical SN)
- High X-ray opacity of neutral matter normally prevents X-rays from escaping SN ejecta for  $>$  decades after explosion
- But... high X-ray luminosity of millisecond PWNe can re-ionize ejecta within months of optical peak (Ionization Break-Out)
- Pulsar wind  $\Rightarrow e^{+/-}$  pairs  $\Rightarrow$  X-rays  $\Rightarrow$  observer
- Luminous X-rays from SCP 06F consistent with ionization break-out  $\Rightarrow$  evidence that some SLSNe are engine-driven.