

Physics 42200 Waves & Oscillations

Lecture 9 – French, Chapter 4

Spring 2013 Semester

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$$m\ddot{x} + b\dot{x} + kx = F_0 \cos \omega t$$

> Frequency of free oscillations:

$$\omega_0 = \sqrt{\frac{k}{m}} \qquad \qquad \omega_{free} = \sqrt{\frac{k}{m} - \frac{b^2}{2m^2}}$$

- \triangleright Driving frequency: ω
- \triangleright Steady state solution ($t \gg 1/\gamma$):

$$x(t) = A \cos(\omega t - \delta)$$

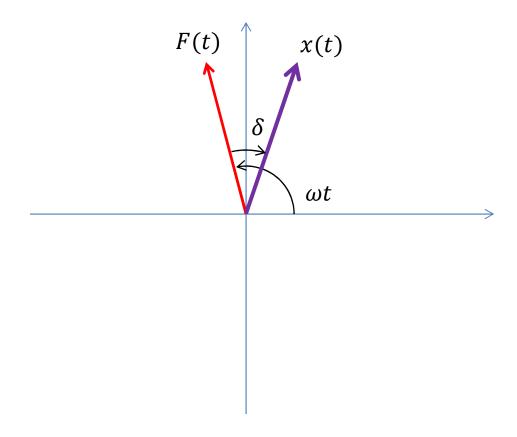
> Amplitude of steady-state oscillations:

$$A = \frac{F_0/m}{\sqrt{\left((\omega_0)^2 - \omega^2\right)^2 + (\omega\gamma)^2}}$$

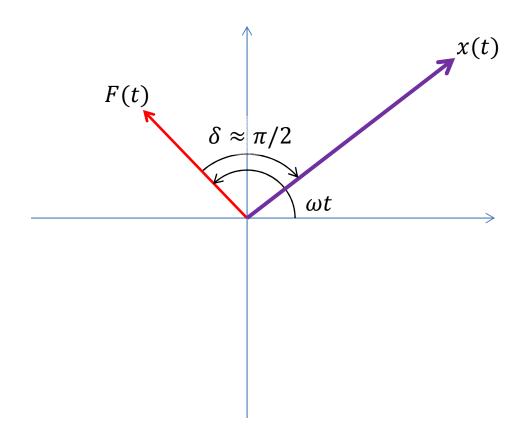
Phase difference:

$$\delta = \tan^{-1} \left(\frac{\omega \gamma}{(\omega_0)^2 - \omega^2} \right)$$

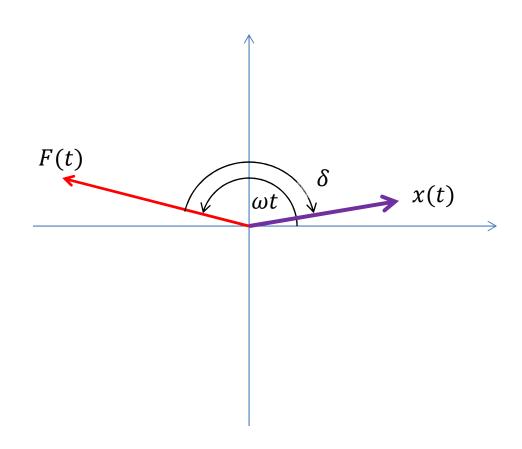
• Phasor diagram: $\omega < \omega_0$



• Phasor diagram: $\omega \approx \omega_0$

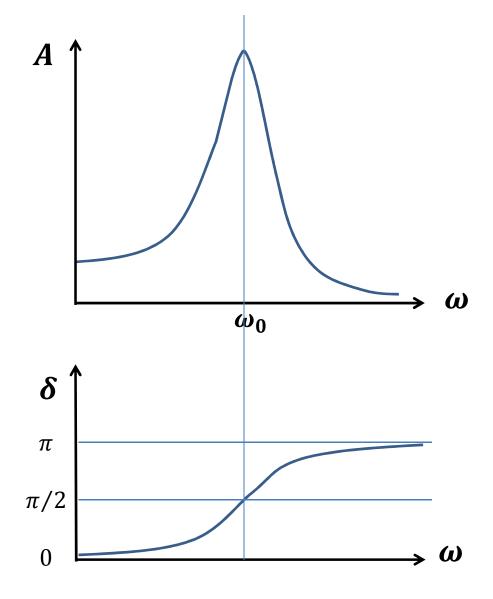


• Phasor diagram: $\omega > \omega_0$



Resonance

- The peak occurs at a frequency that is close to, but not exactly equal to ω_0 .
- At resonance, the phase shift is $\delta = \pi/2$.
- The force pushes the mass in the direction it is already moving adding energy to the system.



"Quality Factor"

- Instead of using $\gamma = b/m$ and $\omega_0 = \sqrt{k/m}$, it is convenient to describe the shape of the resonance curve using the variables ω_0 and $Q = \omega_0/\gamma$.
- $Q = \omega_0/\gamma$ is called the "quality factor".
- Written in terms of ω_0 and Q, the amplitude is

$$A = \frac{F_0/m}{\sqrt{\left((\omega_0)^2 - \omega^2\right)^2 + (\omega\omega_0)^2/Q^2}}$$

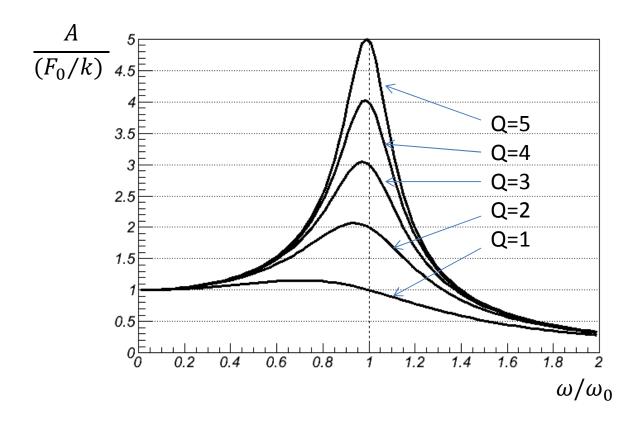
$$= \frac{F_0}{k} \frac{\omega_0/\omega}{\left[\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}\right]^{1/2}}$$

Quality Factor

$$A(\omega) = \frac{F_0}{k} \frac{\omega_0/\omega}{\left[\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}\right]^{1/2}}$$

- Why is this a convenient form?
 - Dimensionless quantities are easier to analyze
 - The scale of the amplitude is determined by F_0/k
 - The shape of the curve is determined by the dimensionless quantities ω/ω_0 and Q

Quality Factor



The normalized height is approximately QThe maximum occurs when $\omega/\omega_0\approx 1$ At resonance, the motion is amplified by the factor Q.

Energy

- An oscillator stores energy
- The driving force adds energy to the system
- The damping force dissipates energy
- Instantaneous rate at which energy is added:

$$P = \frac{dW}{dt} = F\frac{dx}{dt} = F\dot{x}$$

$$F(t) = F_0 \cos \omega t$$

$$x(t) = A \cos(\omega t - \delta)$$

$$\dot{x}(t) = -A\omega \sin(\omega t - \delta)$$

$$P = -F_0 A\omega \cos(\omega t) \sin(\omega t - \delta)$$

Average rate at which energy is added:

$$\bar{P}(\omega) = \frac{1}{2} F_0 A \omega \sin \delta$$

• Maximal when $\delta = \pi/2$

Energy

$$\bar{P}(\omega) = \frac{1}{2} F_0 A \omega \sin \delta$$

Some algebra:

$$A(\omega) = \frac{F_0}{k} \frac{\omega_0/\omega}{\left[\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}\right]^{1/2}}$$

$$\sin \delta = \frac{1/Q}{\left[\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}\right]^{1/2}}$$

Average power:

$$\overline{P}(\omega) = \frac{(F_0)^2 \omega_0}{2kQ} \frac{1}{\left(\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \frac{1}{Q^2}}$$

Energy

• When $\omega \approx \omega_0$ we can simplify further:

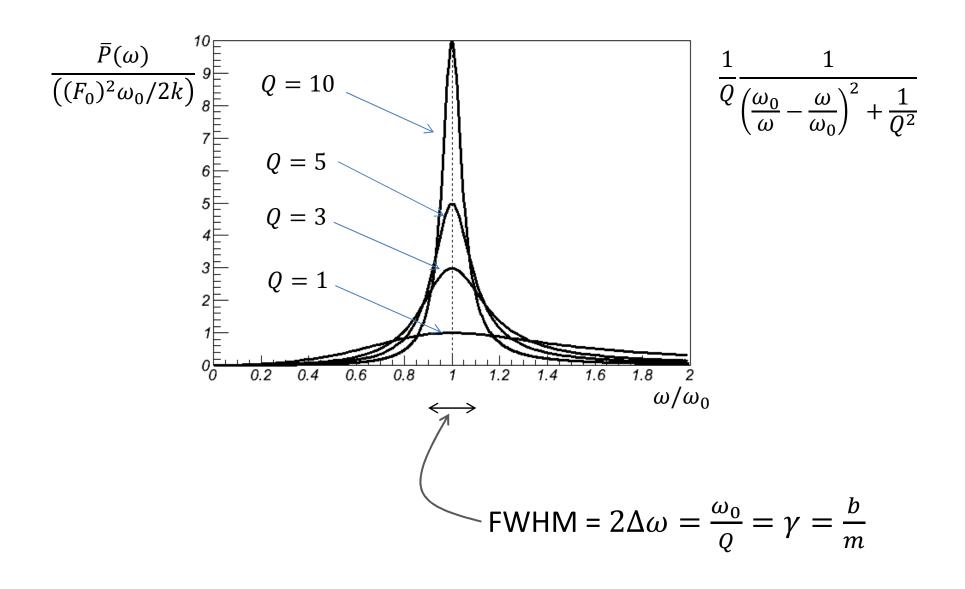
$$\bar{P}(\omega) = \frac{\omega_0 + \Delta\omega}{\frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}} \approx \frac{2}{\omega_0} \Delta\omega$$

$$\bar{P}(\omega) = \frac{(F_0)^2 (\omega_0/Q)}{2k/(\omega_0)^2} \frac{1}{4(\Delta\omega)^2 + (\omega_0/Q)^2}$$

• What value of $\Delta \omega$ will reduce the peak power by a factor of ½?

$$\frac{1}{4(\Delta\omega)^2 + (\omega_0/Q)^2} = \frac{1}{2} \frac{1}{(\omega_0/Q)^2} \implies 2\Delta\omega = \omega_0/Q$$

Power Resonance Shape



Resonance Curves

- General properties:
 - Amplitude at resonance: Static displacement x Q
 - FWHM power bandwidth: $\gamma = \omega_0/Q$
 - When Q is large, a small force at the resonant frequency produces large oscillations
 - Large amplitudes persist only when the driving force is near the oscillation frequency

 So far, we only considered the form of solutions when t was very large:

$$x_1(t) = A\cos(\omega t - \delta)$$

- What is the form of the solution when t is small?
- The solution with no forcing term was $x_2(t) = \mathbf{B}e^{-\gamma t/2}\cos\omega_0 t + \mathbf{C}e^{-\gamma t/2}\sin\omega_0 t$
- Complete solution:

$$x(t) = x_1(t) + x_2(t)$$

Initial conditions determine B and C in the complete solution.

• Suppose a mass is already in motion at t = 0:

$$x(0) = A_0$$

$$\dot{x}(0) = 0$$

• Suppose that γ is small, so that this motion persists for a long time: ignore the $e^{-\gamma t/2}$ terms

$$x_2(t) = B \cos \omega_0 t$$

Steady state solution:

$$x_1(t) = A\cos(\omega t - \delta), \qquad A = \frac{F_0/m}{\sqrt{((\omega_0)^2 - \omega^2)^2 + (\omega \omega_0)^2/Q^2}}$$

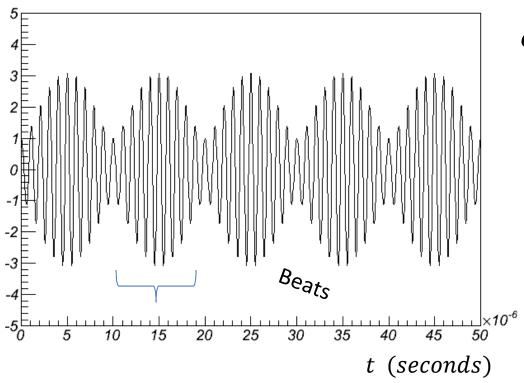
Complete solution:

$$x(t) = B\cos\omega_0 t + A\cos(\omega t - \delta)$$

- At t=0, $\dot{x}(t)=0$: $\dot{x}(t)=-B\omega_0\sin\omega_0t-A\omega\sin(\omega t-\delta)$
- The phase of the driving force must be 0 or π .
- Amplitude at t = 0:

$$A_0 = B + A$$

• A and A_0 are given, so $B = A_0 - A$ $x(t) = (A_0 - A) \cos \omega_0 t + A \cos \omega t$ $= A_0 \cos \omega_0 t + A(\cos \omega t - \cos \omega_0 t)$ $= A_0 \cos \omega_0 t + 2A \sin(\overline{\omega}t) \sin\left(\frac{1}{2}\Delta\omega t\right)$

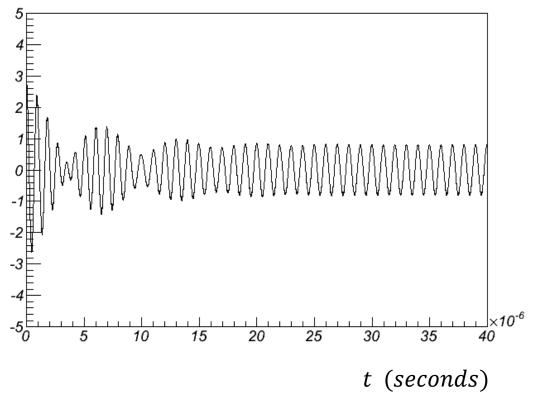


$$\omega_0 = 2\pi \times (1.0 \text{ Mhz})$$

$$\omega = 2\pi \times (1.1 \text{ Mhz})$$

$$Q = 2$$

Accounting for energy dissipation looks like this:



$$\omega_0 = 2\pi \times (1.0 \text{ Mhz})$$

$$\omega = 2\pi \times (1.15 \text{ Mhz})$$

$$Q = 20$$

$$\frac{1}{\gamma} = \frac{Q}{\omega_0} = 3.2 \,\mu\text{s}$$