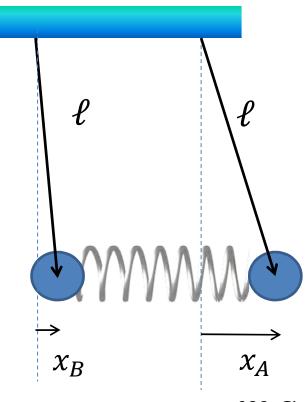


Physics 42200 Waves & Oscillations

Lecture 11 – French, Chapter 5

Spring 2015 Semester

Matthew Jones



- The spring is stretched by the amount $x_A x_B$
- Restoring force on pendulum A:

$$F_A = -k(x_A - x_B)$$

Restoring force on pendulum B:

$$F_B = k(x_A - x_B)$$

$$m\ddot{x}_A + \frac{mg}{\ell}x_A + k(x_A - x_B) = 0$$

$$m\ddot{x}_B + \frac{mg}{\ell}x_B - k(x_A - x_B) = 0$$

$$\ddot{x}_A + [(\omega_0)^2 + (\omega_c)^2] x_A - (\omega_c)^2 x_B = 0$$

$$\ddot{x}_B + [(\omega_0)^2 + (\omega_c)^2] x_B - (\omega_c)^2 x_A = 0$$

$$\omega_0 = \sqrt{g/\ell}, \ \omega_c = \sqrt{k/m}$$

Add equations for A and B together:

$$\frac{d^2}{dt^2}(x_A + x_B) + (\omega_0)^2(x_A + x_B) = 0$$

Subtract equations A and B:

$$\frac{d^2}{dt^2}(x_A - x_B) + [(\omega_0)^2 + 2(\omega_c)^2](x_A - x_B) = 0$$

$$\frac{d^2}{dt^2}(x_A + x_B) + (\omega_0)^2(x_A + x_B) = 0$$

$$\frac{d^2}{dt^2}(x_A - x_B) + (\omega')^2(x_A - x_B) = 0$$

$$\omega_0 = \sqrt{g/\ell}, \, \omega' = \sqrt{(\omega_0)^2 + 2(\omega_c)^2}$$

Normal coordinates:

$$q_1 = x_A + x_B$$
$$q_2 = x_A - x_B$$

Decoupled equations:

$$\ddot{q}_1 + (\omega_0)^2 q_1 = 0$$

$$\ddot{q}_2 + (\omega')^2 q_2 = 0$$

Decoupled equations:

$$\ddot{q}_1 + (\omega_0)^2 q_1 = 0$$

$$\ddot{q}_2 + (\omega')^2 q_2 = 0$$

Solutions are

$$q_1(t) = A\cos(\omega_0 t + \alpha)$$

$$q_2(t) = B\cos(\omega' t + \beta)$$

• The variables q_1 and q_2 are called "normal coordinates".

Initial Conditions

Suppose we had the initial conditions:

$$x_A = A_0$$
 $\dot{x}_A = 0$
 $x_B = 0$ $\dot{x}_B = 0$

• These can be satisfied with $\alpha = \beta = 0$:

$$x_A(t) = \frac{1}{2}(q_1 + q_2) = \frac{1}{2}A\cos\omega_0 t + \frac{1}{2}B\cos\omega' t$$

$$x_B(t) = \frac{1}{2}(q_1 - q_2) = \frac{1}{2}A\cos\omega_0 t - \frac{1}{2}B\cos\omega' t$$

• At time t = 0,

$$\frac{1}{2}(A+B) = A_0 \qquad \frac{1}{2}(A-B) = 0$$

• Now we know that $A = B = A_0$.

Initial Conditions

Velocity:

$$\dot{x}_A(t) = -\frac{1}{2}A_0\omega_0\sin\omega_0 t - \frac{1}{2}A_0\omega'\sin\omega' t$$

$$\dot{x}_B(t) = -\frac{1}{2}A_0\omega_0\sin\omega_0 t + \frac{1}{2}A_0\omega'\sin\omega' t$$

• Initial conditions are satisfied at t=0.

Initial Conditions

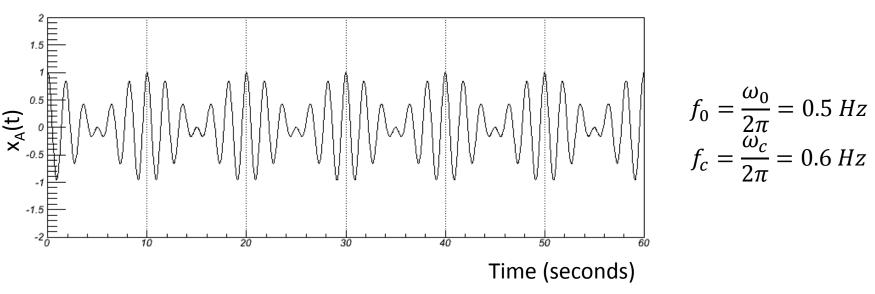
Complete solution:

$$x_A(t) = \frac{1}{2} A_0(\cos \omega_0 t + \cos \omega' t)$$

$$= A_0 \cos \left(\frac{\omega' - \omega_0}{2}t\right) \cos \left(\frac{\omega' + \omega_0}{2}t\right)$$

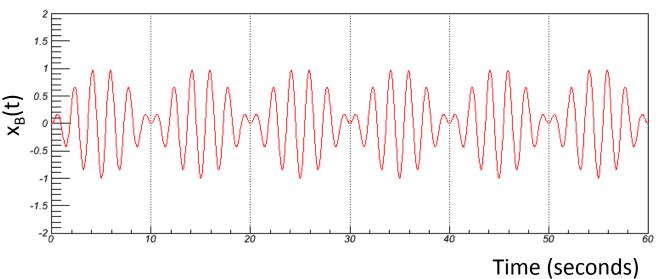
$$x_B(t) = \frac{1}{2} A_0(\cos \omega_0 t - \cos \omega' t)$$

$$= A_0 \sin \left(\frac{\omega' - \omega_0}{2}t\right) \sin \left(\frac{\omega' + \omega_0}{2}t\right)$$



$$f_0 = \frac{\omega_0}{2\pi} = 0.5 \text{ Hz}$$

$$f_c = \frac{\omega_c}{2\pi} = 0.6 \text{ Hz}$$



- This procedure worked, but the problem was very simple. How can we apply this in general?
- Procedure:
 - 1. Construct the set of coupled differential equations
 - 2. Assume solutions are of the form

$$q_i(t) = A_i \cos(\omega t + \varphi_i)$$

- 3. Substitute into the differential equations
- 4. Find the values of ω that satisfy the resulting matrix equation (eigenvalues).
- 5. Solve for constants of integration

$$\ddot{x}_A + [(\omega_0)^2 + (\omega_c)^2] x_A - (\omega_c)^2 x_B = 0$$

$$\ddot{x}_B + [(\omega_0)^2 + (\omega_c)^2] x_B - (\omega_c)^2 x_A = 0$$

$$\omega_0 = \sqrt{g/\ell}, \ \omega_c = \sqrt{k/m}$$

- \triangleright Let $x_A(t) = A\cos(\omega t + \alpha)$
- \triangleright Let $x_B(t) = B\cos(\omega t + \beta)$
- Second derivatives:

$$\Rightarrow \ddot{x}_A(t) = -A\omega^2 \cos(\omega t + \alpha) = -\omega^2 x_A$$

$$\Rightarrow \ddot{x}_B(t) = -B\omega^2 \cos(\omega t + \beta) = -\omega^2 x_B$$

$$\begin{pmatrix} (\omega_0)^2 + (\omega_c)^2 - \omega^2 & -(\omega_c)^2 \\ -(\omega_c)^2 & (\omega_0)^2 + (\omega_c)^2 - \omega^2 \end{pmatrix} \begin{pmatrix} x_A \\ x_B \end{pmatrix} = 0$$

A very important result:

If the matrix equation

$$A\vec{v}=0$$

for any vector \vec{v} , then

$$\det A = 0$$

 You are expected to be able to calculate the determinant of an arbitrary 2x2 or 3x3 matrix!

$$\begin{pmatrix} (\omega_0)^2 + (\omega_c)^2 - \omega^2 & -(\omega_c)^2 \\ -(\omega_c)^2 & (\omega_0)^2 + (\omega_c)^2 - \omega^2 \end{pmatrix} \begin{pmatrix} x_A \\ x_B \end{pmatrix} = 0$$

The determinant of the matrix is:

$$[(\omega_0)^2 + (\omega_c)^2 - \omega^2]^2 - (\omega_c)^4 = 0$$

• Expand the polynomial in $\lambda = \omega^2$:

$$\lambda^2 - 2((\omega_0)^2 + (\omega_c)^2) + [(\omega_0)^2 + (\omega_c)^2]^2 - (\omega_c)^4 = 0$$

• Use the quadratic formula:

$$\lambda = ((\omega_0)^2 + (\omega_c)^2) \pm \sqrt{(\omega_c)^4}$$

Oscillation frequencies are

$$\omega^2 = (\omega_0)^2$$
$$\omega'^2 = (\omega_0)^2 + 2(\omega_c)^2$$

Eigenvalue problem:

$$\begin{pmatrix} (\omega_0)^2 + (\omega_c)^2 & -(\omega_c)^2 \\ -(\omega_c)^2 & (\omega_0)^2 + (\omega_c)^2 \end{pmatrix} \begin{pmatrix} \chi_A \\ \chi_B \end{pmatrix} = \omega^2 \begin{pmatrix} \chi_A \\ \chi_B \end{pmatrix}$$

• First eigenvector: substitute $\omega^2 = (\omega_0)^2$

$$\begin{pmatrix} (\omega_c)^2 & -(\omega_c)^2 \\ -(\omega_c)^2 & (\omega_c)^2 \end{pmatrix} \begin{pmatrix} x_A \\ x_B \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \Rightarrow x_B = x_A$$

• Second eigenvector: substitute $\omega'^2 = (\omega_0)^2 + 2(\omega_c)^2$

$$\begin{pmatrix} -(\omega_c)^2 & -(\omega_c)^2 \\ -(\omega_c)^2 & -(\omega_c)^2 \end{pmatrix} \begin{pmatrix} x_A \\ x_B \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \Rightarrow x_B = -x_A$$

Normal Coordinates

 The first normal mode of vibration corresponds to the first eigenvector:

$$\vec{q}_1(t) = A \begin{pmatrix} 1 \\ 1 \end{pmatrix} \cos(\omega_0 t + \alpha)$$

 The second normal mode of vibration corresponds to the second eigenvector:

$$\vec{q}_2(t) = B \begin{pmatrix} 1 \\ -1 \end{pmatrix} \cos(\omega' t + \beta)$$

Arbitrary motion:

$${x_A(t) \choose x_B(t)} = \mathbf{A} {1 \choose 1} \cos(\omega_0 t + \mathbf{\alpha}) + \mathbf{B} {1 \choose -1} \cos(\omega' t + \mathbf{\beta})$$

• Initial conditions determine constants of integration.

Forced Coupled Oscillator

 What happens when a driving force is applied to one of the oscillators?

$$\ddot{x}_A + [(\omega_0)^2 + (\omega_c)^2] x_A - (\omega_c)^2 x_B = F_0/m \cos \omega t$$
$$\ddot{x}_B + [(\omega_0)^2 + (\omega_c)^2] x_B - (\omega_c)^2 x_A = 0$$

Normal coordinates:

$$q_1 = x_A + x_B$$
$$q_2 = x_A - x_B$$

Equations of motion:

$$\ddot{q}_1 + (\omega_0)^2 q_1 = F_0/m \cos \omega t$$

$$\ddot{q}_2 + (\omega')^2 q_2 = F_0/m \cos \omega t$$

Decoupled equations which we know how to solve.

Forced Coupled Oscillators

Steady state amplitudes:

$$A_1(\omega) = \frac{F_0/m}{(\omega_0)^2 - \omega^2}$$
$$A_2(\omega) = \frac{F_0/m}{(\omega')^2 - \omega^2}$$

Motion of individual masses:

$$x_A(t) = \frac{1}{2}(q_1(t) + q_2(t))$$

Amplitude of steady state oscillations:

$$A(\omega) = \frac{F_0}{2m} \left(\frac{1}{(\omega_0)^2 - \omega^2} + \frac{1}{(\omega')^2 - \omega^2} \right)$$

$$= \frac{F_0}{2m} \frac{(\omega')^2 + (\omega_0)^2 - 2\omega^2}{\left((\omega_0)^2 - \omega^2\right)\left((\omega')^2 - \omega^2\right)} = \frac{F_0}{m} \frac{\left((\omega_0)^2 + (\omega_c)^2\right) - \omega^2}{\left((\omega_0)^2 - \omega^2\right)\left((\omega')^2 - \omega^2\right)}$$

Forced Coupled Oscillators

Motion of individual masses:

$$x_B(t) = \frac{1}{2}(q_1(t) - q_2(t))$$

Amplitude of steady state oscillations:

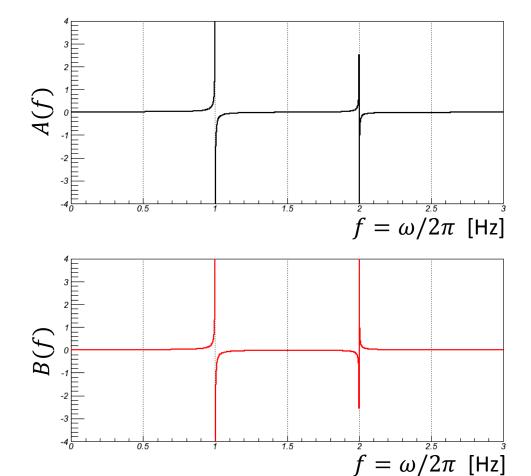
$$B(\omega) = \frac{F_0}{2m} \left(\frac{1}{(\omega_0)^2 - \omega^2} - \frac{1}{(\omega')^2 - \omega^2} \right)$$

$$= \frac{F_0}{2m} \frac{(\omega')^2 - (\omega_0)^2}{((\omega_0)^2 - \omega^2)((\omega')^2 - \omega^2)}$$

$$= \frac{F_0}{m} \frac{(\omega_c)^2}{((\omega_0)^2 - \omega^2)((\omega')^2 - \omega^2)}$$

• Now the system resonates at two frequencies: ω_o and ω' .

Forced Coupled Oscillators



Phase reverses when crossing each of the resonant frequencies.