

Physics 42200
Waves & Oscillations

Lecture 6 – French, Chapter 3

Spring 2015 Semester

Matthew Jones

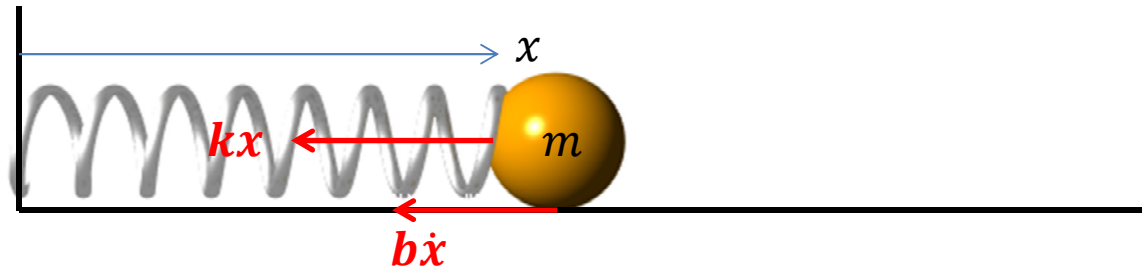
Damping

- Most oscillating physical systems dissipate their energy over time
- We will consider the special cases where the force is a function of velocity

$$F = -b_1 v - b_2 v^2$$

- The drag force is in the opposite direction of the velocity
- Typical of an object moving through a fluid
 - Moving quickly through air: turbulent drag ($b_2 v^2$ is important)
 - Moving slowly through water: viscous drag ($b_1 v$ is important)
- When v is small enough, or b_2 is small enough, only the first term is important.

Oscillating System with Drag



- Newton's second law:

$$m\ddot{x} = -kx - b\dot{x}$$
$$m\ddot{x} + b\dot{x} + kx = 0$$

- What is the solution to this differential equation?
- Let's try this function: $x(t) = Ae^{\alpha t}$
 - then $\dot{x}(t) = \alpha Ae^{\alpha t}$ and $\ddot{x}(t) = \alpha^2 Ae^{\alpha t}$
- Substitute it into the differential equation:

$$(\alpha^2 m + \alpha b + k)Ae^{\alpha t} = 0$$

Oscillating Systems with Drag

$$(\alpha^2 m + \alpha b + k)Ae^{\alpha t} = 0$$

This is true for any value of t only when

$$\alpha^2 m + \alpha b + k = 0$$

Use the quadratic formula:

$$\alpha = \frac{-b \pm \sqrt{b^2 - 4km}}{2m} = -\frac{b}{2m} \pm \sqrt{\frac{b^2}{4m^2} - \frac{k}{m}}$$

which we write as

$$\alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

Oscillating Systems with Damping

$$m\ddot{x} + b\dot{x} + kx = 0$$

$$x(t) = \mathbf{A}e^{-\frac{\gamma}{2}t}e^{t\sqrt{\gamma^2/4-(\omega_0)^2}} + \mathbf{B}e^{-\frac{\gamma}{2}t}e^{-t\sqrt{\gamma^2/4-(\omega_0)^2}}$$

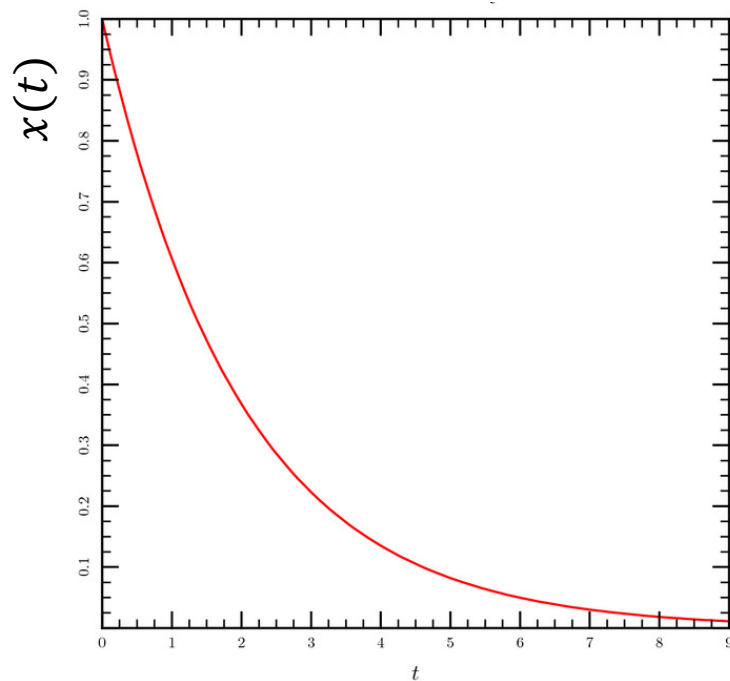
Three cases to consider:

1. $\frac{\gamma^2}{4} - (\omega_0)^2 > 0$ Real roots
2. $\frac{\gamma^2}{4} - (\omega_0)^2 < 0$ Imaginary roots
3. $\frac{\gamma^2}{4} - (\omega_0)^2 = 0$ Degenerate roots

Oscillating Systems with Damping

$$x(t) = Ae^{-\frac{\gamma}{2}t} e^{t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}} + Be^{-\frac{\gamma}{2}t} e^{-t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}}$$

When $\frac{\gamma^2}{4} - (\omega_0)^2 > 0$ both exponents are real and negative



The mass does not oscillate

It gradually approaches the equilibrium position at $x = 0$.

Oscillating Systems with Damping

$$x(t) = Ae^{-\frac{\gamma}{2}t} e^{t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}} + Be^{-\frac{\gamma}{2}t} e^{-t\sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}}$$

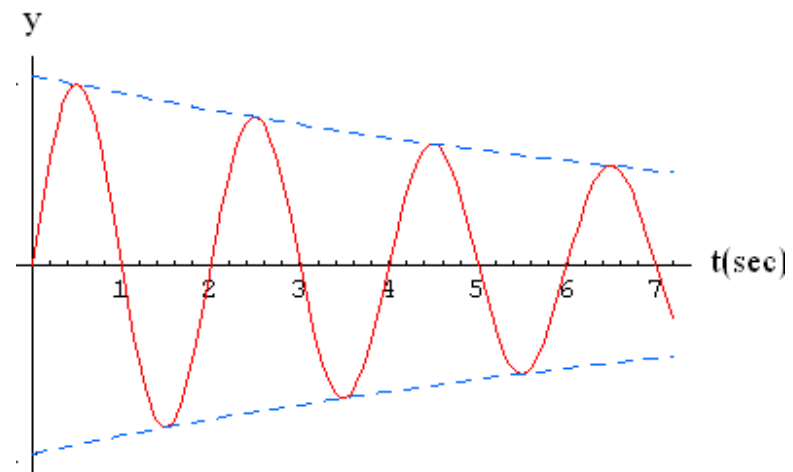
When $\frac{\gamma^2}{4} - (\omega_0)^2 < 0$ we can write:

$$\begin{aligned} x(t) &= Ae^{-\frac{\gamma}{2}t} e^{it\sqrt{(\omega_0)^2 - \frac{\gamma^2}{4}}} + Be^{-\frac{\gamma}{2}t} e^{-it\sqrt{(\omega_0)^2 - \frac{\gamma^2}{4}}} \\ &= Ce^{-\frac{\gamma}{2}t} \cos(\omega t + \varphi) \end{aligned}$$

where $\omega = \sqrt{(\omega_0)^2 - \frac{\gamma^2}{4}}$

Oscillation frequency gradually decreases.

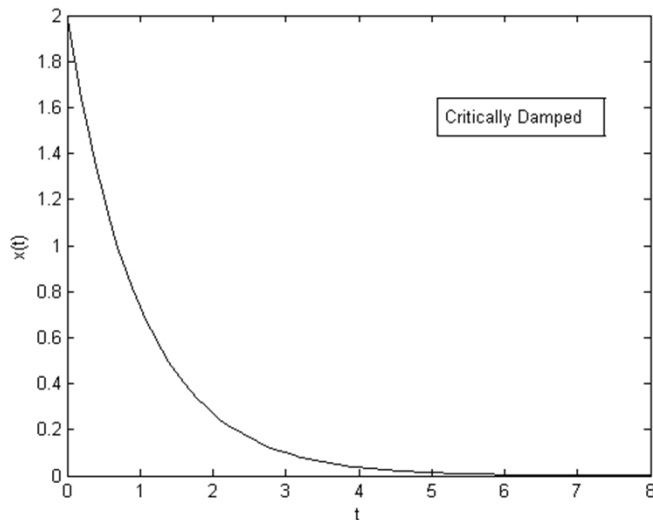
Oscillation frequency is slightly lower than the un-damped oscillator.



Oscillating Systems with Damping

- Third possibility: $\frac{\gamma}{2} - \omega_0 = 0$
- In this case the solution is slightly different:

$$x(t) = (A + Bt)e^{-\frac{\gamma}{2}t}$$



fastest return to equilibrium position without oscillating.

Notation

- We just introduced a lot of notation:

$$\omega_0 = \sqrt{k/m}$$

$$\gamma = \frac{b}{m}$$

$$\omega = \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

- None of these express a “fundamental physical law”
- They are just definitions

- If we had defined $\gamma' = \frac{b}{2m}$ then we could write

$$x(t) = Ae^{-\gamma't} e^{t\sqrt{(\gamma')^2 - (\omega_0)^2}} + Be^{-\gamma't} e^{-t\sqrt{(\gamma')^2 - (\omega_0)^2}}$$

- Maybe this is simpler or easier to remember, but it is still rather arbitrary and chosen only for convenience.

Suggestion

- Do not memorize these formulas...
- Instead, memorize this procedure:

$$m\ddot{x} + b\dot{x} + kx = 0$$

$$x(t) = Ae^{\alpha t}$$

$$(\alpha^2 m + \alpha b + k)Ae^{\alpha t} = 0$$

$$\alpha^2 m + \alpha b + k = 0$$

$$\alpha = \frac{-b \pm \sqrt{b^2 - 4km}}{2m} = -\frac{b}{2m} \pm \sqrt{\frac{b^2}{4m^2} - \frac{k}{m}}$$

- Then define any new variables you introduce:

$$x(t) = \textcolor{red}{C} e^{-\frac{\gamma}{2}t} \cos(\omega t + \textcolor{red}{\varphi}) \quad \left(\text{when } \frac{b^2}{4m^2} - \frac{k}{m} < 0\right)$$

- Recognize that there are three possible forms of the solution.

Example

- Suppose a 1 kg mass oscillates with frequency f and the amplitude of oscillations decreases by a factor of $\frac{1}{2}$ in time T . What differential equation describes the motion?

$$m\ddot{x} + b\dot{x} + kx = 0$$

$$m = 1 \text{ kg}$$

$$b = \frac{2m \log 2}{T}$$

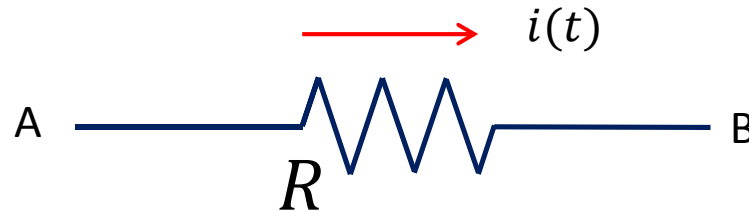
$$k = m \left(4\pi^2 f^2 + \left(\frac{\log 2}{T} \right)^2 \right)$$

RLC Circuits

Review of electricity and magnetism:

- Capacitors store energy in an electric field
- Resistors dissipate energy by heating
- Inductors store energy in a magnetic field
- Voltage
 - Energy per unit charge
 - SI units: *Volt = J/C*
- Current
 - Charge passing a point in a circuit per unit time
 - SI units: *Ampere = C/s*

Resistors

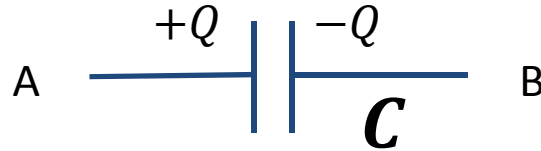


- Current flows from A to B in the direction indicated by the arrow.
- Charges at point B have less energy than at point A because some of their energy was dissipated as heat.
- Potential difference:

$$\Delta V = V_A - V_B = i(t) R$$

- Resistance, R , is measured in *ohms* in SI units

Capacitors



- Potential difference:

$$\Delta V = V_A - V_B = \frac{Q}{C}$$

- No current can flow across the capacitor, so any charge that flows onto the plates accumulates there:

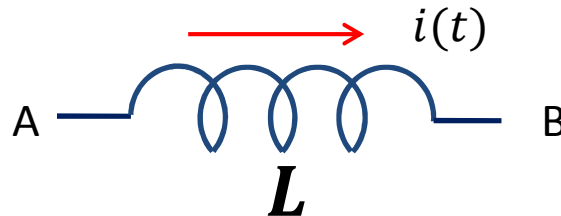
$$Q(t) = Q_0 + \int_0^t i(t) dt$$

- Potential difference:

$$\Delta V = \frac{1}{C} \int_0^t i(t) dt$$

- SI units for capacitance: *farad*

Inductors



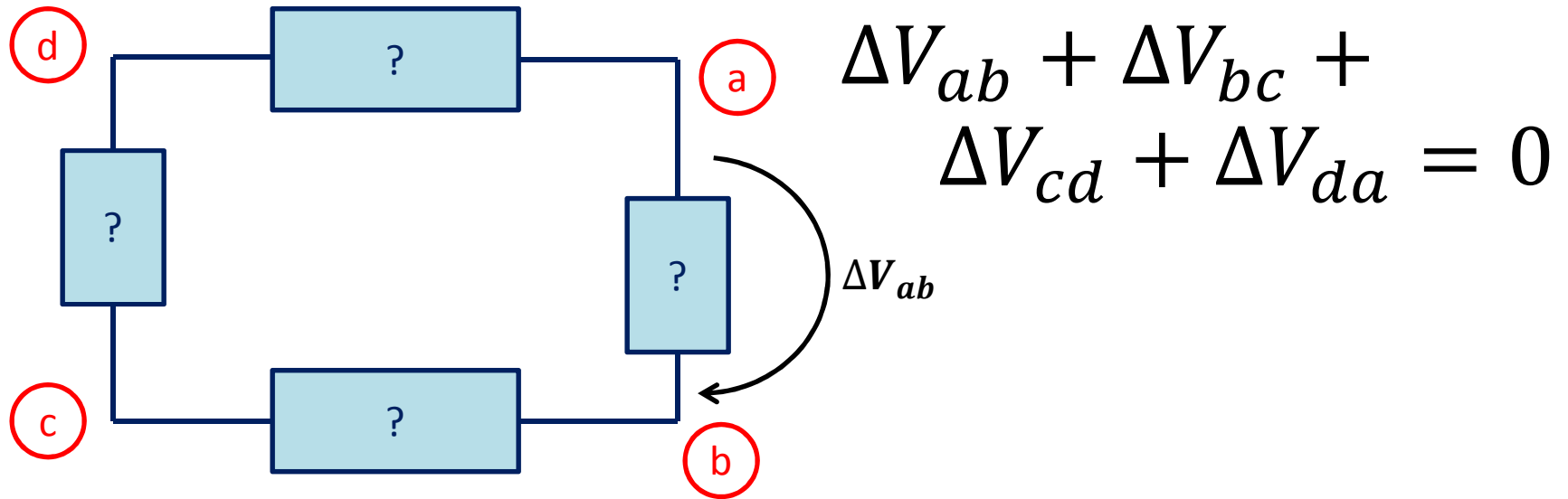
- The inductor will establish a potential difference that opposes any change in current.
 - If $di/dt < 0$ then $V_B > V_A$
 - magnetic field is being converted to energy
 - If $di/dt > 0$ then $V_B < V_A$
 - energy is being stored in the magnetic field

$$\Delta V = V_A - V_B = L \frac{di}{dt}$$

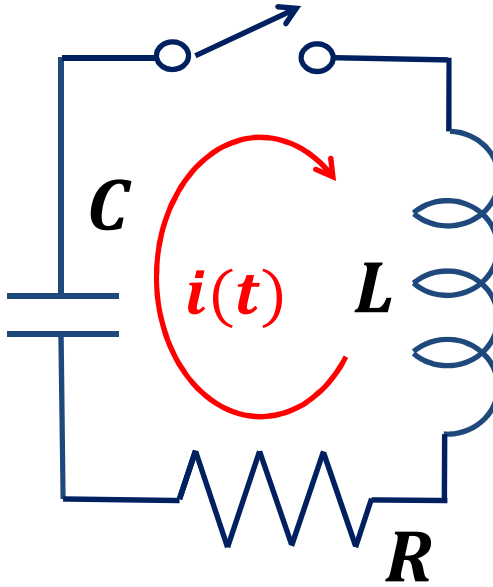
- SI units for inductance: *henry*

Kirchhoff's Loop Rule

- The sum of the potential differences around a loop in a circuit must equal zero:



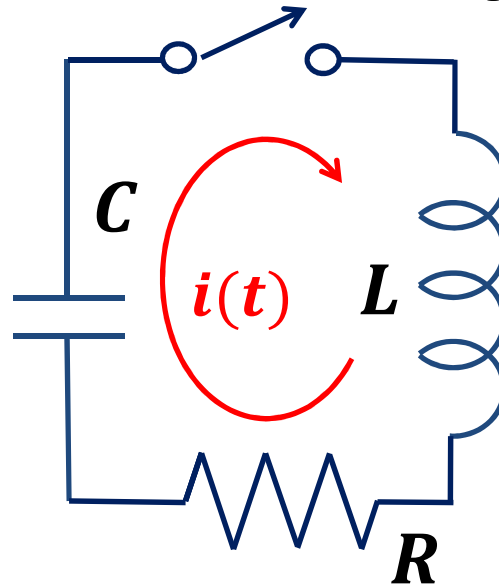
Kirchhoff's Loop Rule



Sum of potential differences:

$$-L \frac{di}{dt}$$

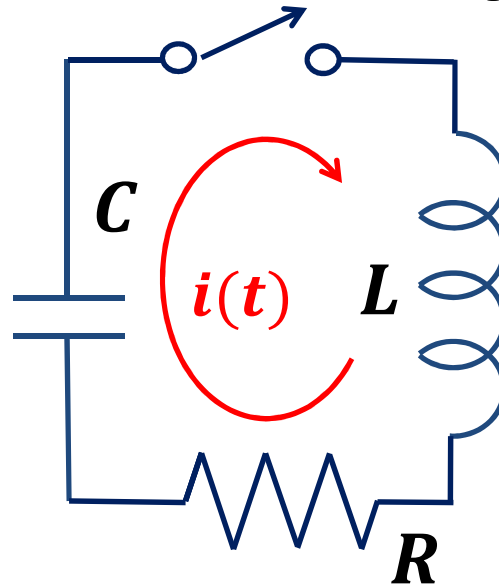
Kirchhoff's Loop Rule



Sum of potential differences:

$$-L \frac{di}{dt} - i(t)R$$

Kirchhoff's Loop Rule



Sum of potential differences:

$$-L \frac{di}{dt} - i(t)R - \frac{1}{C} \left(Q_0 + \int_0^t i(t) dt \right) = 0$$

Initial charge, Q_0 , defines the initial conditions.

Kirchhoff's Loop Rule

$$L \frac{di}{dt} + i(t)R + \frac{1}{C} \left(Q_0 + \int_0^t i(t) dt \right) = 0$$

Differentiate once with respect to time:

$$L \frac{d^2 i}{dt^2} + R \frac{di}{dt} + \frac{1}{C} i(t) = 0$$

This is of the same form as the equation for a damped harmonic oscillator:

$$m \frac{d^2 x}{dt^2} + b \frac{dx}{dt} + kx(t) = 0$$

Solutions

$$L \frac{d^2 i}{dt^2} + R \frac{di}{dt} + \frac{1}{C} i(t) = 0$$

Suppose $i(t) = Ae^{\alpha t}$

Then $\frac{di}{dt} = \alpha Ae^{\alpha t}$ and $\frac{d^2 i}{dt^2} = \alpha^2 Ae^{\alpha t}$

Substitute into the differential equation:

$$\left(\alpha^2 L + \alpha R + \frac{1}{C} \right) Ae^{\alpha t} = 0$$

True for any t only if $\alpha^2 L + \alpha R + \frac{1}{C} = 0$.

Solutions

$$\alpha^2 L + \alpha R + \frac{1}{C} = 0$$

Roots of the polynomial:

$$\alpha = -\frac{R}{2L} \pm \sqrt{\frac{R^2}{4L^2} - \frac{1}{LC}}$$

Define some new symbols:

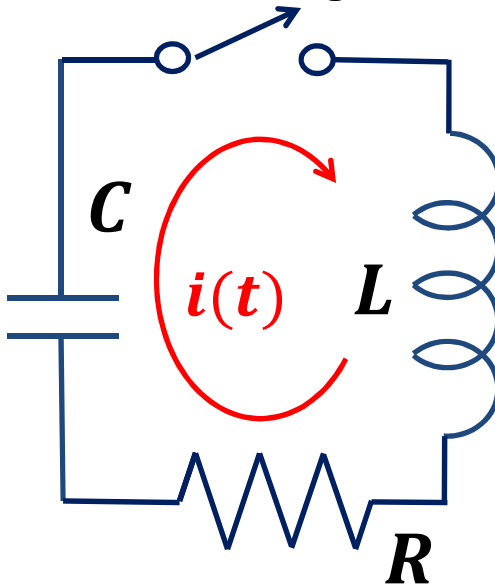
$$\omega_0 = \sqrt{1/LC}$$

$$\gamma = R/L$$

Then the roots can be written

$$\alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

Example



- If $L = 2.2 \mu H$ and $C = 10 nF$ what is the frequency of oscillations when $R = 0$?
- What is the largest value of R that will still allow the circuit to oscillate?

Example

$$\omega_0 = \sqrt{1/LC} = \sqrt{\frac{1}{(2.2 \mu H)(0.01 \mu F)}} = 6.74 \times 10^6 \text{ s}^{-1} = 2.15 \text{ MHz}$$

$$\gamma = R/L$$

$$\alpha = -\frac{\gamma}{2} \pm \sqrt{\frac{\gamma^2}{4} - (\omega_0)^2}$$

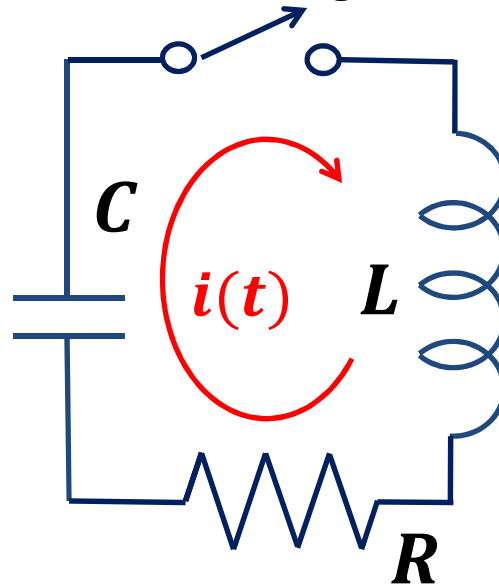
- Critical damping:

$$\frac{\gamma^2}{4} - (\omega_0)^2 = 0$$

$$\gamma = \frac{R}{L} = 2 \sqrt{\frac{1}{LC}}$$

$$R = 2 \sqrt{\frac{L}{C}} = 2 \sqrt{\frac{2.2 \mu H}{0.01 \mu F}} = 29.7 \Omega$$

Example



$$L = 2.2 \mu H$$

$$C = 10 nF$$

$$R = 2 \Omega$$

- Suppose the initial charge on the capacitor was 10 nC... What voltage is measured across R as a function of time?

Example

- Calculate the discriminant:

$$\frac{R^2}{4L^2} - \frac{1}{LC} = \frac{(2\ \Omega)^2}{4(2.2\ \mu H)^2} - \frac{1}{(2.2\ \mu H)(0.01\ \mu F)} < 0$$

- The circuit will oscillate with frequency:

$$f = \frac{1}{2\pi} \sqrt{\frac{1}{LC} - \frac{R^2}{4L^2}} = 1.07\ MHz$$

- Time constant:

$$\frac{\gamma}{2} = \frac{R}{2L} = \frac{(2\ \Omega)}{2(2.2\ \mu H)} = 4.55 \times 10^5\ s^{-1}$$

- Current will be:

$$i(t) = i_0 e^{-\gamma t/2} \cos(\omega t + \varphi)$$

Example

$$i(t) = i_0 e^{-\gamma t/2} \cos(\omega t + \varphi)$$

- Initial conditions:

- $i(0) = 0$ because the inductor produces a potential difference that opposes the change in current.

- Therefore, $\varphi = \pi/2$

- Initial potential across capacitor:

$$\Delta V = \frac{Q}{C} = \frac{10 \text{ nC}}{10 \text{ nF}} = 1 \text{ Volt}$$

- Initial voltage across inductor:

$$\Delta V = L \frac{di}{dt} = 1 \text{ Volt}$$

$$\frac{di}{dt} = \frac{1 \text{ V}}{2.2 \mu\text{H}} = 4.55 \times 10^5 \text{ A/s} = i_0 \omega$$

- $i_0 = \frac{4.55 \times 10^5 \text{ A/s}}{6.73 \times 10^6 /s} = 68 \text{ mA}$

Example

- Current in circuit:

$$i(t) = i_0 e^{-\gamma t/2} \sin \omega t$$

where $\gamma/2 = 4.55 \times 10^5 \text{ s}^{-1}$ and $i_0 = 68 \text{ mA}$

- Potential difference across the resistor:

$$\Delta V = i(t)R$$

$$v(t) = v_0 e^{-\gamma t/2} \sin \omega t$$

where $v_0 = 136 \text{ mV}$.

