

## Physics 42200 Waves & Oscillations

Lecture 6 – French, Chapter 3

Spring 2016 Semester

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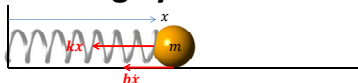
### Damping

- Most oscillating physical systems dissipate their energy over time
- We will consider the special cases where the force is a function of velocity

$$F = -b_1 v - b_2 v^2$$

- The drag force is in the opposite direction of the velocity
- Typical of an object moving through a fluid
  - Moving quickly through air: turbulent drag ( $b_2 v^2$  is important)
  - Moving slowly through water: viscous drag ( $b_1 v$  is important)
- When  $v$  is small enough, or  $b_2$  is small enough, only the first term is important.

### Oscillating System with Drag



- Newton's second law:

$$m\ddot{x} = -kx - b\dot{x}$$

$$m\ddot{x} + b\dot{x} + kx = 0$$

- What is the solution to this differential equation?
- Let's try this function:  $x(t) = Ae^{\alpha t}$ 
  - then  $\dot{x}(t) = \alpha Ae^{\alpha t}$  and  $\ddot{x}(t) = \alpha^2 Ae^{\alpha t}$
- Substitute it into the differential equation:
 
$$(\alpha^2 m + \alpha b + k)Ae^{\alpha t} = 0$$

## Oscillating Systems with Drag

$$(\alpha^2 m + \alpha b + k)Ae^{\alpha t} = 0$$

This is true for any value of  $t$  only when

$$\alpha^2 m + \alpha b + k = 0$$

Use the quadratic formula:

$$\alpha = \frac{-b \pm \sqrt{b^2 - 4km}}{2m} = -\frac{b}{2m} \pm \sqrt{\frac{b^2}{4m^2} - \frac{k}{m}}$$

which we write as

$$\alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

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## Oscillating Systems with Damping

$$m\ddot{x} + b\dot{x} + kx = 0$$

$$x(t) = \mathbf{A}e^{-\frac{\gamma}{2}t}e^{t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}} + \mathbf{B}e^{-\frac{\gamma}{2}t}e^{-t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}}$$

Three cases to consider:

1.  $\frac{\gamma^2}{4} - (\omega_0)^2 > 0$       Real roots
2.  $\frac{\gamma^2}{4} - (\omega_0)^2 < 0$       Imaginary roots
3.  $\frac{\gamma^2}{4} - (\omega_0)^2 = 0$       Degenerate roots

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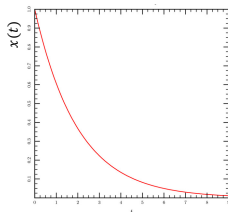
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## Oscillating Systems with Damping

$$x(t) = Ae^{-\frac{\gamma}{2}t}e^{t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}} + Be^{-\frac{\gamma}{2}t}e^{-t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}}$$

When  $\frac{\gamma^2}{4} - (\omega_0)^2 > 0$  both exponents are real and negative



The mass does not oscillate

It gradually approaches the equilibrium position at  $x = 0$ .

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## Oscillating Systems with Damping

$$x(t) = Ae^{-\frac{\gamma}{2}t} e^{it\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}} + Be^{-\frac{\gamma}{2}t} e^{-it\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}}$$

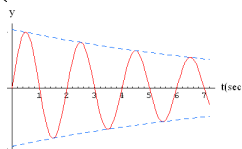
When  $\frac{\gamma^2}{4} - (\omega_0)^2 < 0$  we can write:

$$x(t) = Ae^{-\frac{\gamma}{2}t} e^{it\sqrt{(\omega_0)^2 - \frac{\gamma^2}{4}}} + Be^{-\frac{\gamma}{2}t} e^{-it\sqrt{(\omega_0)^2 - \frac{\gamma^2}{4}}} \\ = Ce^{-\frac{\gamma}{2}t} \cos(\omega t + \varphi)$$

where  $\omega = \sqrt{(\omega_0)^2 - \frac{\gamma^2}{4}}$

Oscillation frequency gradually decreases.

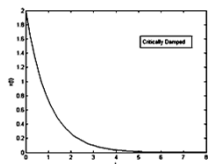
Oscillation frequency is slightly lower than the un-damped oscillator.



## Oscillating Systems with Damping

- Third possibility:  $\frac{\gamma}{2} - \omega_0 = 0$
- In this case the solution is slightly different:

$$x(t) = (A + Bt)e^{-\frac{\gamma}{2}t}$$



fastest return to equilibrium position without oscillating.

## Notation

- We just introduced a lot of notation:

$$\omega_0 = \sqrt{k/m}$$

$$\gamma = \frac{b}{m}$$

$$\omega = \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

- None of these express a "fundamental physical law"
- They are just definitions
- If we had defined  $\gamma' = \frac{b}{2m}$  then we could write
 
$$x(t) = Ae^{-\gamma't} e^{it\sqrt{(\gamma')^2 - (\omega_0)^2}} + Be^{-\gamma't} e^{-it\sqrt{(\gamma')^2 - (\omega_0)^2}}$$
- Maybe this is simpler or easier to remember, but it is still rather arbitrary and chosen only for convenience.

### Suggestion

- Do not memorize these formulas...
- Instead, memorize this procedure:

$$m\ddot{x} + b\dot{x} + kx = 0$$

$$x(t) = Ae^{\alpha t}$$

$$(\alpha^2 m + \alpha b + k)Ae^{\alpha t} = 0$$

$$\alpha^2 m + \alpha b + k = 0$$

$$\alpha = \frac{-b \pm \sqrt{b^2 - 4km}}{2m} = -\frac{b}{2m} \pm \sqrt{\frac{b^2}{4m^2} - \frac{k}{m}}$$

- Then define any new variables you introduce:

$$x(t) = Ce^{-\frac{b}{2m}t} \cos(\omega t + \phi) \text{ (when } \frac{b^2}{4m^2} - \frac{k}{m} < 0)$$

- Recognize that there are three possible forms of the solution.

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### Example

- Suppose a 1 kg mass oscillates with frequency  $f$  and the amplitude of oscillations decreases by a factor of  $\frac{1}{2}$  in time  $T$ . What differential equation describes the motion?

$$m\ddot{x} + b\dot{x} + kx = 0$$

$$m = 1 \text{ kg}$$

$$b = \frac{2m \log 2}{T}$$

$$k = m \left( 4\pi^2 f^2 + \left( \frac{\log 2}{T} \right)^2 \right)$$

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### RLC Circuits

Review of electricity and magnetism:

- Capacitors store energy in an electric field
- Resistors dissipate energy by heating
- Inductors store energy in a magnetic field
- Voltage
  - Energy per unit charge
  - SI units: **Volt** = J/C
- Current
  - Charge passing a point in a circuit per unit time
  - SI units: **Ampere** = C/s

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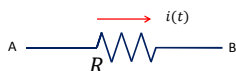
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## Resistors



- Current flows from A to B in the direction indicated by the arrow.
- Charges at point B have less energy than at point A because some of their energy was dissipated as heat.
- Potential difference:  

$$\Delta V = V_A - V_B = i(t) R$$
- Resistance,  $R$ , is measured in *ohms* in SI units

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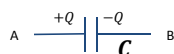
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## Capacitors



- Potential difference:  

$$\Delta V = V_A - V_B = \frac{Q}{C}$$
- No current can flow across the capacitor, so any charge that flows onto the plates accumulates there:  

$$Q(t) = Q_0 + \int_0^t i(t) dt$$
- Potential difference:  

$$\Delta V = \frac{1}{C} \int_0^t i(t) dt$$
- SI units for capacitance: *farad*

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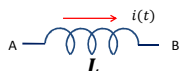
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## Inductors



- The inductor will establish a potential difference that opposes any change in current.
  - If  $di/dt < 0$  then  $V_B > V_A$ 
    - magnetic field is being converted to energy
  - If  $di/dt > 0$  then  $V_B < V_A$ 
    - energy is being stored in the magnetic field

$$\Delta V = V_A - V_B = L \frac{di}{dt}$$

- SI units for inductance: *henry*

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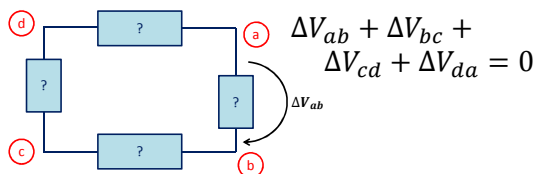
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### Kirchhoff's Loop Rule

- The sum of the potential differences around a loop in a circuit must equal zero:




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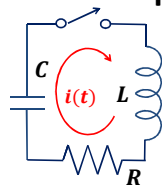
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### Kirchhoff's Loop Rule



Sum of potential differences:

$$-L \frac{di}{dt}$$

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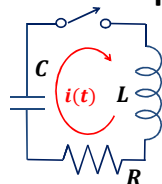
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### Kirchhoff's Loop Rule



Sum of potential differences:

$$-L \frac{di}{dt} - i(t)R$$

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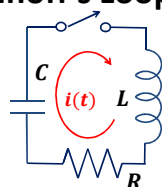
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**Kirchhoff's Loop Rule**

Sum of potential differences:

$$-L \frac{di}{dt} - i(t)R - \frac{1}{C} \left( Q_0 + \int_0^t i(t) dt \right) = 0$$

Initial charge,  $Q_0$ , defines the initial conditions.

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**Kirchhoff's Loop Rule**

$$L \frac{di}{dt} + i(t)R + \frac{1}{C} \left( Q_0 + \int_0^t i(t) dt \right) = 0$$

Differentiate once with respect to time:

$$L \frac{d^2i}{dt^2} + R \frac{di}{dt} + \frac{1}{C} i(t) = 0$$

This is of the same form as the equation for a damped harmonic oscillator:

$$m \frac{d^2x}{dt^2} + b \frac{dx}{dt} + kx(t) = 0$$

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**Solutions**

$$L \frac{d^2i}{dt^2} + R \frac{di}{dt} + \frac{1}{C} i(t) = 0$$

Suppose  $i(t) = Ae^{\alpha t}$

Then  $\frac{di}{dt} = \alpha Ae^{\alpha t}$  and  $\frac{d^2i}{dt^2} = \alpha^2 Ae^{\alpha t}$

Substitute into the differential equation:

$$\left( \alpha^2 L + \alpha R + \frac{1}{C} \right) Ae^{\alpha t} = 0$$

True for any  $t$  only if  $\alpha^2 L + \alpha R + \frac{1}{C} = 0$ .

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## Solutions

$$\alpha^2 L + \alpha R + \frac{1}{C} = 0$$

Roots of the polynomial:

$$\alpha = -\frac{R}{2L} \pm \sqrt{\frac{R^2}{4L^2} - \frac{1}{LC}}$$

Define some new symbols:

$$\omega_0 = \sqrt{1/LC}$$

$$\gamma = R/L$$

Then the roots can be written

$$\alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

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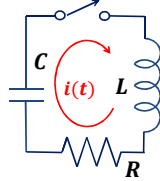
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## Example



- If  $L = 2.2 \mu\text{H}$  and  $C = 10 \text{ nF}$  what is the frequency of oscillations when  $R = 0$ ?
- What is the largest value of  $R$  that will still allow the circuit to oscillate?

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## Example

$$\omega_0 = \sqrt{1/LC} = \sqrt{\frac{1}{(2.2 \mu\text{H})(10 \text{ nF})}} = 6.74 \times 10^6 \text{ s}^{-1} = 2.15 \text{ MHz}$$

$$\gamma = R/L$$

$$\alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

- Critical damping:

$$\frac{\gamma^2}{4} - (\omega_0)^2 = 0$$

$$\gamma = \frac{R}{L} = 2 \sqrt{\frac{1}{LC}}$$

$$R = 2 \sqrt{\frac{L}{C}} = 2 \sqrt{\frac{2.2 \mu\text{H}}{10 \text{ nF}}} = 29.7 \Omega$$

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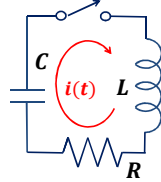
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**Example**

$$\begin{aligned} L &= 2.2 \mu H \\ C &= 10 \text{ nF} \\ R &= 2 \Omega \end{aligned}$$

- Suppose the initial charge on the capacitor was 10 nC... What voltage is measured across  $R$  as a function of time?

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**Example**

- Calculate the discriminant:

$$\frac{R^2}{4L^2} - \frac{1}{LC} = \frac{(2 \Omega)^2}{4(2.2 \mu H)^2} - \frac{1}{(2.2 \mu H)(0.01 \mu F)} < 0$$

- The circuit will oscillate with frequency:

$$f = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}} = 1.07 \text{ MHz}$$

- Time constant:

$$\frac{\gamma}{2} = \frac{R}{2L} = \frac{(2 \Omega)}{2(2.2 \mu H)} = 4.55 \times 10^5 \text{ s}^{-1}$$

- Current will be:

$$i(t) = i_0 e^{-\gamma t/2} \cos(\omega t + \varphi)$$

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**Example**

$$i(t) = i_0 e^{-\gamma t/2} \cos(\omega t + \varphi)$$

- Initial conditions:

➤  $i(0) = 0$  because the inductor produces a potential difference that opposes the change in current.

➤ Therefore,  $\varphi = \pi/2$

➤ Initial potential across capacitor:

$$\Delta V = \frac{Q}{C} = \frac{10 \text{ nC}}{10 \text{ nF}} = 1 \text{ Volt}$$

➤ Initial voltage across inductor:

$$\Delta V = L \frac{di}{dt} = 1 \text{ Volt}$$

$$\frac{di}{dt} = \frac{1 \text{ V}}{2.2 \mu H} = 4.55 \times 10^5 \text{ A/s} = i_0 \omega$$

➤  $i_0 = \frac{4.55 \times 10^5 \text{ A/s}}{6.73 \times 10^6 \text{ /s}} = 68 \text{ mA}$

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### Example

- Current in circuit:

$$i(t) = i_0 e^{-\gamma t/2} \sin \omega t$$

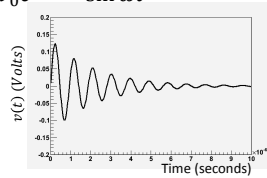
where  $\gamma/2 = 4.55 \times 10^5 \text{ s}^{-1}$  and  $i_0 = 68 \text{ mA}$

- Potential difference across the resistor:

$$\Delta V = i(t)R$$

$$v(t) = v_0 e^{-\gamma t/2} \sin \omega t$$

where  $v_0 = 136 \text{ mV}$ .




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