

Physics 42200 Waves & Oscillations

Lecture 35 – Interference

Spring 2015 Semester

Matthew Jones

Optical Path Length

$$\vec{E}(z,t) = \vec{E}_0 \cos(kz - \omega t)$$

• What is the phase after propagating a distance d?

$$\delta = kd = \frac{2\pi d}{\lambda_0}$$

- $-\lambda_0$ is the wavelength in vacuum
- In a material with index of refraction n, $\lambda = \frac{\lambda_0}{n}$
- Phase is now

$$\delta' = k'd = \frac{2\pi nd}{\lambda_0}$$

Optical Path Length

- Consider light with wavelength λ_0 in vacuum that is split into two optical paths.
- Phase advance over each path:

$$\delta_1 = \frac{2\pi n_1 d_1}{\lambda_0} \qquad \qquad \delta_2 = \frac{2\pi n_2 d_2}{\lambda_0}$$

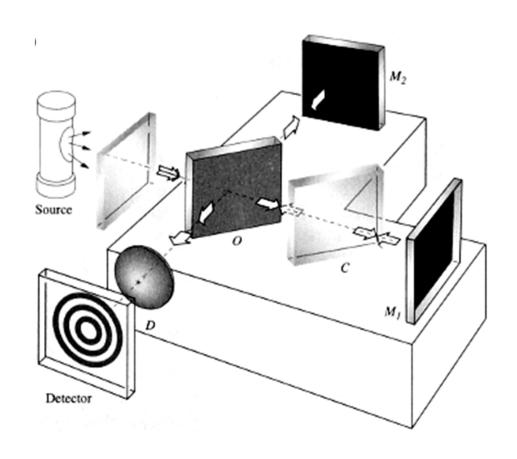
• Phase difference between two paths:

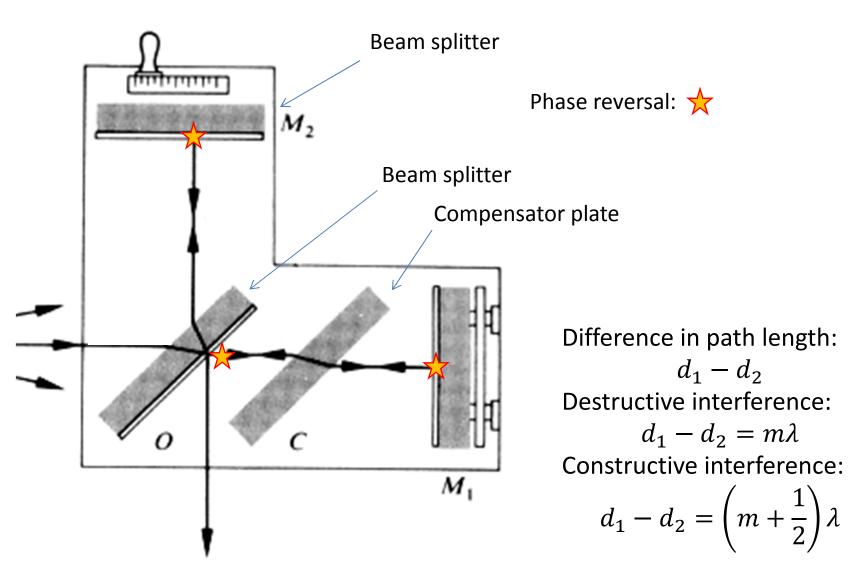
$$\Delta \delta = \frac{2\pi}{\lambda_0} (n_1 d_1 - n_2 d_2)$$

Constructive interference when

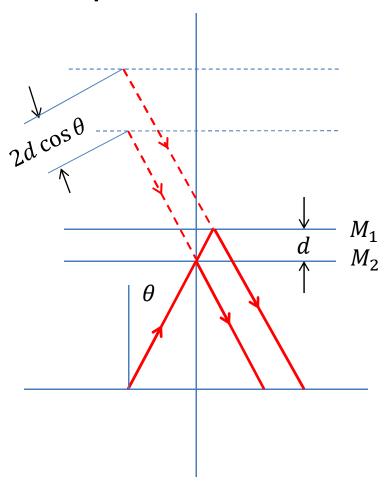
$$n_1 d_1 - n_2 d_2 = m \lambda_0$$
 where $m = 0,1,2,...$

• The *optical path length* is defined $\ell = nd$.



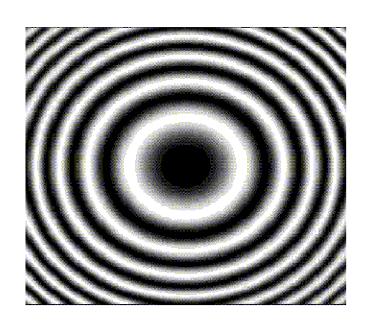


• Equivalent optics:



Bright fringes occur when

$$\delta = 2\pi \left(\frac{2d}{\lambda \cos \theta} + \frac{1}{2} \right) = 2\pi m$$



 How does the position of a fringe change when the path length changes?

$$\frac{2d}{\lambda \cos \theta} = m + \frac{1}{2}$$

$$2d = \lambda \cos \theta \left(m + \frac{1}{2} \right)$$

$$2\Delta d = -\lambda \sin \theta \left(m + \frac{1}{2} \right) \Delta \theta$$

$$\frac{\Delta \theta}{\Delta d} = -\frac{2}{\left(m + \frac{1}{2} \right) \lambda \sin \theta}$$

- Application: Consider two closely spaced wavelengths, λ and λ'
- Bright fringes from one wavelength occur when

$$\frac{2d}{\lambda} = m$$

Bright fringes from the other wavelength occur when

$$\frac{2d}{\lambda'} = m'$$

The two fringes will coincide when

$$\frac{2d}{\lambda} = \frac{2d}{\lambda'} + N$$

 Adjust the position of the movable mirror so that the next set of fringes coincide

$$\frac{2d'}{\lambda} = \frac{2d'}{\lambda'} + N + 1$$

Subtract these:

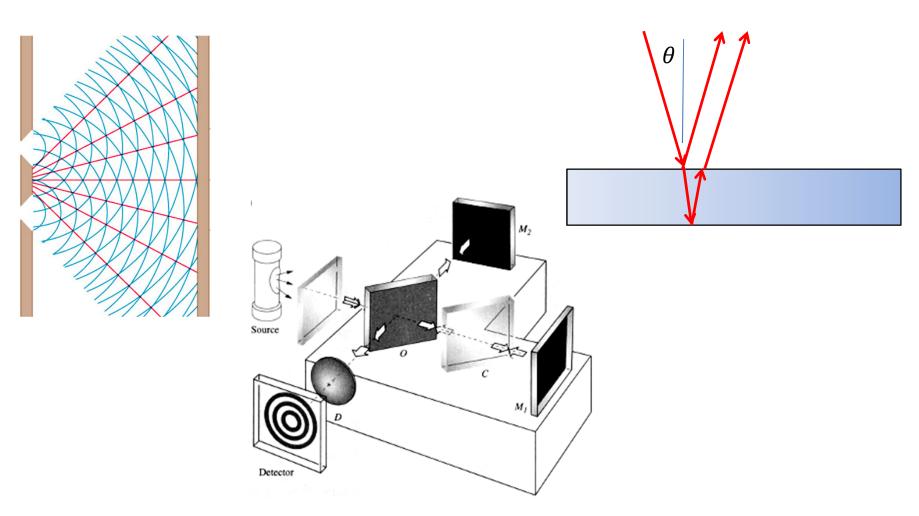
$$\frac{2d'}{\lambda} - \frac{2d}{\lambda} = \frac{2d'}{\lambda'} - \frac{2d}{\lambda'} + 1$$
$$\lambda' - \lambda = \frac{\lambda \lambda'}{2\Delta d} \approx \frac{\lambda^2}{2\Delta d}$$

For the yellow sodium line,

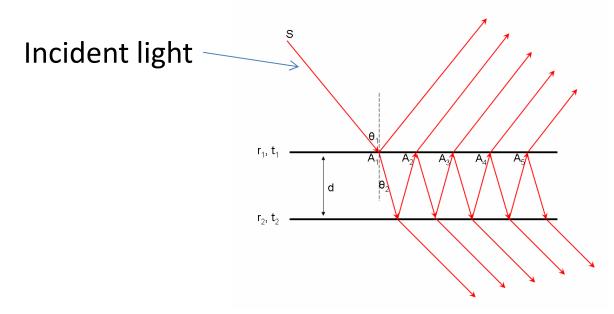
$$\lambda = 588.991 \ nm$$
 $\lambda' = 589.595 \ nm$
 $\Delta \lambda = 0.604 \ nm$
 $\Delta d = \lambda^2 / 2\Delta \lambda = 287,472 \ nm = 0.287 \ mm$



• Previously we considered only two interfering beams:



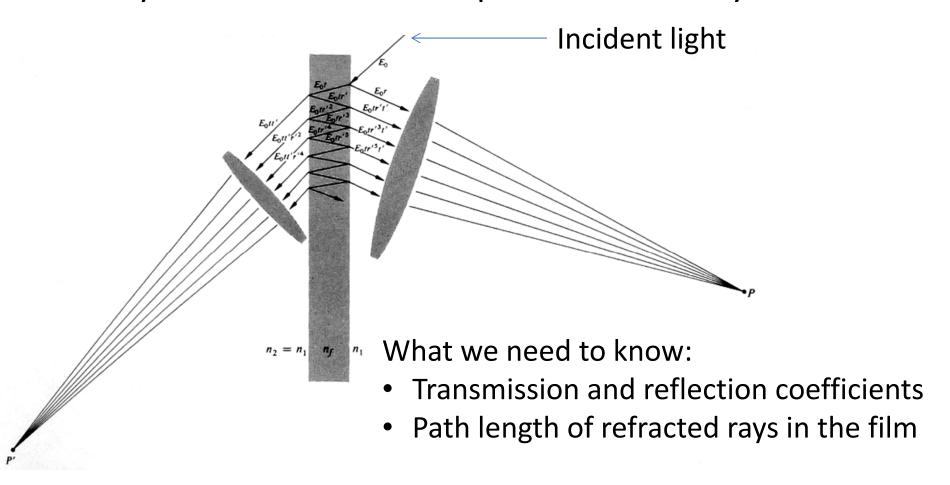
- In many situations, a coherent beam can interfere with itself multiple times
- Consider a beam incident on a thin film
 - Some component of the light will be reflected at each surface and some will be transmitted



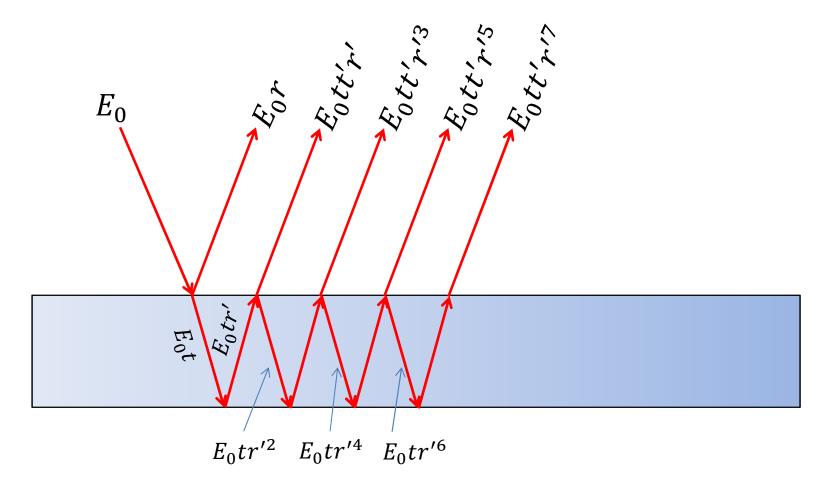
Each transmitted beam will have a different phase relative to the adjacent beams.

What is the total intensity of the reflected light?

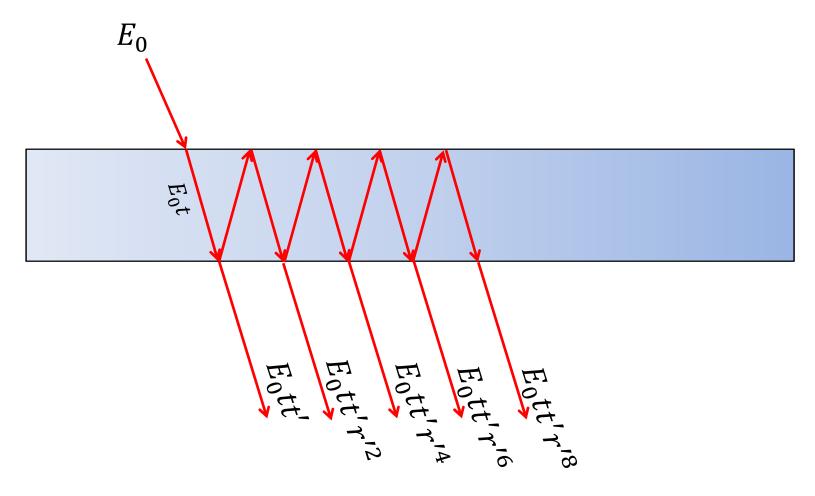
- All transmitted and reflected rays will be parallel
- They can be focused onto points P and P' by lenses:



- Reflection coefficients: r and r'
- Transmission coefficients: t and t'



- Reflection coefficients: r and r'
- Transmission coefficients: t and t'



The additional phase in the film is always the same:

$$\delta = \frac{2nkd}{\cos\theta_t}$$

If the initial phase is zero, then

$$E_{1r} = E_0 r e^{i\omega t}$$

$$E_{2r} = E_0 t t' r' e^{i(\omega t - \delta)}$$

$$E_{3r} = E_0 t t' r'^3 e^{i(\omega t - 2\delta)}$$

$$E_{4r} = E_0 t t' r'^5 e^{i(\omega t - 3\delta)}$$

• In general:

$$E_{Nr} = E_0 e^{i\omega t} t t' r'^{2N-3} e^{-i(N-1)\delta}$$
$$= E_0 e^{-i\omega t} t t' r' e^{-i\delta} (r'^2 e^{-i\delta})^{N-2}$$

The total electric field on one side of the film:

$$E_{r} = E_{0}e^{i\omega t}r + E_{0}e^{i\omega t}tt'r'e^{-i\delta} \times \left[1 + r'^{2}e^{-i\delta} + (r'^{2}e^{-i\delta})^{2} + (r'^{2}e^{-i\delta})^{3} + \cdots\right]$$

• This is in infinite sum of the form:

$$\sum_{n=0}^{\infty} z^n = \frac{1}{1-z} \quad \text{(when } |z| < 1)$$

Total electric field:

$$E_r = E_0 e^{i\omega t} \left[r + \frac{tt'r'e^{-i\delta}}{1 - r'^2 e^{-i\delta}} \right]$$

Simplifications:

$$r' = -r$$
$$tt' = 1 - r^2$$

Total electric field:

$$E_{r} = E_{0}e^{i\omega t}r \left[1 - \frac{(1 - r^{2})e^{-i\delta}}{1 - r^{2}e^{-i\delta}} \right]$$

$$= E_{0}e^{i\omega t}r \left[\frac{1 - r^{2}e^{-i\delta} - e^{-i\delta} + r^{2}e^{-i\delta}}{1 - r^{2}e^{-i\delta}} \right]$$

$$= E_{0}e^{i\omega t}r \left[\frac{1 - e^{-i\delta}}{1 - r^{2}e^{-i\delta}} \right]$$

• The intensity of the light is $I_r \propto |E_r|^2$

$$I_{r} = I_{0} \left\{ r \left[\frac{1 - e^{-i\delta}}{1 - r^{2}e^{-i\delta}} \right] \right\}^{*} \left\{ r \left[\frac{1 - e^{-i\delta}}{1 - r^{2}e^{-i\delta}} \right] \right\}$$

$$= I_{0}r^{2} \frac{(1 - e^{i\delta})(1 - e^{-i\delta})}{(1 - r^{2}e^{i\delta})(1 - r^{2}e^{-i\delta})}$$

$$= I_{0} \frac{2r^{2}(1 - \cos\delta)}{(1 + r^{4}) - 2r^{2}\cos\delta}$$

• The intensity of the transmitted light is $I_t \propto |E_t|^2$

$$I_t = I_0 \frac{1 - r^2}{(1 + r^4) - 2r^2 \cos \delta}$$

One more identity will clean this up a bit:

$$\cos \delta = 1 - 2\sin^2(\delta/2)$$

Reflected intensity:

$$I_r = I_0 \frac{F \sin^2(\delta/2)}{1 + F \sin^2(\delta/2)}$$

• Transmitted intensity:

$$I_t = I_0 \frac{1}{1 + F \sin^2(\delta/2)}$$

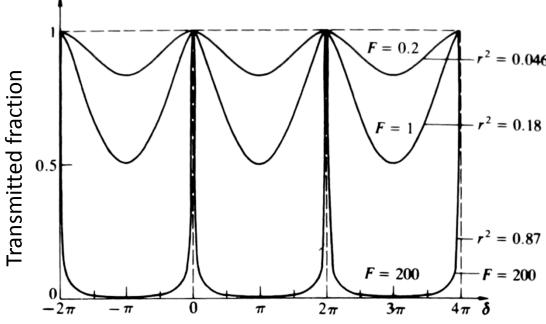
- The parameter $F = \left(\frac{2r}{1-r^2}\right)^2$ is called the *coefficient of finesse*
- Notice that $I_0 = I_r + I_t$
 - We assumed that no energy was lost in the film

The function

$$\mathcal{A}(\theta) = \frac{1}{1 + F \sin^2(\delta/2)}$$

is called the Airy function.

Remember that 8 is a Remember that 8 is a Remember the angle of the angle function of incidence, 0.



- In practice, some fraction of the light will be absorbed
- Absorptance, *A*, is defined by:

$$T + R + A = 1$$

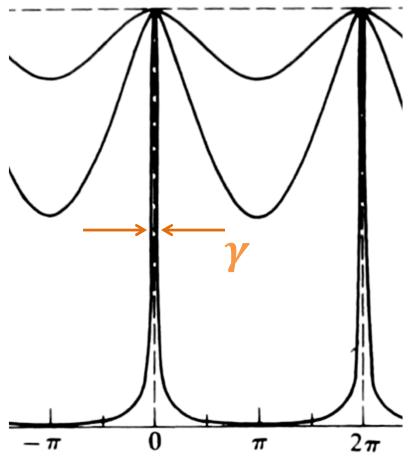
This modifies the transmitted intensity:

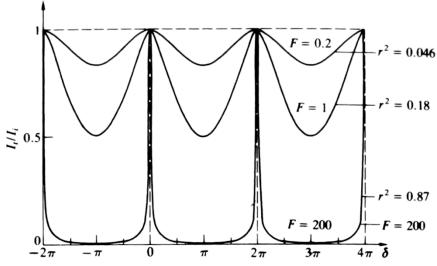
$$I_t = I_0 \left[1 - \frac{A}{1 - R} \right]^2 \mathcal{A}(\theta)$$

Example: silver film, 50 nm thick, deposited on glass

$$R = 0.94, T = 0.01, A = 0.05$$
$$\left[1 - \frac{A}{1 - R}\right]^{2} = 0.0278$$
$$F = 1044$$

How sharp are the peaks?





Width of one line:

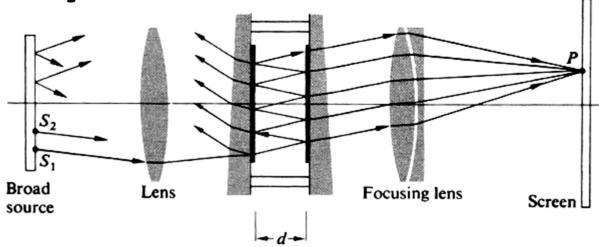
$$\gamma = 4/\sqrt{F}$$

Ratio of line spacing to the width:

$$\mathcal{F} = \frac{2\pi}{\gamma} = \frac{\pi\sqrt{F}}{2}$$

"Finesse" \mathcal{F} , not to be confused with the "coefficient of finesse" F.

Previous example: $\mathcal{F} \approx 50$



Phase difference:

$$\delta = \frac{4\pi n_f}{\lambda_0} d\cos\theta_i = 2\pi m$$
$$m\lambda_0 = 2n_f d\cos\theta_i$$

• Differentiate: $m \Delta \lambda_0 + \Delta m \lambda_0 = 0$ $\frac{\Delta m}{m} = -\frac{\Delta \lambda_0}{\lambda}$

$$\frac{\lambda_0}{\Delta \lambda_0} = \frac{2\pi m}{\Delta \delta}$$

Smallest resolvable wavelength difference:

$$(\Delta \lambda_0)_{min} = \frac{\lambda_0 (\Delta \delta)_{min}}{2\pi m}$$

Minimum resolvable phase shift:

$$(\Delta \delta)_{min} \sim \gamma = 4/\sqrt{F}$$

Chromatic resolving power:

$$\mathcal{R} = \frac{\lambda_0}{(\Delta \lambda_0)_{min}} \approx \mathcal{F} m \approx \mathcal{F} \frac{2n_f d}{\lambda_0}$$

Typical values:

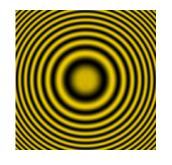
$$\mathcal{F} = 50$$

$$n_f d = 1 cm$$

$$\lambda_0 = 500 nm$$

$$\mathcal{R} = \frac{2 \times 50 \times 1 cm}{500 nm} = 2 \times 10^6$$

Diffraction grating



Michelson interferometer



Fabry-Perot interferometer

 The effective gap between the surfaces can be adjusted by changing the pressure of a gas, or by means of piezoelectric actuators

