

Physics 42200

Waves & Oscillations

Lecture 17 – French, Chapter 6

Spring 2013 Semester

Matthew Jones

Midterm Exam:

Date: Wednesday, March 6th

Time: 8:00 – 10:00 pm

Room: PHYS 203

Material: French, chapters 1-8

Fourier Analysis

- Wave equation:

$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2}$$

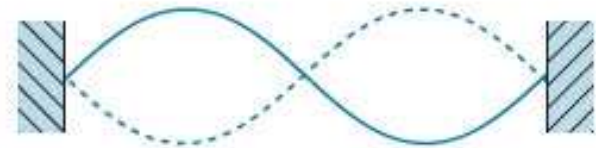
- Boundary conditions at $t = 0$:
 $y(0, t) = y(L, t) = 0$

- Normal modes of oscillation:

$$\omega_n = \frac{n\pi v}{L}$$

- Wavelengths of normal modes:

$$\lambda_n = \frac{2L}{n}$$



Fourier Analysis

- Normal modes of oscillation:

$$y_n(x, t) = \sin\left(\frac{n\pi x}{L}\right) \cos(\omega_n t)$$

- General solution:

$$y(x, t) = \sum_{n=1}^{\infty} A_n \sin\left(\frac{n\pi x}{L}\right) \cos(\omega_n t - \delta_n)$$

- The parameters A_n and δ_n are determined from initial conditions. At $t = 0$,

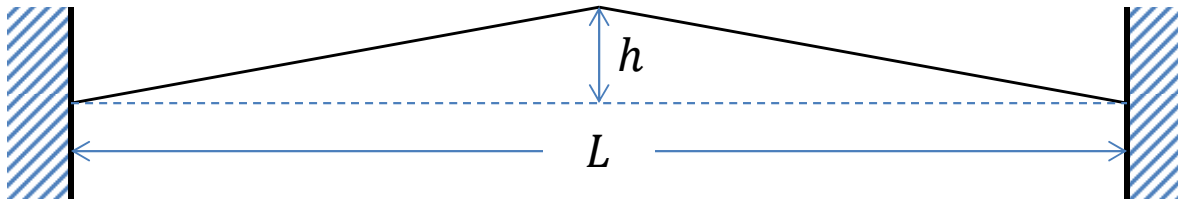
$$y(x, 0) = u(x) = \sum_{n=1}^{\infty} B_n \sin\left(\frac{n\pi x}{L}\right)$$

Fourier Analysis

- Fourier's method:

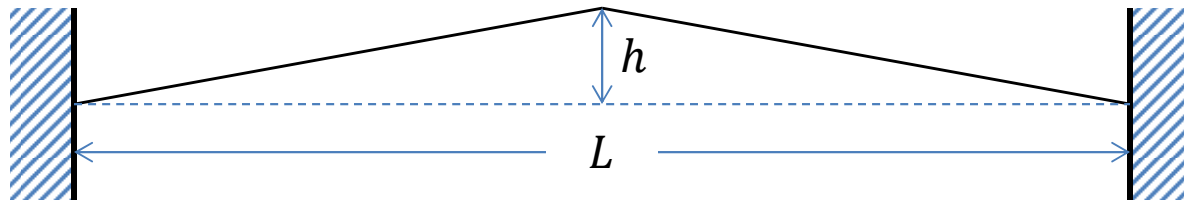
$$B_k = \frac{2}{L} \int_0^L \sin\left(\frac{k\pi x}{L}\right) u(x) dx$$

- Example: (*French 6-12*):



- The string has tension T and mass per unit length μ
- If the string is released from rest, what is $y(x, t)$?

Example



- Initial displacement is described by the function

$$u(x) = \begin{cases} \frac{2hx}{L} & 0 < x < L/2 \\ 2h \left(1 - \frac{x}{L}\right) & L/2 < x < L \end{cases}$$

- The string is released from rest, so $v(x) = 0$.

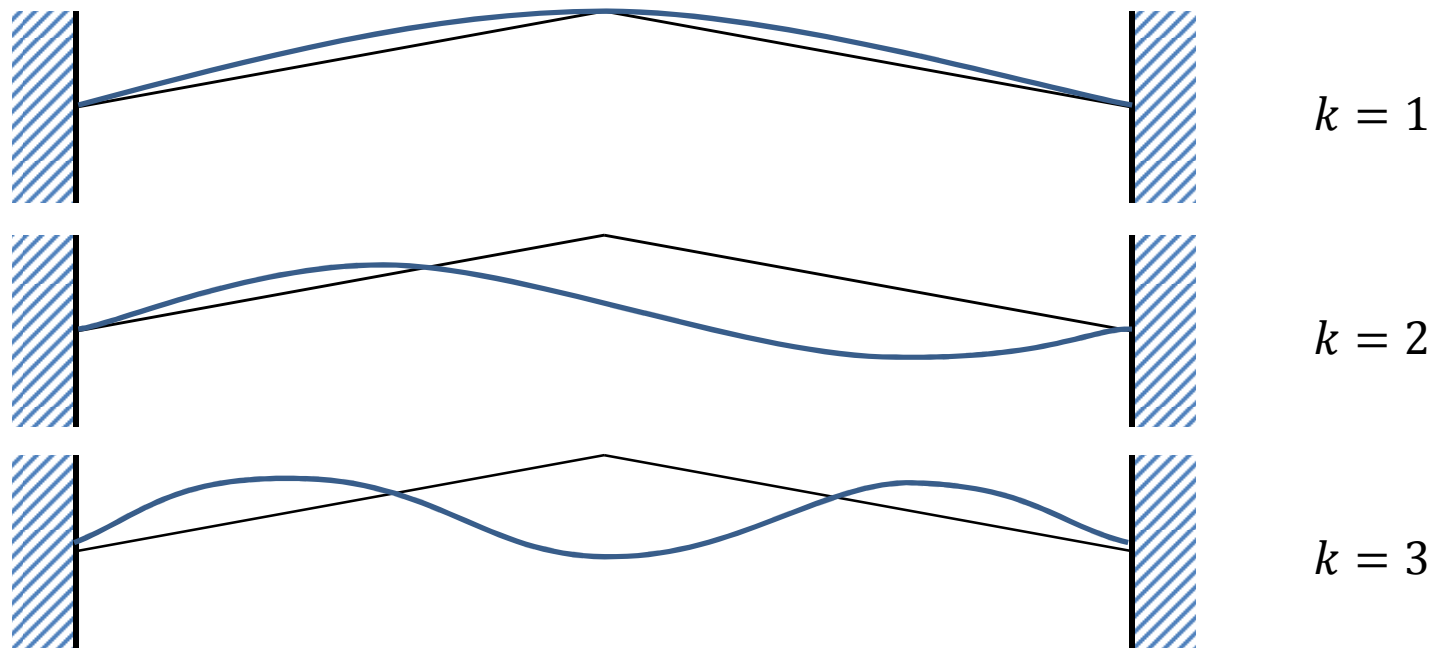
Example

- Fourier coefficients:

$$\begin{aligned} B_k &= \frac{2}{L} \int_0^L \sin\left(\frac{k\pi x}{L}\right) u(x) dx \\ &= \frac{2}{L} \int_0^{L/2} \sin\left(\frac{k\pi x}{L}\right) \left(\frac{2hx}{L}\right) dx \\ &\quad + \frac{2}{L} \int_{L/2}^L \sin\left(\frac{k\pi x}{L}\right) (2h(1 - x/L)) dx \\ &= \frac{4h}{L^2} \int_0^{L/2} x \sin\left(\frac{k\pi x}{L}\right) dx - \frac{4h}{L^2} \int_{L/2}^L x \sin\left(\frac{k\pi x}{L}\right) dx + \frac{4h}{L} \int_{L/2}^L \sin\left(\frac{k\pi x}{L}\right) dx \end{aligned}$$

- This is getting messy...*
- Don't get your table of integrals just yet...*
- Think about the symmetry of the problem.*

Example



- Integrals with even values of k will be zero.
 - This is because $\sin(k\pi x/L)$ is odd but $u(x)$ is even when reflected about the point $x = L/2$.
- When k is odd, we can just double the value of the integral from 0 to $L/2$.

Example

- When k is odd,

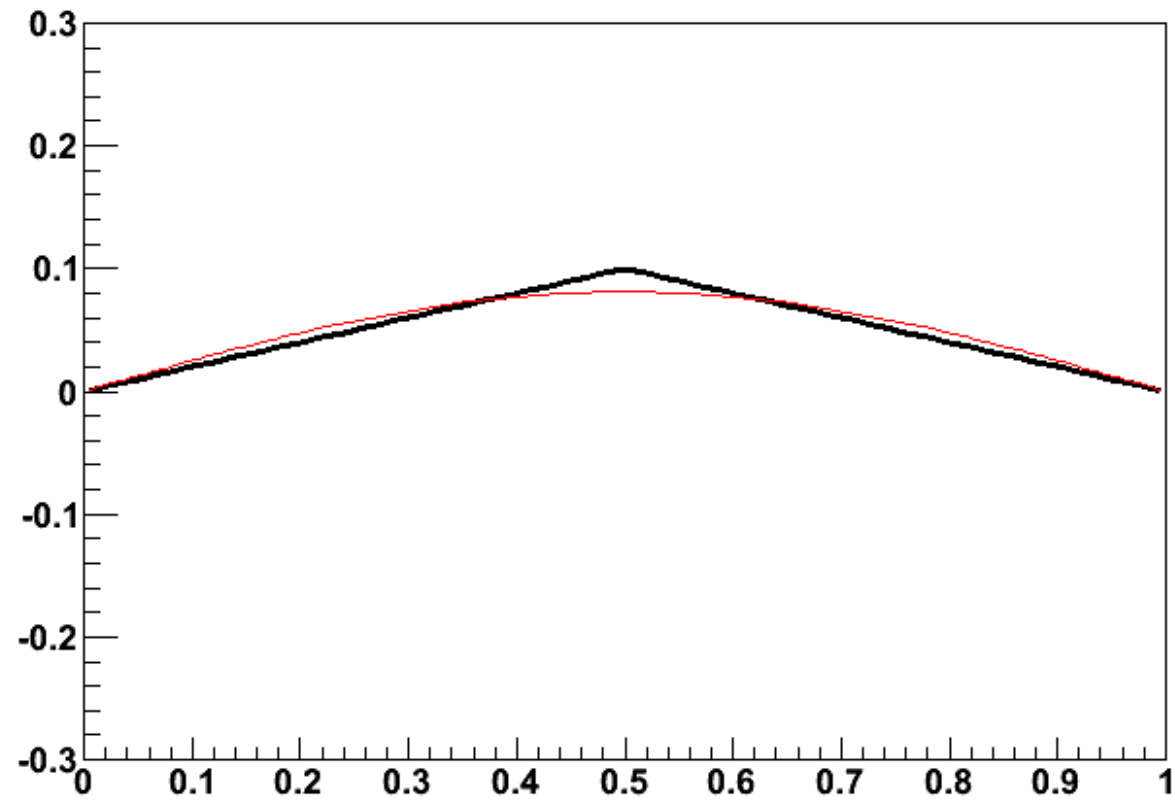
$$\begin{aligned} B_k &= \frac{2}{L} \int_0^L \sin\left(\frac{k\pi x}{L}\right) u(x) dx = \frac{4}{L} \int_0^{L/2} \sin\left(\frac{k\pi x}{L}\right) \left(\frac{2hx}{L}\right) dx \\ &= \frac{8h}{L^2} \int_0^{L/2} x \sin\left(\frac{k\pi x}{L}\right) dx \end{aligned}$$

- Now get your table of integrals:

$$(91) \quad \int x \sin(ax) dx = -\frac{x}{a} \cos ax + \frac{1}{a^2} \sin ax$$

$$\begin{aligned} B_k &= -\frac{8h}{L^2} \frac{Lx}{k\pi} \cos\left(\frac{k\pi x}{L}\right) + \frac{8h}{L^2} \frac{L^2}{k^2 \pi^2} \sin\left(\frac{k\pi x}{L}\right) \Big|_0^{L/2} \\ &= \pm \frac{8h}{k^2 \pi^2} \end{aligned}$$

Example



$$B_1 = \frac{8h}{\pi^2} = 0.8106 h$$

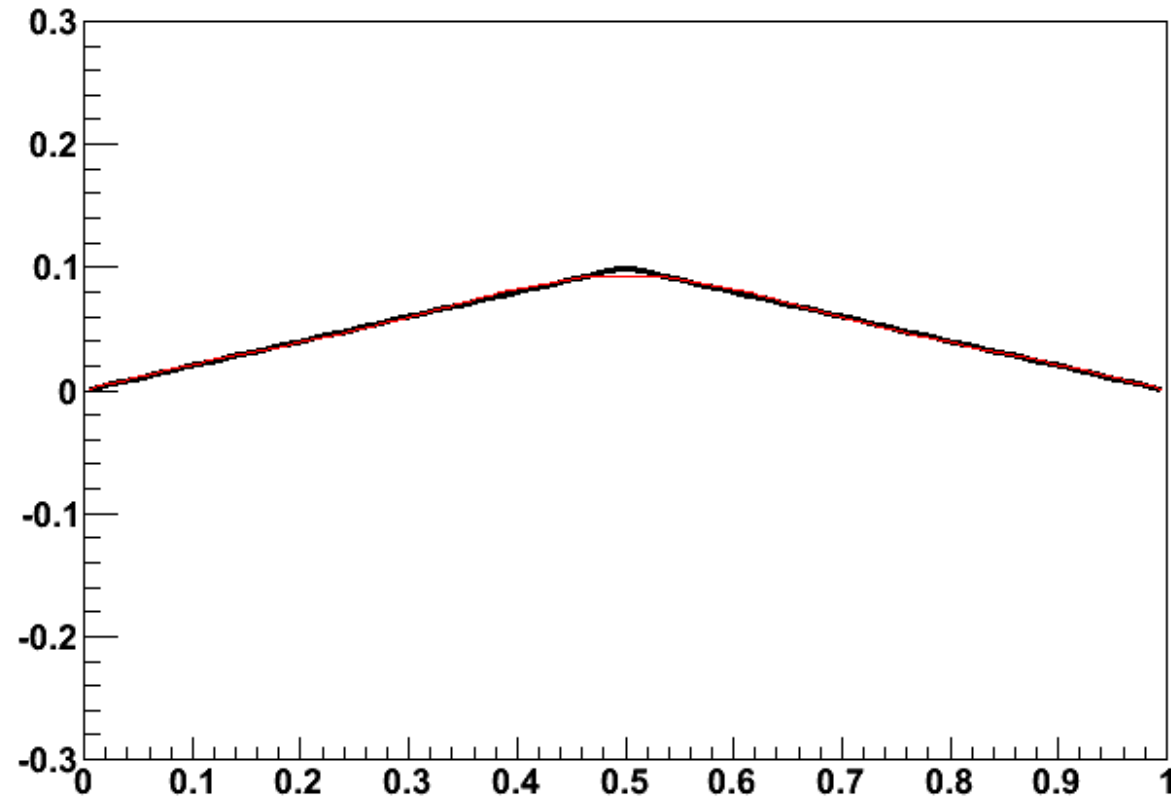
Example



$$B_1 = \frac{8h}{\pi^2} = 0.8106 h$$

$$B_3 = -\frac{8h}{9\pi^2} = -0.0901 h$$

Example



$$B_1 = \frac{8h}{\pi^2} = 0.8106 h$$

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$$B_5 = \frac{8h}{25\pi^2} = 0.0324 h$$

Example



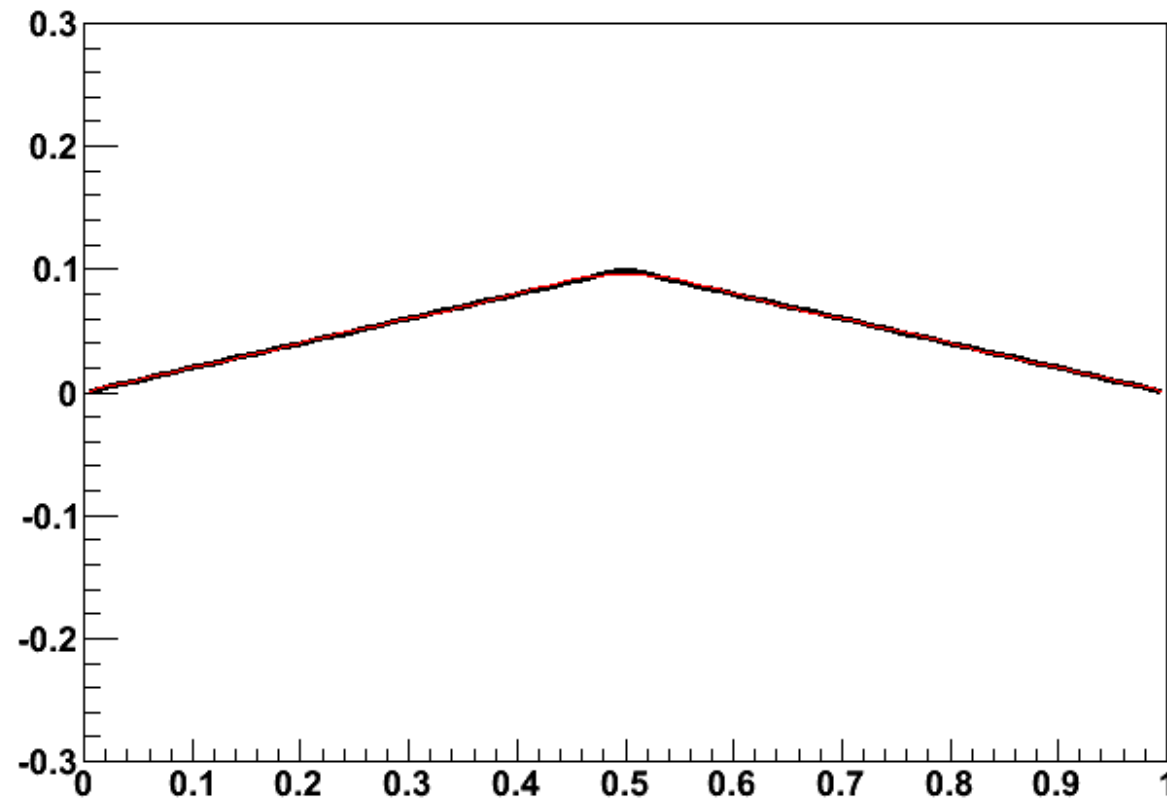
$$B_1 = \frac{8h}{\pi^2} = 0.8106 h$$

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$$B_7 = -\frac{8h}{49\pi^2} = -0.0165 h$$

Example



$$B_1 = \frac{8h}{\pi^2} = 0.8106 h$$

$$B_3 = -\frac{8h}{9\pi^2} = -0.0901 h$$

$$B_5 = \frac{8h}{25\pi^2} = 0.0324 h$$

$$B_7 = -\frac{8h}{49\pi^2} = -0.0165 h$$

$$B_9 = \frac{8h}{81\pi^2} = 0.0100 h$$

Example

$$y(x, t) = \sum_{k=1}^{\infty} B_k \sin\left(\frac{n\pi x}{L}\right) \cos(\omega_k t)$$

$$\omega_k = \frac{k\pi v}{L}$$

- Period of lowest frequency mode:

$$T = \frac{1}{\nu_1} = \frac{2L}{v}$$

- After time $t = T/2 = L/v$,

$$y(x, t) = \sum_{k=1}^{\infty} B_k \sin\left(\frac{n\pi x}{L}\right) \cos(\pi k)$$

$$= -u(x)$$

—1 because $B_k \neq 0$
only when k is odd.

- The string maintains its triangular shape!

Wave Propagation

- The wave equation:

$$\frac{\partial^2 y}{\partial x^2} = \frac{1}{v^2} \frac{\partial^2 y}{\partial t^2}$$

- We worked out solutions that satisfied specific boundary conditions.
- Previously (first week of class) we showed that a general solution is any function that is of the form

$$y(x, t) = f(x \pm vt)$$

- Are these two pictures compatible?

Wave Propagation

- Solutions for normal modes:

$$y_n(x, t) = \sin\left(\frac{n\pi x}{L}\right) \cos(\omega_n t)$$
$$\omega_n = \frac{\pi n v}{L}$$

- Trigonometric identity:

$$\sin \alpha \cos \beta = \frac{1}{2} [\sin(\alpha + \beta) + \sin(\alpha - \beta)]$$

- This gives,

$$y_n(x, t) = \frac{1}{2} \left[\sin\left(\frac{n\pi x}{L} + \omega_n t\right) + \sin\left(\frac{n\pi x}{L} - \omega_n t\right) \right]$$

Wave Propagation

$$y_n(x, t) = \frac{1}{2} \left[\sin \left(\frac{n\pi x}{L} + \omega_n t \right) + \sin \left(\frac{n\pi x}{L} - \omega_n t \right) \right]$$

- Write this as

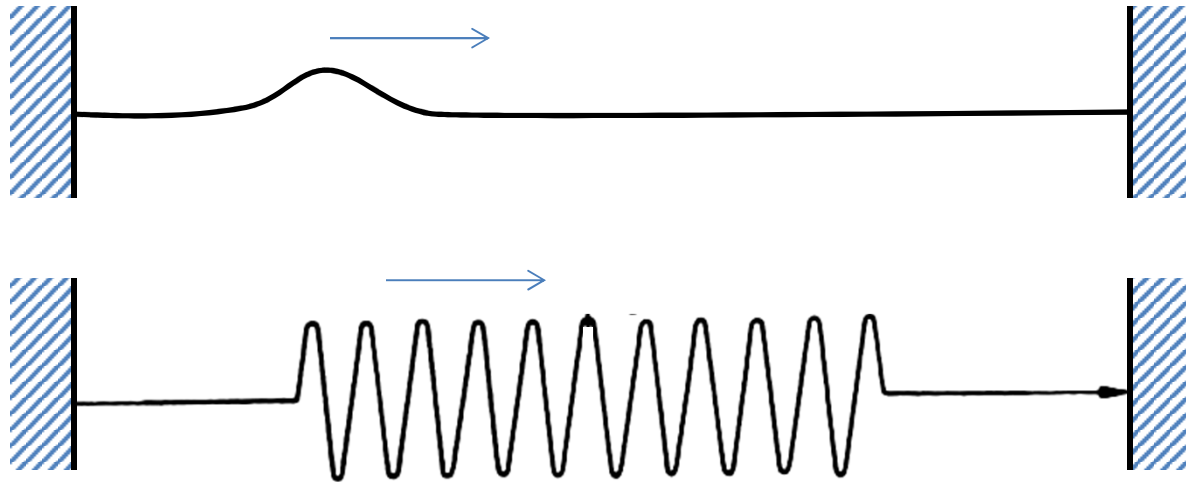
$$y_n(x, t) = \frac{1}{2} [\sin(k(x + vt)) + \sin(k(x - vt))]$$

$$k = \frac{n\pi}{L}$$

- This is the equation for two sine-waves moving in opposite directions.
- The text refers to these as “progressive waves”.
- The “standing waves” that satisfy the boundary conditions are the superposition of “progressive waves” that move in opposite directions.

Wave Propagation

- Waves can propagate in either direction.
- Easiest to visualize in terms of a pulse, or wave packet:



- If this disturbance is far from the ends, the effect is the same as letting $L \rightarrow \infty$

Wave Propagation

$$B_k = \frac{2}{L} \int_0^L \sin\left(\frac{k\pi x}{L}\right) u(x) dx$$

- In the limit where the disturbance is very far from either boundary, the Fourier sine transform is:

$$B(k) = \int_{-\infty}^{\infty} u(x) \sin(kx) dx$$

- Similarly, we can define the Fourier cosine transform:

$$A(k) = \int_{-\infty}^{\infty} u(x) \cos(kx) dx$$

- The original function is represented by:


$$u(x) = \frac{1}{\pi} \int_0^{\infty} A(k) \cos(kx) dk + \frac{1}{\pi} \int_0^{\infty} B(k) \sin(kx) dk$$

Wave Propagation

- Previously, we interpreted the coefficients B_n as the amplitude of the normal mode with frequency ω_n
 - wavelength $\lambda_n = 2L/n$
 - wavenumber $k_n = 2\pi/\lambda_n = \pi n/L$
- Now, we interpret $A(k)$ and $B(k)$ as the amplitude for harmonic waves with wavenumbers between k and $k + dk$.
- It can be important to decompose a pulse into its frequency components because in real materials, the nature of wave propagation can depend on the frequency.

Example

- Consider a pulse that has a Gaussian shape:


$$g(x) = \frac{1}{\sqrt{2\pi}} e^{-x^2/2}$$

- Special case:
 - Peak position is at $x = 0$
 - Width of the peak is $\sigma = 1$
- Other Gaussian functions can be transformed into this special case by linear change of variables.
- What is the continuous Fourier transform?

Example

$$B(k) = \int_{-\infty}^{\infty} g(x) \sin(kx) dx$$

- The Gaussian function $g(x)$ is an even function:

$$g(x) = g(-x)$$

- The function $\sin(kx)$ is an odd function:

$$\sin(-kx) = -\sin(kx)$$

- This integral must vanish...

$$B(k) = 0$$

Example

$$\begin{aligned} A(k) &= \int_{-\infty}^{\infty} g(x) \cos(kx) dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-x^2/2} \cos(kx) dx \end{aligned}$$

- From your table of integrals:

$$\int_{-\infty}^{\infty} e^{-ax^2} \cos bx dx = \sqrt{\frac{\pi}{a}} e^{-b^2/4a}$$

- In this case, $a = 1/2$ and $b = k$

$$A(k) = \frac{1}{\sqrt{2\pi}} \times \sqrt{2\pi} e^{-k^2/2} = e^{-k^2/2}$$

- This is a Gaussian distribution of wavenumbers $k = \omega/v$.