

Physics 42200 Waves & Oscillations

Lecture 2 – French, Chapter 1

Spring 2016 Semester Matthew Jones

- Mass-spring system:
 - Force given by Hooke's law:

$$F = -kx$$

– Newton's second law:

$$F = ma = m\ddot{x} = m\frac{d^2x}{dt^2}$$

– Equation of motion:

$$m\ddot{x} + kx = 0$$
$$\ddot{x} + \omega^2 x = 0$$
where $\omega = \sqrt{k/m}$

Differential equation:

$$\frac{d^2x}{dt^2} + \omega^2 x = 0$$

Solutions can be written in various ways:

$$x(t) = A \cos(\omega t + \varphi)$$

$$x(t) = A \sin \omega t + B \cos \omega t$$
(and many others...)

 Two constants of integration need to be determined from initial conditions or other information.

- How do we know that these are solutions?
- Compute the derivatives:

$$x(t) = A \cos(\omega t + \varphi)$$

$$\dot{x}(t) = -A\omega \sin(\omega t + \varphi)$$

$$\ddot{x}(t) = -A\omega^2 \cos(\omega t + \varphi) = -\omega^2 x(t)$$

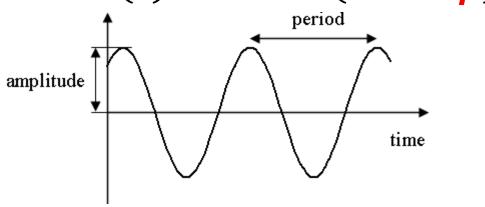
Substitute into the differential equation:

$$\ddot{x} + \omega^2 x = \left(-\omega^2 x(t)\right) + \omega^2 x(t) = 0$$

- Mathematical details:
 - How could we deduce that this was a solution if we didn't already know it was?

Properties of the solution:

$$x(t) = \mathbf{A} \, \cos(\boldsymbol{\omega}t + \boldsymbol{\varphi})$$



Notation: Amplitude: A

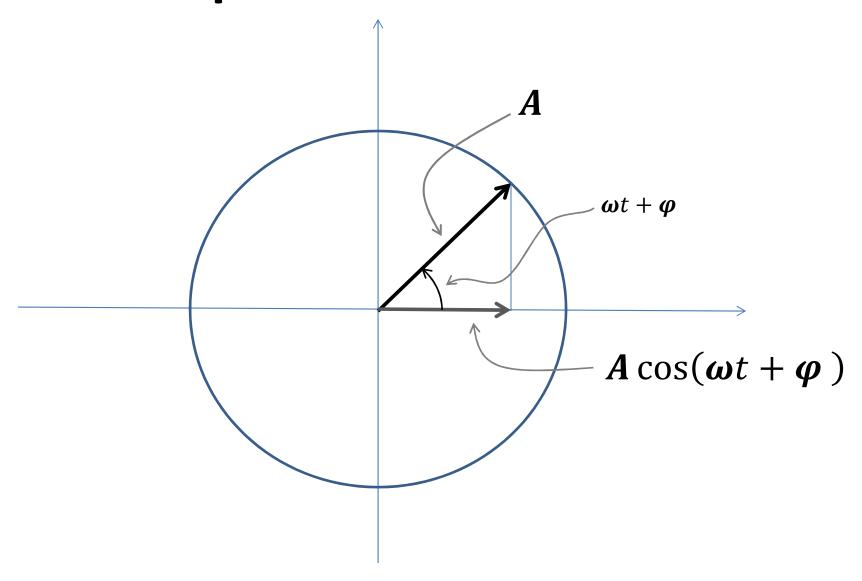
Initial phase: ϕ

Angular frequency: ••

Frequency: $f = \omega/2\pi$

Period: $T = 1/f = 2\pi/\omega$

Descriptions of Harmonic Motion

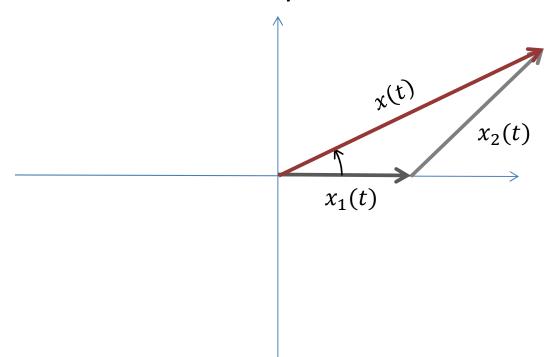


Uniform Circular Motion

- The graph of $x(t) = A \cos(\omega t + \varphi)$ is the same as the projection onto the x-axis of a vector of length A, rotating with angular frequency ω .
- This is a useful geometric description of the motion.
 - Two-component vectors are introduced only for convenience (we call them "phasors")
 - The solution we are interested in is just the projection onto the x-axis.
- Example...

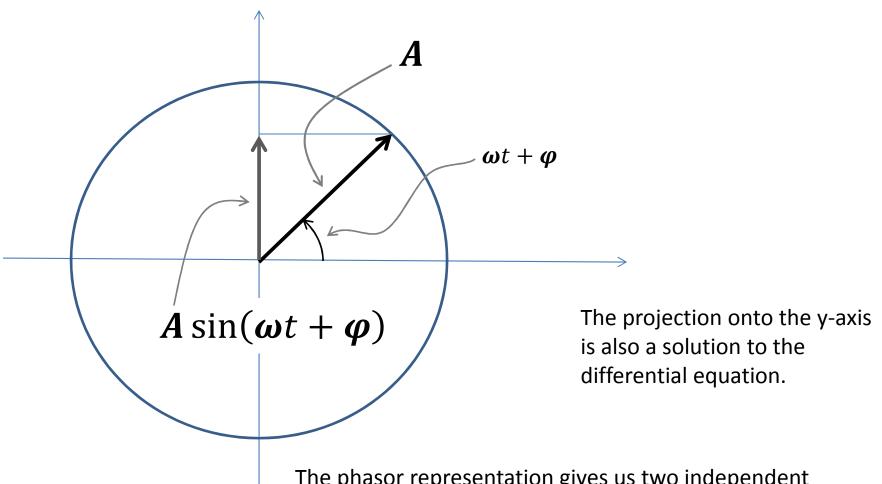
Uniform Circular Motion

- The differential equation $\ddot{x} + \omega^2 x = 0$ is linear:
 - Suppose $x_1(t)$ and $x_2(t)$ are both solutions
 - Then the function $x(t) = a x_1(t) + b x_2(t)$ is also a solution for any real numbers a and b.



Actually, the functions $x_1(t)$, $x_2(t)$ and x(t) are the projections of these vectors onto the x-axis.

Uniform Circular Motion



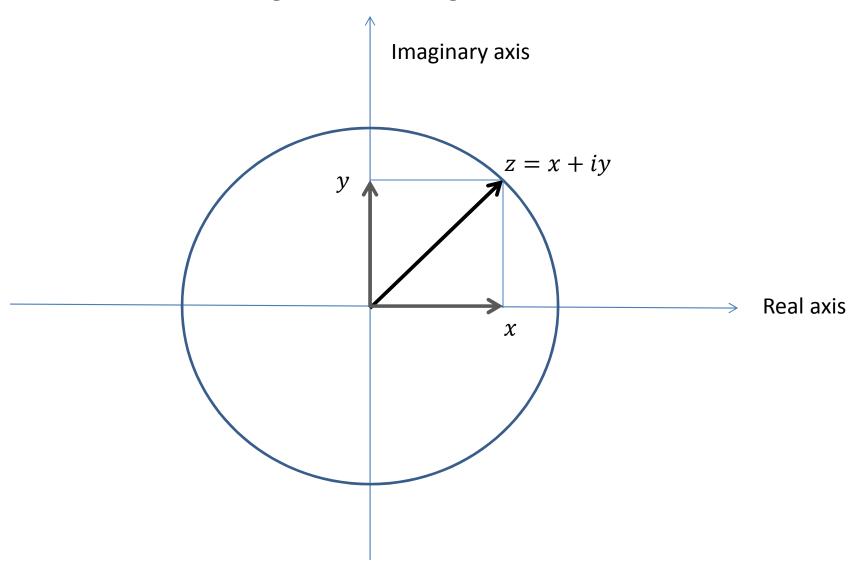
The phasor representation gives us two independent solutions, even though we might only want to use only one of them to describe the motion.

Phasor Representation

- The phasor provides all the information we need to describe the motion
 - If we just knew the value of x at one time t, we still don't know what A and φ are.
 - But if we know x and y at time t then we have enough information to calculate both A and φ .
- The more general description of the motion can be useful for analyzing problems even if the "physical" solution to the equations of motion is just one of its projections.

Basic definitions:

$$i=\sqrt{-1}$$
 $i^2=-1$
 $z=x+iy$
 $z^*=x-iy$
 $|z|=\sqrt{z^*z}=\sqrt{x^2+y^2}$
 $Re(z)=x=(z+z^*)/2$
 $Im(z)=y=(z-z^*)/2i$
(where x and y are real numbers)



- But complex numbers are way better...
- Euler's identity:

$$e^{i\theta} = \cos\theta + i\sin\theta$$

Complex numbers in this form satisfy:

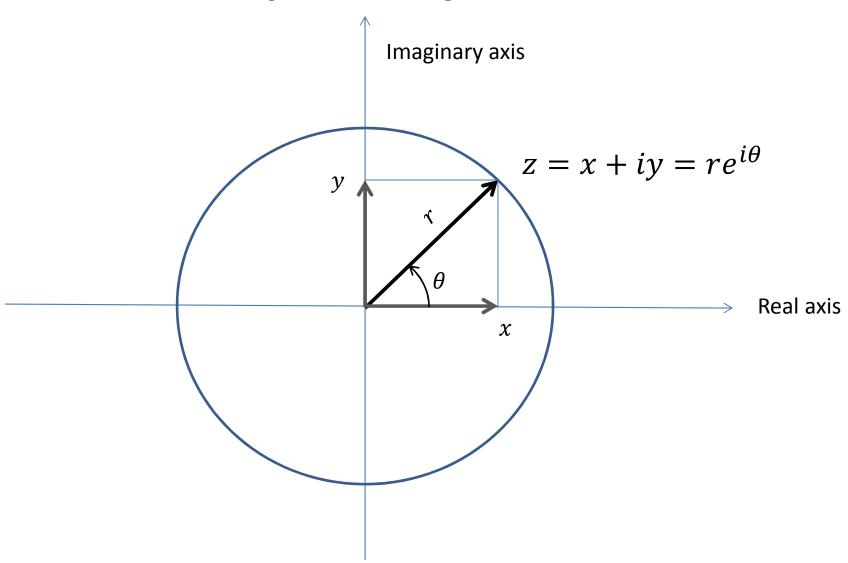
$$(e^{i\theta})^* = e^{-i\theta}$$

$$e^{i\theta}e^{-i\theta} = (\cos\theta + i\sin\theta)(\cos\theta - i\sin\theta)$$

$$= \cos^2\theta + \sin^2\theta = 1$$

In general, we can always write

$$z = x + iy = re^{i\theta}$$



The other way we will describe solutions to

$$\ddot{x} + \omega^2 x = 0$$

will be using complex numbers...

$$-\operatorname{Let} x(t) = \mathbf{r}e^{i(\omega t + \mathbf{\varphi})} = (\mathbf{r}e^{i\mathbf{\varphi}})e^{i\omega t} = \mathbf{c}e^{i\omega t}$$

This part is just a constant.

– Derivatives are:

$$\dot{x}(t) = i\omega \, \mathbf{c}e^{i\omega t}$$
$$\ddot{x}(t) = (i\omega)^2 \mathbf{c}e^{i\omega t} = -\omega^2 x(t)$$

It is a solution:

$$\ddot{x} + \omega^2 x = \left(-\omega^2 x(t)\right) + \omega^2 x(t) = 0$$

- The *physical* displacement of the mass must be a real number.
- The displacement as a function of time is given by the real component, Re[x(t)].
- The complex representation contains more information than is present in just the function describing the physical displacement.
 - It provides both amplitude and phase information

A solution to the differential equation

$$\frac{d^2x}{dt^2} + \omega^2 x = 0$$

can be written

$$x(t) = \mathbf{A} \cos(\omega t + \mathbf{\varphi})$$

- We need more information if have to also determine the constants of integration A and φ .
- This information is usually given as *initial conditions* (or *boundary conditions*).

Example:

"A mass of 100 g is attached to a spring with spring constant of 1 N/m that is initially stretched to a length of 2 cm. What is the solution to the equation of motion if it is released from rest at time t=0?"

- 1. Analyze the problem in general don't use the numbers yet. If you make a mistake you will never find it in a big mess of numbers.
- Let m = 100 g be the mass, k = 1 N/m be the spring constant and x_0 be the initial displacement.
- 2. Introduce other symbols if it is convenient to do so:

Let
$$\omega = \sqrt{k/m}$$
.

3. Now describe the system in terms of the relevant physical principles...

The force acting on the mass is described by Hooke's law:

$$F(x) = -k x$$

and the resulting motion is given by Newton's second law:

$$m \frac{d^2x}{dt^2} = -k x...$$

4. Write it in a "standard form" where you know what the solution is:

... which can be written

$$\frac{d^2x}{dt^2} + \omega^2 x = 0$$

and has solutions of the form,

$$x(t) = A \cos(\omega t + \varphi)$$

where A and φ are constants of integration that remain to be determined.

5. Calculate the initial conditions at t=0:

At t = 0, the initial displacement is $x(0) = A \cos \varphi = x_0$

and the initial velocity is

$$\dot{x}(0) = \frac{dx}{dt} \Big|_{t=0} = -A\omega \sin \varphi = 0.$$

6. What values of A and φ will satisfy these two equations?

$$x(0) = A\cos\varphi = x_0$$

$$\dot{x}(0) = \frac{dx}{dt}\Big|_{t=0} = -A\omega\sin\varphi = 0$$

- $A \neq 0$ and $\omega \neq 0$ so we must have $\varphi = 0$.
- Then $A = x_0$.

The initial conditions are satisfied by $A=x_0$ and $\varphi=0$, so the solution to the initial value problem is

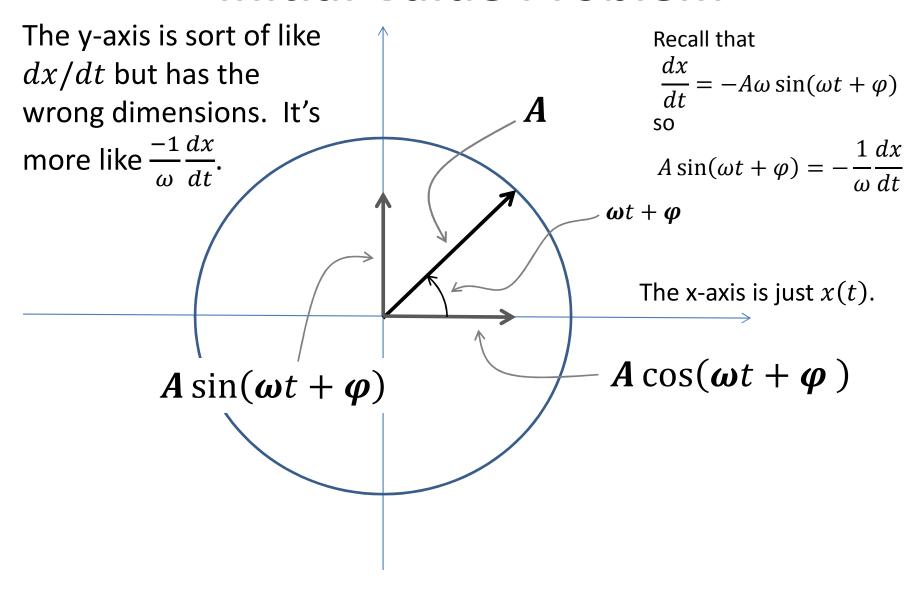
$$x(t) = x_0 \cos \omega t$$

Now you can substitute in the numbers if you want...

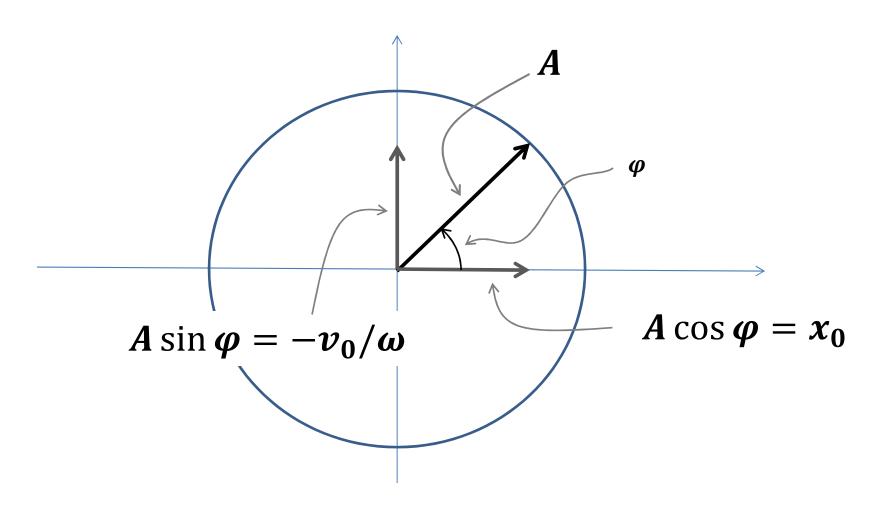
$$\omega = \sqrt{\frac{k}{m}} = \sqrt{\frac{1 N/m}{0.1 kg}} = \sqrt{\frac{1 [kg \cdot m/s^2]/m}{0.1 kg}}$$
$$= 3.162 s^{-1}$$
$$f = \omega/2\pi = 0.503 s^{-1}$$
$$x(t) = (2 cm) \cos[(3.162 s^{-1}) \times t]$$

Make sure you write the units everywhere. That way it is clear that x(t) is in cm and t is in seconds.

- What if the initial conditions were more complicated?
 - Initial displacement: x_0
 - Initial velocity: v_0
- How can we determine the constants of integration?
 - You could use lots of trigonometric identities
 - You could look at the problem in a geometric way to see the solution...



At t=0, the x-component is $x(0)=x_0$ and the y-component is $-\frac{\dot{x}(0)}{\omega}=-v_0/\omega$.



• Solve for A and φ :

$$\frac{A \sin \varphi}{A \cos \varphi} = \tan \varphi = \frac{-v_0}{\omega x_0}$$

$$\varphi = \tan^{-1} \left(\frac{-v_0}{\omega x_0}\right)$$

$$A^2 \sin^2 \varphi + A^2 \cos^2 \varphi = A^2 = x_0^2 + v_0^2/\omega^2$$

$$A = \pm \sqrt{x_0^2 + \frac{v_0^2}{\omega^2}}$$

You still need the initial conditions to get the right sign.