

SM matter fermion masses can be easily obtained from the superpotential - as they have only a good SUSY contribution. There will be an identical good SUSY contribution to the scalar masses, in addition to their soft-SUSY breaking contributions. In terms of the general formula the fermion mass terms have the form

$$\mathcal{L}_{\text{mass, Yuk}} = 2 \cdot \mathcal{Z}^a \frac{\delta^2 W(A)}{\delta A^a \delta A^b} \Big|_{\text{vev}} \mathcal{Z}^b + \text{h.c.}$$

(p. 53)

$$= 2H_u \cdot g_{Y_u} u^c + 2H_d \cdot g_{Y_d} d^c$$

$$+ 2H_d \cdot g_{Y_e} e^c + \text{h.c.}$$

$$= -\bar{u}_R \Gamma^{ut} H_u g_L + \bar{g}_L \Gamma^{u+} H_u u_R$$

$$- \bar{d}_R \Gamma^{dt} H_d g_L + \bar{g}_L \Gamma^{d+} H_d d_R$$

$$- \bar{e}_R \Gamma^{et} H_d l_L + \bar{l}_L \Gamma^{e+} H_d e_R$$

$$= +\bar{u}_R \Gamma^{ut} \frac{v_u}{\sqrt{2}} u_L + \bar{u}_L \Gamma^{u+} \frac{v_u}{\sqrt{2}} u_R$$

$$- \bar{d}_R \Gamma^{dt} \frac{v_d}{\sqrt{2}} d_L - \bar{d}_L \Gamma^{d+} \frac{v_d}{\sqrt{2}} d_R$$

$$- \bar{e}_R \Gamma^{et} \frac{v_d}{\sqrt{2}} e_L - \bar{e}_L \Gamma^{e+} \frac{v_d}{\sqrt{2}} e_R$$

mass  
→  
terms

$$\Rightarrow \mathcal{L}_{\text{matter mass}} \equiv -\bar{u}_L M^u u_R - \bar{d}_L M^d d_R - \bar{e}_L M^e e_R + \text{h.c.}$$

where

$$M^u \equiv -\Gamma^u \frac{N^u}{\sqrt{2}} = -\frac{1}{\sqrt{2}} \Gamma^u N \sin\beta$$

$$M^d \equiv +\Gamma^d \frac{N^d}{\sqrt{2}} = +\frac{1}{\sqrt{2}} \Gamma^d N \cos\beta$$

$$M^e \equiv +\Gamma^e \frac{N^e}{\sqrt{2}} = +\frac{1}{\sqrt{2}} \Gamma^e N \cos\beta$$

As in the SM these matrices can be diagonalized with the mixing given by the CKM matrix.

So far we have found the matter fermion masses — as in the SM & the gauge  $W^\pm, Z$  masses and the massless photon  $A_\mu$  & gluons  $G_\mu^m$  as SM used.

Let's turn next to the scalar Higgs particle masses. There are  $H_u^+, H_u^0, H_d^-, H_d^0$  — 4 complex scalar fields —> 8 independent real fields. Recall the Higgs potential  $V_H$   
 pages -55- to -56-