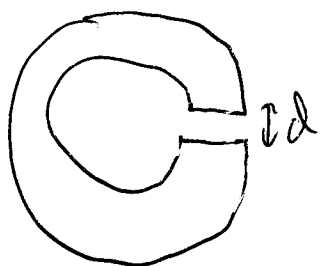


X.) i) 2. Torus with gap:



linear material $\vec{B} = \mu \vec{H}$

$$\oint_C \vec{H} \cdot d\vec{l} = NI + \underline{\vec{B}} \cdot \underline{\vec{C}}$$

$\Rightarrow B = \text{same everywhere}$

$$B = \frac{\mu \mu_0 NI}{\mu_0 2\pi r + d(\mu - \mu_0)}$$

End of Statics

XI.) Electromagnetic Induction:

$$\frac{\partial \vec{B}}{\partial t} \neq 0.$$

a) $\mathcal{E}_{\text{emf}} = \mathcal{E} = \oint_C \vec{E} \cdot d\vec{l} \neq 0$ when $\frac{\partial \vec{B}}{\partial t} \neq 0.$

b.) Faraday's law new exp. law: 1831

$$\mathcal{E} = - \frac{d\phi}{dt}$$

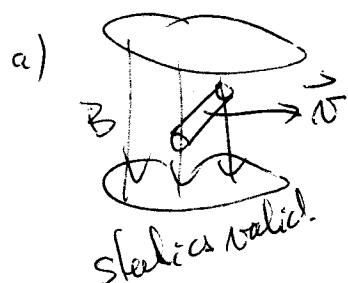
$$\phi = \int_S \vec{B} \cdot d\vec{S} \quad \text{integral form.}$$

$$\vec{\nabla} \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}$$

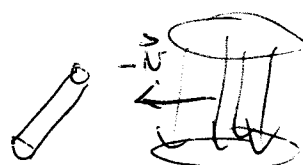
diff. form.

(- sign = Lenz's law)

XI.c) ex. 1) Wire moving with $\vec{v}$ in $\vec{B}$



b.)



use Faraday's law to get same result.

d) Self - & Mutual - Inductance

1) $\phi = \phi(I)$ so $\frac{d\phi}{dt} = \frac{d\phi}{dI} \frac{dI}{dt}$
 rigid circuit.

Single circuit.

$L = \frac{d\phi}{dI}$ self-ind.

$\mathcal{E} = -L \frac{dI}{dt}$

ex. toroidal coil



$L = \frac{\mu_0 N^2 A}{l}$

2) many circuits: $\phi_i = \sum_{j=1}^n \phi_{ij}$

$\frac{d\phi_{ij}}{dt} = \frac{d\phi_{ij}}{dI_j} \frac{dI_j}{dt} \equiv M_{ij} \frac{dI_j}{dt} = -\mathcal{E}_i$
 $\downarrow$
 mutual Ind.

$M_{ii} = L_i$

XI. d.2) ex. twice wrapped toroid coil.

$$M_{12} = M_{21} = \frac{\mu_0 N_1 N_2 A}{l}$$

3.) Neumann Formula: $M_{ij} = M_{ji}$

B.-S. law $\Rightarrow$

$$M_{12} = \frac{\mu_0}{4\pi} \oint_{C_1} \oint_{C_2} \frac{d\vec{l}_1 \cdot d\vec{l}_2}{|\vec{r}_{12} - \vec{r}_{11}|}$$

c.) Magnetic Energy

1) Electrostatic: $U = \frac{1}{2} \int \vec{D} \cdot \vec{E} dV$

2) $dU = -E dq$ $\oint dq = \vec{J} \cdot d\vec{S} dt$; $q = \int_S \vec{J} \cdot d\vec{S} dt$

$$dU = - \int E \vec{J} \cdot d\vec{S} dt$$

$$= \left(\int_S \vec{J} \cdot d\vec{S} \frac{d\phi}{dt} \right) dt$$

Now $\vec{J}(t) = \begin{cases} \frac{1}{T} \vec{J} & 0 \leq t \leq T \\ \vec{J} & t \geq T \end{cases}$

But $\phi \sim \vec{J}$ B.-S. law

$$\Rightarrow \phi(t) = \begin{cases} \frac{1}{T} \phi & 0 \leq t \leq T \\ \phi & t \geq T \end{cases}$$

XI.) e.2.) S_0

$$U = \int_0^T \left(\int_S \frac{\partial \phi}{\partial t} \vec{J} \cdot d\vec{s} \right) dt$$

$$\boxed{U = \frac{1}{2} \int_S \phi \vec{J} \cdot d\vec{s}} = \frac{1}{2} \vec{J} \cdot \vec{I} \cdot \phi_i$$

$$\text{But } \phi = \int_{\Gamma} \vec{B} \cdot d\vec{s} = \int_{\partial \Gamma} \vec{A} \cdot d\vec{l}$$

$$S_0 \int_S \phi \vec{J} \cdot d\vec{s} = \int_S \int_{\partial \Gamma} \vec{A} \cdot d\vec{l} \vec{J} \cdot d\vec{s} = \int_V \vec{J} \cdot \vec{A} dV$$

$$\Rightarrow \boxed{U = \frac{1}{2} \int_V \vec{J} \cdot \vec{A} dV}$$

$$\text{finally } \vec{J} = \vec{\nabla} \times \vec{H} \quad \vec{B} = \vec{\nabla} \times \vec{A}$$

$$\Rightarrow \boxed{U = \frac{1}{2} \int_V \vec{B} \cdot \vec{H} dV}$$

XII.) Maxwell's equations

a) Not only can $\frac{\partial \vec{B}}{\partial t} \neq 0$ but also $\vec{\nabla} \cdot \vec{J} \neq 0$.

$$\text{Now } \vec{\nabla} \cdot \vec{J} = -\frac{\partial \rho}{\partial t} = -\vec{\nabla} \cdot \frac{\partial \vec{B}}{\partial t} \neq 0,$$

$$\text{Since } \vec{\nabla} \cdot \vec{B} = \rho.$$

XII.) a) So, Ampere's law needs to be modified
consistent with charge conservation
i.e. continuity eq. $\vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} = 0$.

i.e. $\vec{\nabla} \times \vec{H} = \vec{J}$ take $\vec{\nabla} \cdot$; $\vec{\nabla} \cdot \vec{\nabla} \times \vec{H} \equiv 0$!
 $\Rightarrow \vec{\nabla} \cdot \vec{J} = 0$!
 not true always.

So

Maxwell said

$$\boxed{\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}}$$

modified form of
Ampere's law.

So

$$0 = \vec{\nabla} \cdot \vec{\nabla} \times \vec{H} = \vec{\nabla} \cdot \vec{J} + \frac{\partial}{\partial t} \vec{\nabla} \cdot \vec{D} = \vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} !$$

Charge conservation.

b.) All E-M phenomena described by
the 4 Maxwell's equations:

1) $\vec{\nabla} \cdot \vec{D} = \rho$: Gauss's law

2) $\vec{\nabla} \cdot \vec{B} = 0$: no magnetic monopoles

3) $\vec{\nabla} \times \vec{E} = - \frac{\partial \vec{B}}{\partial t}$: Faraday's law

4) $\vec{\nabla} \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t}$: Ampere's law.

XI) b.) In addition we need def. of $\vec{D}$ & $\vec{H}$

$$\begin{array}{ll}
 1) \vec{D} = \epsilon_0 \vec{E} + \vec{P} & \vec{P}(\vec{E}) \\
 2) \vec{H} = \frac{1}{\mu_0} \vec{B} - \vec{M} & \vec{M}(\vec{B}) \\
 3) \vec{J} = \vec{J}(\vec{E}) & \vec{J}(\vec{E})
 \end{array}
 \left. \begin{array}{l} \\ \\ \end{array} \right\} \begin{array}{l} \text{we further} \\ \text{need} \end{array}$$

For linear, isotropic material

$$\begin{aligned}
 \vec{D} &= \epsilon \vec{E} \\
 \vec{B} &= \mu \vec{H}
 \end{aligned}$$

$$\vec{J} = \sigma \vec{E} \quad (\text{Ohmic})$$

c.) Matter, i.e. charges, react to E-M fields via the Lorentz force law

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B})$$

UA