

Relativistic, under-dense outflows

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Outline

- Introduction
- Two-fluid simulations
- Analytic model of an electromagnetic precursor

Under-dense vs. over-dense

- Can the plasma screen out the electric field?
- Answer depends on the electron density and time available

$$\omega/\omega_{pe} \ll 1 \qquad \text{YES} - \textit{over-dense}$$

$$\omega/\omega_{pe} > 1 \qquad \text{NO} - \textit{under-dense}$$

- ... and on the amplitude $\omega_{pe} = (4\pi n_e e^2 / \langle \gamma \rangle m)^{1/2}$
- **Over-dense:** only MHD modes (magnetosonic, Alfvén), phase speeds subluminal
 ⇒ use ideal-MHD, force-free...
- **Under-dense:** high-frequency, transverse electromagnetic modes possible, with superluminal phase speed
 ⇒ use (at least) two-fluids.

Under-dense zones in a conical e^\pm jet/beam

Three dimensionless jet parameters:

- 1 (Mass-loading)⁻¹ $\mu = L/\dot{M}c^2$
- 2 Magnetization $\sigma_0 = \text{Poynting flux}/\text{K.E. flux}$
- 3 *A parameter describing the jet composition: e/m*
 - Cross-jet potential $\times e/mc^2$: $a_0 = eBr/mc^2$
 - (Dimensionless luminosity/unit solid angle)^{1/2}:

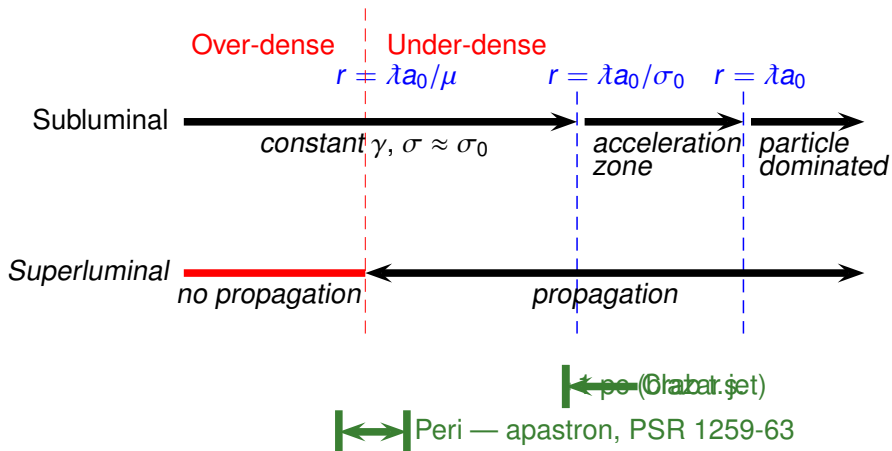
$$a_0 = (4\pi L/\Omega_s)^{1/2} (e^2/m^2c^5)^{1/2}$$

Constraints/Estimates:

- 1 $a_0 = 3.4 \times 10^{14} \sqrt{4\pi L_{46}/\Omega_s}$
- 2 $\sigma_0 \lesssim \mu^{2/3}$ (for a supermagnetosonic jet)
- 3 Pair multiplicity $\kappa_0 = a_0/(4\mu) > 1$

Under-dense zones in a conical jet

Fluctuation wavelength λ , $a_0 \gg \mu \gg \sigma \gg 1$

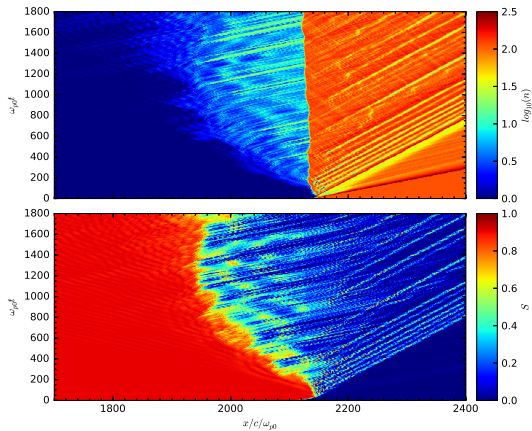


Two-fluid simulations

Simplest description that includes superluminal, electromagnetic modes is one with two charged fluids [Amano & Kirk ApJ \(2013\)](#)

- Relativistic, finite temperature electron & positron fluids
- 1D in space, 3D in momentum and EM fields
- Initial conditions:
 - Left half: circularly polarized, cold, static shear, $\gamma = 40$, $\sigma = 10$, $\lambda \approx \lambda_g/4$
 - Right half: shocked (R-H conditions) unmagnetized plasma

Time evolution

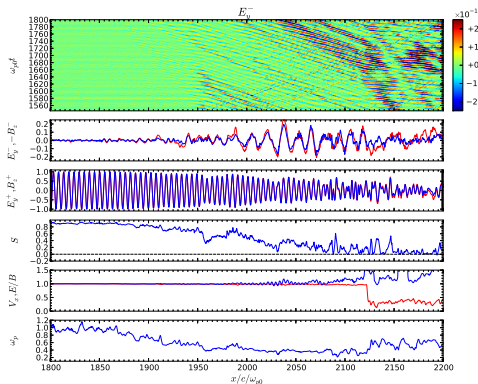


$$\Gamma = 40$$

$$\sigma = 10$$

$$\omega = 1.2\omega_p$$

Wave helicity



Positive helicity injected wave (E^+ , B^+).

Backwards propagating, negative helicity waves generated.

$E > B$ in precursor and downstream ($v_{\text{wave}} = B/E$).

Simulation Results and Implications

- Poynting flux dissipated completely
- A precursor containing strong electromagnetic waves is formed
- A hydrodynamic shock remains
- Particles are not tied to magnetic field lines in the downstream medium

Radiation-damped precursor

Analytical model of wave deceleration by radiation damping

Mochol & Kirk *ApJ* (2013).

Radiation mechanisms:

- synchro-Compton radiation (but with $\gamma \sim a$)
 - $\Omega_{\text{pulsar}} \rightarrow eB_{\text{eff}}/\gamma mc$
 - Frequency $\propto \gamma^2 B_{\text{eff}} \propto \gamma^3 \propto 1/r^3$
 - Power $\propto \gamma^2 B_{\text{eff}}^2 \propto \gamma^4 \propto 1/r^4$
- enhanced inverse-Compton radiation, in the presence of target photons
 - Frequency $\propto \gamma^2 v_{\text{target}} \propto 1/r^2$
 - Power $\propto \gamma^2 U_{\text{target}} \propto 1/r^2$

⇒ Gamma-ray binaries

Be star-pulsar binary

X-rays from interacting winds
Tavani & Arons (1997)

Be star wind:

$$\dot{M} \sim 10^{-8} M_{\odot} \text{ yr}^{-1}$$

$$v_{\text{wind}} \sim 10^3 \text{ km s}^{-1}$$

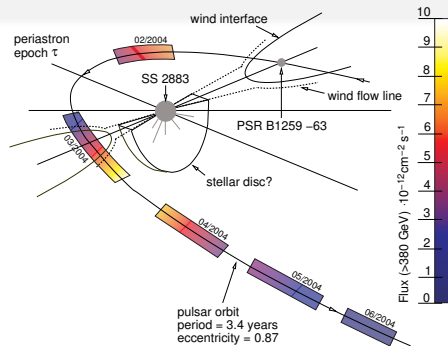
Pulsar wind:

$$L_{\text{s.d.}} \approx 8 \times 10^{35} \text{ erg s}^{-1}$$

Momentum balance:

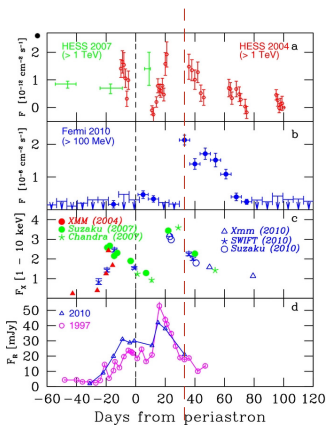
$$\frac{r_{\text{Be}}}{r_{\text{p}}} = \sqrt{\frac{L_{\text{s.d.}}}{\dot{M} v_{\text{wind}} c}} \sim 0.7$$

But $L_{\text{s.d.}} \gg L_{\text{Be wind}}$



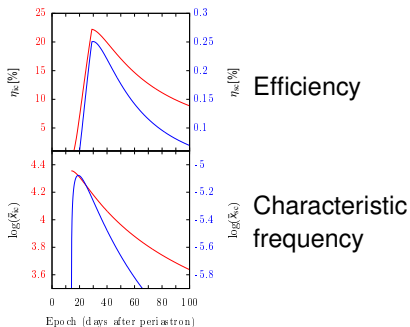
Based on cartoon by O. de Jager

GeV flare from PSR B1259-63



Abdo et al ApJL 2011

Stationary EM precursor



Inverse Compton
Synchro-Compton

Conclusions

- Under-dense flows present in pulsar winds and possibly also in AGN jets
- Their interactions with obstacles differ substantially from those of MHD flows.
- An observable signature is predicted from “thermal” particles in an electromagnetic precursor (gamma-ray binaries)

Future work: particle acceleration properties. . .