

Multiphase Transport Phenomena

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- > Interdisciplinary research and education challenges
- Learning theory and Zen on learning and confusion
- Engagement
- Build on background

Approach

- Connect to other lectures
- > Some things classical and some things not
- Facilitate abstraction
- Provide for the needy

Scales: molecular, micro, macro, meso, and mega

- Phases: single-fluid phase, multiple fluid phase systems
- Flow versus species transport
- Continuum scale modeling

Some Terminology

- >Closure, or constitutive, relations
- >Forward versus inverse problems

Overview

- Modeling process
- Conventional macroscale, or porous medium continuum scale, modeling approach
- >Macroscale multiphase flow
- >Macroscale multiphase flow and transport
- >Multiscale modeling
- Microscale modeling approaches
- >Example applications









Conservation Equations

- Averaging procedures were alluded to as a means to develop and close the single-phase flow equation
- •Averaging procedures were also discussed as a means to develop a species transport equation

Species Balance Equation:

 $\frac{\partial}{\partial t} \left(\epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \right) = -\nabla \cdot (\mathbf{j}^{\iota \alpha} + \epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\iota \alpha} + \mathcal{R}^{\iota \alpha} + \mathcal{S}^{\iota \alpha}$

Species-Summed Flow Equation: $\frac{\partial}{\partial t} (\epsilon^{\alpha} \rho^{\alpha}) = -\nabla \cdot (\epsilon^{\alpha} \rho^{\alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\alpha} + S^{\alpha}$

What set of constraints can be developed with respect to the quantities that appear in the species conservation equation?

Ponderables

➢ Using these constraints, show how the flow equation can be derived from the species mass conservation equation

Conservation Equations and Constraints

Species Balance Equation:

 $\frac{\partial}{\partial t} \left(\epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \right) = -\nabla \cdot (\mathbf{j}^{\iota \alpha} + \epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\iota \alpha} + \mathcal{R}^{\iota \alpha} + \mathcal{S}^{\iota \alpha}$

Species-Summed Flow Equation:

$$\frac{\partial}{\partial t} (\epsilon^{\alpha} \rho^{\alpha}) = -\nabla \cdot (\epsilon^{\alpha} \rho^{\alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\alpha} + S^{\alpha}$$
$$\sum_{\alpha} \epsilon^{\alpha} = \mathbf{1}, \ \sum_{\alpha} \mathcal{I}^{\iota \alpha} = \mathbf{0},$$
$$\sum_{\iota} \omega^{\iota \alpha} = \mathbf{1}, \ \sum_{\iota} \mathbf{j}^{\iota \alpha} = \mathbf{0}, \ \sum_{\iota} \mathcal{R}^{\iota \alpha} = \mathbf{0}$$
$$\sum_{\iota} \mathcal{I}^{\iota \alpha} = \mathcal{I}^{\alpha} \ \sum_{\iota} S^{\iota \alpha} = S^{\alpha}$$

For single-phase flow, what is the closure problem?

Ponderables

- Consider Darcy's experiments and law and use as an alternative to computational or theoretical approaches to solve the closure problem and note all other assumptions used
- Derive the single-phase flow model for the case in which porosity is constant and the fluid is incompressible, note the consequences, and assess if these are reasonable

Conservation Equations and Constraints

Species Balance Equation:

 $\frac{\partial}{\partial t} \left(\epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \right) = -\nabla \cdot (\mathbf{j}^{\iota \alpha} + \epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\iota \alpha} + \mathcal{R}^{\iota \alpha} + \mathcal{S}^{\iota \alpha}$

Species-Summed Flow Equation:

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$$\sum_{\iota} \mathcal{I}^{\iota \alpha} = \mathcal{I}^{\alpha} \ \sum_{\iota} S^{\iota \alpha} = S^{\alpha}$$



Traditional Single-Phase Flow Model

$$\frac{\partial}{\partial t} (\epsilon^{\alpha} \rho^{\alpha}) = -\nabla \cdot (\epsilon^{\alpha} \rho^{\alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\alpha} + S^{\alpha}$$
$$\rho^{a} = \rho^{a} (p^{a})$$
$$\mathbf{q} = -\mathbf{K} \cdot \nabla h$$
$$\partial h$$

$$S_s \frac{\partial n}{\partial t} = \nabla \cdot (\mathbf{K} \cdot \nabla h) + S$$



Equilibrium state





$$p_n - p_w = \frac{2\gamma}{R} \cos \theta$$

Physics of Multiphase Porous Medium Systems

- In a multiphase porous medium system, fluids move in response to viscous, capillary, and gravity forces
- This balance of forces is influenced by properties of the medium and the fluids: morphology of the pore space, contact angle, interfacial tensions, densities, and viscosities
- These forces result in very complex patterns of flow and entrapment of residual nonwetting phases can result

Multiphase Flow and Species Transport

Multiphase flow---more than one fluid occupying the pore space

- Water infiltration
- Short time scale NAPL infiltration
- Petroleum exploration

Multiphase flow and species transport---more than one fluid and species mass fractions or concentrations are important

- Pesticide transport
- > BTEX problems from petroleum spills

DNAPL Behavior in Heterogeneous Porous Media







Species Balance Equation:

Multiphase Conservation Equations

 $\frac{\partial}{\partial t} \left(\epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \right) = -\nabla \cdot (\mathbf{j}^{\iota \alpha} + \epsilon^{\alpha} \rho^{\alpha} \omega^{\iota \alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\iota \alpha} + \mathcal{R}^{\iota \alpha} + \mathcal{S}^{\iota \alpha}$

Species-Summed Flow Equation:

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- > Define the closure problem for two-phase flow in a macroscale porous medium system
- How might one investigate approaches to produce a closed model?
- > What sorts of assumptions are implicit in traditional closure approaches?



Assuming an immobile solid phase

 $\frac{\partial}{\partial t} \left(\epsilon^{\alpha} \rho^{\alpha} \right) = -\nabla \cdot \left(\epsilon^{\alpha} \rho^{\alpha} \mathbf{v}^{\alpha} \right) + \mathcal{I}^{\alpha} + \mathcal{S}^{\alpha}$

For a two-fluid system, this results in five unknowns for each phase or a total of ten unknowns in two equations

Capillary Pressure Saturation Relations



•C-109 sand experiment •Key features to note:

- •Primary drainage
- •Entry pressure
- •Uniformity effects
- Main imbibition
- •Non-wetting phase trapped
- •Wetting scanning curves
- •Hysteresis
- •Quasi-static experiments

Examples of Common Closure Assumptions

- Saturation is solely determined based upon capillary pressure and its history in a quasi-static sense
- Darcy's law can be extended with modification to multiphase systems
- Relative permeability is solely dependent upon the saturation of the respective phase and its history
- Rigorous connection with microscale quantities can be ignored

Write a general form for a closed multiphase flow model

Ponderables

Consider an air-water system and write a closed model for this special case noting the reasoning steps used in the simplification Closed Multiphase Flow Model

$$\begin{split} \frac{\partial}{\partial t} (\epsilon^{\alpha} \rho^{\alpha}) &= -\nabla \cdot (\epsilon^{\alpha} \rho^{\alpha} \mathbf{v}^{\alpha}) + \mathcal{I}^{\alpha} + \mathcal{S}^{\alpha} \\ \mathbf{q}^{\alpha} &= \epsilon^{\alpha} \mathbf{v}^{\alpha} = -\frac{kk^{r\alpha}}{\mu^{\alpha}} \nabla \left(p^{\alpha} + \rho^{\alpha} g z \right) \\ S^{\alpha} &= f \left(p^{\beta}(t) \right), \text{ for } \beta = 1, ..., n_{f} \\ k^{r\alpha} &= f \left(S^{\beta}(t) \right), \text{ for } \beta = 1, ..., n_{f} \\ \mathbf{q}^{\alpha} &= \epsilon^{\alpha} \mathbf{v}^{\alpha} = -\frac{\mathbf{k}^{\alpha}}{\mu^{\alpha}} \cdot \left(\nabla p^{\alpha} + \rho^{\alpha} g \nabla z \right) \\ \rho^{\alpha} &= \rho^{\alpha}(p^{\alpha}) \\ \sum_{\alpha} \epsilon^{\alpha} &= 1 \end{split}$$

Example Closure Relations

van Genuchten P-S Relation:

$$S_e = \frac{\epsilon^a - \epsilon^r}{\epsilon^s - \epsilon^r} = (1 + |\alpha_v \psi|^{n_v})^{-m_v} \text{ for } \psi < 0$$
$$= 1 \text{ for } \psi \ge 0$$

Mualem S-K Relation:

$$k^{rw}(S_e) = S_e^{1/2} \left[1 - \left(1 - S_e^{1/m_v} \right)^{m_v} \right]^2$$

$$k^{rn}(S_e) = (1 - S_e)^{1/2} \left(1 - S_e^{1/m_v} \right)^{2m_v}$$

Air is much more mobile than water, therefore pressure gradients must be small for the air phase-assume zero

- Porosity is assumed constant, thus changes in water and air volume fractions are inversely related
- Common multiphase extension of Darcy's law applies
- Quasi-static pressure-saturation-relative permeability relations apply

Air-Water System

Spatial gradients of aqueous-phase density can be ignored

$$\widehat{Pichards' Equation}$$

$$\frac{\partial}{\partial t} (\epsilon^{a} \rho^{a}) = -\nabla \cdot (\rho^{a} q^{a})$$

$$\left(\frac{\epsilon^{a}}{\rho^{a}}\right) \frac{\partial \rho^{a}}{\partial t} + \frac{\partial \epsilon^{a}}{\partial t} = -\nabla \cdot q^{a}$$

$$S_{s} S^{a} \frac{\partial \psi}{\partial t} + \frac{\partial \epsilon^{a}}{\partial t} = \nabla \cdot [K^{a} \nabla (\psi + z)]$$

$$S_{s} S^{a} \frac{\partial \psi}{\partial t} + \frac{\partial \epsilon^{a}}{\partial t} = \frac{\partial}{\partial z} \left[K^{a} \left(\frac{\partial \psi}{\partial z} + 1\right)\right]$$

$$\left(S_{s} S^{a} + \frac{\partial \epsilon^{a}}{\partial \psi}\right) \frac{\partial \psi}{\partial t} = \frac{\partial}{\partial z} \left[K^{a} \left(\frac{\partial \psi}{\partial z} + 1\right)\right]$$

> Often the problem of concern

Commonality with single-phase systems that transport model requires solution of the flow model for closure

Multiphase Flow and Transport

Commonality with single-phase flow model as well for implications of reaction form on size and formal type of resultant system of conservation equations

Ponderables

- Formulate a model to describe the transport and fate of contaminants resulting from the spill of a refined petroleum product in the unsaturated zone
- > Describe the flow model
- > What transport processes are of concern?
- Without simplifying assumptions, what is the size of the system of conservation equations?

How can separation of time scales be used to simplify the system of equations?

Ponderables

- > How can the number of species be reduced?
- What assumptions are used to support the notion of natural attenuation for this class of problem?

Examples of Current Multiphase Research

- > Multiscale inspired
- Single-phase flow
- Pressure-saturation relations
- > NAPL dissolution fingering
- Viscous coupling of fluids
- Multiscale NAPL dissolution
- > Thermodynamically constrained averaging theory
- DNAPL remediation revisited

Multiscale Porous Medium Systems





Pore scale



Lab scale

Field scale



A predictive tool to determine constitutive relations for standard continuum-scale models

A significant means to close new continuumscale theories for multiphase flow

An important way of understanding the fundamental pore-scale processes











0.5

- 0.5

Pore-network models



Simulate fluids as microscopic particles that move along a lattice and collide with each other

> Fully recover Navier-Stokes equation

Relatively easy implementation of boundary condition on complex geometries

> Suitable for massively parallel computers


$$f_i(\vec{x} + \vec{e}_i, t+1) - f_i(\vec{x}, t) = \frac{1}{\tau} \Big[f_i^{(eq)}(\vec{x}, t) - f_i(\vec{x}, t) \Big]$$

 $\rho' = \sum f_i$: density

 $\rho' \vec{u} = \sum_{i} f_i \vec{e}_i$: momentum





| | RSP1 | RSP23 |
|---------------------------|-------|-------|
| < <i>D</i> > (mm) | 0.20 | 0.19 |
| $\sigma < D >$ | 0.5% | 66% |
| L(mm) | 4.250 | 6.192 |
| φ | 0.442 | 0.334 |
| $\mathbf{N}_{\mathbf{s}}$ | 10328 | 14380 |



























Onset of Non-Darcy flow







Two-Phase Pressure Saturation Relations



•C-109 sand experiment •Key features to note:

- •Primary drainage
- •Entry pressure
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- •Non-wetting phase trapped
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Experimental and Simulated Properties

| | GB1b | Simulated GB1b |
|----------------------------|-----------------|----------------|
| <d> (mm)</d> | 0.1156 | 0.1149 |
| $\sigma_{\rm D}({\rm mm})$ | 0.0121 | 0.0116 |
| Porosity \phi | 0.356 ± 0.002 | 0.356 |
| Domain Length L (mm) | 2.35 | 2.35 |
| Number of Spheres | | 9532 |
| NWP-WP | Dyed PCE- water | |

<D>: Arithmetic mean diameter

 σ_{D} : Arithmetic standard deviation of grain diameter

Density ratio between fluids

- Viscosity ratio
- Interfacial tension (fluid fluid interaction)

Calibration of LB Multiphase Model

- Wettability (fluid solid interaction)
- Boundary conditions



Initial state

Equilibrium state







Interfacial tension force profile

Pressure profile









Laplace's Law:

$$p_n - p_w = \frac{2\gamma}{R} \cos\theta$$

















Micromodel TCE Residual

- •Two-dimensional glass bead micromodel
- •TCE dyed with Oil Red O
- •Water saturated followed by DNAPL displacement and then water flushing
- •TCE residual saturation results
- •Large range of sizes of trapped TCE
- •Largest features contain the majority of the TCE mass and are the most difficult to remove





















































NAPL Dissolution Tailing for TCE



•Column brought to residual saturation with TCE

•Water flusing in an attempt to obtain drinking water standard concentrations of TCE

•Large TCE residual feature determines clean-up time

•Eventually complex TCE region breaks up and drinking water standards reached

•Reference: Imhoff et al. [ES&T, 32(16), 1998]



$\sum \text{Linear model}$ MassFlux=k_la_{na}(C_s-C)

$$K_{I}=k_{I}a_{na}$$

> Important dimensionless groupings

$$Sh = \frac{K_1 d_p^2}{D_m}, Re = \frac{v_a \rho_a d_p}{\mu_a}, Sc = \frac{\mu_a}{\rho_a D_m}, \theta_n$$






➢ As the dissolution front moved downward through the medium, preferential flow paths developed.



Dissolution Experiment Results



t = 0 hrs





Dissolution Experiment: Simulated Results















IMPERMEABLE LAYER

 $\log(\theta_n)$



DNAPLs in Heterogeneous Systems





Nonaqueous Phase Dissolution





Relative permeabilitysaturation (*ks*) relation:

Viscous Coupling

Viscous coupling effect







Deficiencies in Traditional Models

- Models are often based upon ill-defined variables and empirically derived closure approximations lacking in theoretical support and precise knowledge of limitations
- Empirical closure relations are routinely extended beyond their level of experimental support

Deficiencies in Traditional Models

- Rigorous linkages among scales is usually absent
- Important phenomena are often not included naturally in multiphase models (e.g., wettability)
- Standard porous medium models are not typically constrained to obey the second law of thermodynamics

Deficiencies in Traditional Models

- Quantities of interest, such as interfacial areas, do not explicitly arise in standard models
- Standard models are often built upon assumptions well-known to be violated (e.g. quasi-equilibrium states)



Standard approaches lack a rigorous structure in which to examine simplifying assumptions

Form general conservation equations

TCAT Approach

- Use general conservation equations to formulate specific conservation equations for mass, momentum, angular momentum, energy, and entropy
- Specify thermodynamic dependence of internal energy and independent system variables

> Derive a total system entropy inequality

TCAT Approach

- Constrain the entropy inequality with the product of Lagrange multipliers and specific conservation equations, thermodynamic relations, and other constraints
- Solve for Lagrange multipliers to simplify entropy inequality

TCAT Approach

- Exploit entropy inequality to guide development of closure relations
- Use sub-scale theory, computation, or experiment to guide final form of closure relations
- Compare model systems to experimental observations and use to guide experimental design

$$\mathcal{E} = \{\Omega_{\iota} | \iota \in \mathcal{I}\} = \{\Omega_{w}, \Omega_{s}, \Omega_{ws}\}$$
$$\mathcal{I} = \{w, s, ws\}$$
$$\mathcal{E}_{c\iota} = \{\Omega_{\kappa} | (\bar{\Omega}_{\iota} \cap \bar{\Omega}_{\kappa} \neq \emptyset) \land (\bar{\Omega}_{\iota} \neq \bar{\Omega}_{\kappa}), \forall \Omega_{\kappa} \in \mathcal{E}\}$$
$$\mathcal{E}_{cw} = \{\Omega_{ws}\}$$
$$\mathcal{E}_{cs} = \{\Omega_{ws}\}$$
$$\mathcal{E}_{cws} = \{\Omega_{w}, \Omega_{s}\}$$

$$\begin{split} \begin{array}{l} \overbrace{D^{\bar{\imath}}\eta^{\bar{\imath}}}{Dt} + \eta^{\bar{\imath}}\mathbf{l}:\mathbf{d}^{\bar{\imath}} - \nabla \cdot \left(\epsilon^{\iota}\varphi^{\bar{\imath}}\right) - \epsilon^{\iota}b^{\iota} - \sum_{\kappa \in \mathfrak{I}_{c\iota}} \left(\stackrel{\kappa \to \iota}{M_{\eta}} + \stackrel{\kappa \to \iota}{\Phi}\right) = \Lambda^{\bar{\imath}}, \quad \text{for } \iota \in \mathfrak{I} \\ \langle \mathfrak{P}_{i} \rangle_{\Omega_{j},\Omega_{k},w} &= \frac{\int_{\Omega_{j}} w \mathfrak{P}_{i} \, \mathrm{d}\mathfrak{r}}{\int_{\Omega_{k}} w \, \mathrm{d}\mathfrak{r}} & \eta^{\bar{\imath}} = \langle \eta_{\iota} \rangle_{\Omega_{\iota},\Omega} \\ \mathbf{v}^{\bar{\imath}} = \langle \mathbf{v}_{\iota} \rangle_{\Omega_{\iota},\Omega_{\iota},\rho_{\iota}} \\ \mathbf{d}^{\bar{\imath}} = \frac{1}{2} \left[\nabla \mathbf{v}^{\bar{\imath}} + \left(\nabla \mathbf{v}^{\bar{\imath}} \right)^{\mathrm{T}} \right] \\ \rho^{\bar{\imath}} = \langle \varphi_{\iota} \rangle_{\Omega_{\iota},\Omega_{\iota}} - \left\langle \eta_{\iota} \left(\mathbf{v}_{\iota} - \mathbf{v}^{\bar{\imath}} \right) \right\rangle_{\Omega_{\iota},\Omega_{\iota}} \\ b^{\iota} = \langle b_{\iota} \rangle_{\Omega_{\iota},\Omega} \\ \Lambda^{\bar{\imath}} = \langle \Lambda_{\iota} \rangle_{\Omega_{\iota},\Omega} \end{split}$$

TCAT---Conservation Equations

$$\mathcal{M}^{\iota} = \frac{\mathbf{D}^{\overline{\iota}} \left(\epsilon^{\iota} \rho^{\iota}\right)}{\mathbf{D}t} + \epsilon^{\iota} \rho^{\iota} \mathbf{l} : \mathbf{d}^{\overline{\overline{\iota}}} - \sum_{\kappa \in \mathfrak{I}_{\mathbf{c}\iota}} \overset{\kappa \to \iota}{M} = 0, \quad \text{for } \iota \in \mathfrak{I}$$

$$\begin{aligned} \boldsymbol{\mathcal{P}}^{\iota} &= \frac{\mathbf{D}^{\overline{\iota}} \left(\epsilon^{\iota} \rho^{\iota} \mathbf{v}^{\overline{\iota}} \right)}{\mathbf{D} t} + \epsilon^{\iota} \rho^{\iota} \mathbf{v}^{\overline{\iota}} \mathbf{l} : \mathbf{d}^{\overline{\iota}} - \nabla \boldsymbol{\cdot} \left(\epsilon^{\iota} \mathbf{t}^{\overline{\iota}} \right) - \epsilon^{\iota} \rho^{\iota} \mathbf{g}^{\overline{\iota}} \\ &- \sum_{\kappa \in \mathfrak{I}_{\mathrm{c}\iota}} \left(\begin{pmatrix} \kappa \rightarrow \iota \\ \mathbf{M}_{v} + \mathbf{T}^{\iota} \end{pmatrix} = 0, \quad \text{for } \iota \in \mathfrak{I} \end{aligned}$$

$$\begin{split} \mathcal{E}^{\iota} &= \frac{\mathbf{D}^{\overline{\iota}} \left[E^{\overline{\overline{\iota}}} + \epsilon^{\iota} \rho^{\iota} \left(\frac{1}{2} \mathbf{v}^{\overline{\iota}} \cdot \mathbf{v}^{\overline{\iota}} + K_{E}^{\overline{\overline{\iota}}} + \psi^{\overline{\iota}} \right) \right]}{\mathbf{D}t} \\ &+ \left[E^{\overline{\iota}} + \epsilon^{\iota} \rho^{\iota} \left(\frac{1}{2} \mathbf{v}^{\overline{\iota}} \cdot \mathbf{v}^{\overline{\iota}} + K_{E}^{\overline{\iota}} + \psi^{\overline{\iota}} \right) \right] \mathbf{l} : \mathbf{d}^{\overline{\iota}} - \nabla \cdot \left(\epsilon^{\iota} \mathbf{t}^{\overline{\iota}} \cdot \mathbf{v}^{\overline{\iota}} + \epsilon^{\iota} \mathbf{q}^{\overline{\iota}} \right) \\ &- \epsilon^{\iota} h^{\iota} - \sum_{\kappa \in \mathfrak{I}_{c\iota}} \left(K_{E}^{\kappa \to \iota} + K_{v}^{\kappa \to \iota} + K_{v}^{\epsilon \to \iota} \right) = 0, \quad \text{for } \iota \in \mathfrak{I} \end{split}$$

$$\begin{split} \mathbf{\nabla} \mathbf{T} \mathbf{C} \mathbf{A} \mathbf{T} \textbf{---Averaged Thermodynamics}} \\ \mathcal{V} &= \left\{ \epsilon^{\iota}, \rho^{\iota}, \mathbf{v}^{\overline{\iota}}, \overset{ws \to \kappa}{M}, \mathbf{t}^{\overline{t}}, \overset{ws \to \kappa}{\mathbf{T}}, E^{\overline{t}}, K_{E}^{\overline{t}}, \psi^{\overline{\iota}}, \mathbf{q}^{\overline{t}}, h^{\iota}, \overset{ws \to \kappa}{Q} \right\} \\ \mathcal{T}^{w} &= \frac{\mathbf{D}^{\overline{w}} E^{\overline{w}}}{\mathbf{D} t} - \theta^{\overline{w}} \frac{\mathbf{D}^{\overline{w}} \eta^{\overline{w}}}{\mathbf{D} t} - \mu^{\overline{w}} \frac{\mathbf{D}^{\overline{w}} (\epsilon^{w} \rho^{w})}{\mathbf{D} t} + p^{w} \frac{\mathbf{D}^{\overline{w}} \epsilon^{w}}{\mathbf{D} t} \\ &+ \left\langle \eta_{w} \frac{\mathbf{D}^{\overline{w}} \left(\theta_{w} - \theta^{\overline{w}}\right)}{\mathbf{D} t} + \rho_{w} \frac{\mathbf{D}^{\overline{w}} \left(\mu_{w} - \mu^{\overline{w}}\right)}{\mathbf{D} t} - \frac{\mathbf{D}^{\overline{w}} \left(p_{w} - p^{w}\right)}{\mathbf{D} t} \right\rangle_{\Omega_{w},\Omega} = 0 \\ \mathbf{v}^{\overline{w}} &= \mathbf{v}^{\overline{s}} = \mathbf{v}^{\overline{ws}} = \text{constant} \\ &\theta^{\overline{w}} = \theta^{\overline{s}} = \theta^{\overline{ws}} = \text{constant} \\ &\frac{\mathbf{D}^{\overline{w}} \epsilon^{w}}{\mathbf{D} t} = \frac{\mathbf{D}^{\overline{s}} \epsilon^{s}}{\mathbf{D} t} = \frac{\mathbf{D}^{\overline{ws}} \epsilon^{ws}}{\mathbf{D} t} = 0 \\ &\langle p_{w} + \mathbf{n}_{s} \cdot \mathbf{t}_{s} \cdot \mathbf{n}_{s} + \gamma_{ws} \nabla' \cdot \mathbf{n}_{s} - \rho_{ws} \mathbf{g}_{ws} \cdot \mathbf{n}_{s} \rangle_{\Omega_{ws},\Omega_{ws}} = 0 \\ &\Lambda &= \sum_{i} J_{i} F_{i} + \sum_{j} \mathbf{J}_{j} \cdot \mathbf{F}_{j} + \sum_{k} \mathbf{J}_{k} : \mathbf{F}_{k} \end{split}$$



 $\sum \left(\mathcal{S}^{\iota} + \lambda^{\iota}_{\mathcal{M}} \mathcal{M}^{\iota} + \boldsymbol{\lambda}^{\iota}_{\boldsymbol{\mathcal{P}}} \boldsymbol{\cdot} \boldsymbol{\mathcal{P}}^{\iota} + \lambda^{\iota}_{\mathcal{E}} \mathcal{E}^{\iota} + \lambda^{\iota}_{\mathcal{T}} \mathcal{T}^{\iota} \right) = \Lambda \geq 0$ $\iota \in \mathcal{I}$

$$\begin{split} \sum_{\iota \in \mathfrak{I}} & \left[\mathcal{S}_{D}^{\iota} + \frac{1}{\theta^{\overline{\imath}}} \left(K_{E}^{\overline{\imath}} + \mu^{\overline{\imath}} + \psi^{\overline{\imath}} - \frac{\left(\mathbf{v}^{\overline{\imath}} \cdot \mathbf{v}^{\overline{\imath}} \right)}{2} \right) \mathcal{M}_{D}^{\iota} \right] \\ & + \sum_{\iota \in \mathfrak{I}} & \left[\frac{\mathbf{v}^{\overline{\imath}}}{\theta^{\overline{\imath}}} \cdot \mathcal{P}_{D}^{\iota} - \frac{1}{\theta^{\overline{\imath}}} \left(\epsilon^{\iota} \rho^{\iota} \frac{\mathbf{D}^{\overline{\imath}} \left(K_{E}^{\overline{\imath}} + \psi^{\overline{\imath}} \right)}{\mathbf{D}t} + \mathcal{E}_{D}^{\iota} \right) \right] \\ & + \frac{1}{\theta^{\overline{w}}} \left(p^{w} \frac{\mathbf{D}^{\overline{w}} \epsilon^{w}}{\mathbf{D}t} + \mathcal{T}_{r}^{w} \right) + \frac{\mathcal{T}_{r}^{s}}{\theta^{\overline{s}}} - \frac{1}{\theta^{\overline{ws}}} \left(\gamma^{ws} \frac{\mathbf{D}^{\overline{ws}} \epsilon^{ws}}{\mathbf{D}t} - \mathcal{T}_{r}^{ws} \right) = \Lambda \ge 0 \end{split}$$

$$\begin{aligned} & \frac{e^{w}}{\theta^{\overline{w}}} \left(\mathbf{t}^{\overline{w}} + p^{w} \mathbf{I} \right) : \mathbf{d}^{\overline{w}} + \frac{e^{s}}{\theta^{\overline{s}}} \left(\mathbf{t}^{\overline{s}} - \mathbf{t}^{s} \right) : \mathbf{d}^{\overline{s}} \\ & + \frac{e^{w}}{\theta^{\overline{w}s}} \left[\mathbf{t}^{\overline{ws}} - \gamma^{ws} \left(\mathbf{I} - \mathbf{G}^{ws} \right) \right] : \mathbf{d}^{\overline{ws}} + \frac{e^{w} \mathbf{q}^{\overline{w}}}{\left(\theta^{\overline{w}} \right)^{2}} \cdot \nabla \theta^{\overline{w}} + \frac{e^{ws} \mathbf{q}^{\overline{ws}}}{\left(\theta^{\overline{ws}} \right)^{2}} \cdot \nabla \theta^{\overline{ws}} \\ & + \frac{1}{\left(\theta^{\overline{s}} \right)^{2}} \left[e^{s} \mathbf{q}^{\overline{s}} - \left\langle \left(\mathbf{t}_{s} - \boldsymbol{\sigma}_{s} : \frac{\mathbf{C}_{s}}{j_{s}} \mathbf{I} \right) \cdot \left(\mathbf{v}_{s} - \mathbf{v}^{\overline{s}} \right) \right\rangle_{\Omega_{s},\Omega} \right] \cdot \nabla \theta^{\overline{s}} \\ & - \frac{ws \rightarrow w}{M} \frac{1}{\theta^{\overline{ws}}} \left[\left(\mu^{\overline{w}} + K_{\overline{E}}^{\overline{w}} + \psi^{\overline{w}} \right) - \left(\mu^{\overline{ws}} + K_{\overline{E}}^{\overline{ws}} + \psi^{\overline{ws}} \right) \right] \\ & - \frac{ws \rightarrow s}{M} \frac{1}{\theta^{\overline{ws}}} \left\{ \left(\mu^{\overline{s}} + K_{\overline{E}}^{\overline{s}} + \psi^{\overline{s}} \right) + \left\langle \frac{\boldsymbol{\sigma}_{s}}{\rho_{s}} : \frac{\mathbf{C}_{s}}{j_{s}} - \frac{1}{\rho_{s}} \mathbf{n}_{s} \cdot \mathbf{t}_{s} \cdot \mathbf{n}_{s} \right\rangle_{\Omega_{ws},\Omega_{ws}} \\ & - \left(\mu^{\overline{ws}} + K_{\overline{E}}^{\overline{ws}} + \psi^{\overline{ws}} \right) \right\} \\ & - \frac{1}{\theta^{\overline{w}}} \left\{ \frac{ws \rightarrow w}{\mathbf{T}} + \left(\frac{\mathbf{v}^{\overline{w},\overline{s}} - \mathbf{v}^{\overline{ws},\overline{s}}}{2} \right) \frac{ws \rightarrow w}{M} + e^{w} \rho^{w} \mathbf{g}^{\overline{w}} \\ & + e^{w} \rho^{w} \nabla \left(\psi^{\overline{w}} + \mu^{\overline{w}} + K_{\overline{E}}^{\overline{w}} \right) - \nabla \left(e^{w} p^{w} \right) + \eta^{\overline{w}} \nabla \theta^{\overline{w}} \right\} \cdot \mathbf{v}^{\overline{w},\overline{s}} \end{aligned}$$

$$\begin{split} & \mathbf{T}CA\mathbf{T} - -\mathbf{Simplified Entropy Inequality} \\ &+ \frac{1}{\theta^{\overline{ws}}} \Biggl\{ \sum_{\iota \in \Im_{p}} \left[\mathbf{T}^{ws \to \iota} + \left(\mathbf{Y}^{\overline{\iota},\overline{s}} - \mathbf{Y}^{\overline{ws},\overline{s}} \right) \mathbf{M}^{ws \to \iota} \right] - \epsilon^{ws} \rho^{ws} \mathbf{g}^{\overline{ws}} \\ &- \epsilon^{ws} \rho^{ws} \left(\mathbf{I} - \mathbf{G}^{ws} \right) \cdot \nabla \left(\mu^{\overline{ws}} + K_{E}^{\overline{ws}} + \psi^{\overline{ws}} \right) \\ &- \eta^{\overline{ws}} \left(\mathbf{I} - \mathbf{G}^{ws} \right) \cdot \nabla \theta^{\overline{ws}} - \nabla \cdot \left[\epsilon^{ws} \gamma^{ws} \left(\mathbf{I} - \mathbf{G}^{ws} \right) \right] \Biggr\} \cdot \mathbf{v}^{\overline{ws},\overline{s}} \\ &- \frac{1}{\theta^{\overline{ws}}} \langle p_w + \mathbf{n}_s \cdot \mathbf{t}_s \cdot \mathbf{n}_s + \gamma_{ws} \nabla' \cdot \mathbf{n}_s - \rho_{ws} \mathbf{g}_{ws} \cdot \mathbf{n}_s \rangle_{\Omega_{ws},\Omega_{ws}} \frac{\mathbf{D}^{\overline{s}} \epsilon^s}{\mathbf{D} t} \\ &+ \Biggl\{ \frac{ws \to w}{Q} + \left(\frac{E^{\overline{w}}}{\epsilon^w \rho^w} - \mu^{\overline{w}} \right)^{ws \to w} \mathbf{M}^w + \mathbf{v}^{\overline{w},\overline{s}} \cdot \left[\mathbf{T}^w + \left(\frac{\mathbf{v}^{\overline{w},\overline{s}} - \mathbf{v}^{\overline{w}\overline{s},\overline{s}} \right)^{ws \to w} \right] \\ &- \langle p_w \rangle_{\Omega_{ws},\Omega_{ws}} \frac{\mathbf{D}^{\overline{s}} \epsilon^s}{\mathbf{D} t} \Biggr\} \left(\frac{1}{\theta^{\overline{w}}} - \frac{1}{\theta^{\overline{ws}}} \right) \\ &+ \Biggl\{ \frac{ws \to s}{Q} + \left(\frac{E^{\overline{s}}}{\epsilon^s \rho^s} - \mu^{\overline{s}} \right)^{ws \to s} \mathbf{M} - \langle \mathbf{n}_s \cdot \mathbf{t}_s \cdot \mathbf{n}_s \rangle_{\Omega_{ws},\Omega_{ws}} \frac{\mathbf{D}^{\overline{s}} \epsilon^s}{\mathbf{D} t} \\ &- \Biggl\{ \frac{\sigma_s}{\rho_s} : \frac{\mathbf{C}_s}{j_s} - \frac{1}{\rho_s} \mathbf{n}_s \cdot \mathbf{t}_s \cdot \mathbf{n}_s \Biggr\} \frac{ws \to s}{\Omega_{ws,\Omega_{ws}}} \mathbf{M}^w \Biggr\} \left(\frac{1}{\theta^{\overline{s}}} - \frac{1}{\theta^{\overline{ws}}} \right) \\ &= \Lambda \ge 0 \end{split}$$

$$D^{\overline{\iota}}(\epsilon^{\iota}\rho^{\iota}) \qquad \iota \ \nabla \overline{\iota} \quad c = 0$$

$$\frac{\mathrm{D}(\epsilon \ \rho)}{\mathrm{D}t} = -\epsilon^{\iota} \rho^{\iota} \nabla \cdot \mathbf{v}^{\overline{\iota}}, \quad \text{for } \iota \in \mathfrak{I}_p$$

$$\frac{\mathrm{D}^{\overline{\iota}}\left(\epsilon^{w}\rho^{w}\mathbf{v}^{\overline{w}}\right)}{\mathrm{D}t} = -\epsilon^{w}\rho^{w}\mathbf{v}^{\overline{w}}\nabla\cdot\mathbf{v}^{\overline{w}} - \hat{\mathbf{R}}^{w}\cdot\mathbf{v}^{\overline{w},\overline{s}} - \epsilon^{w}\rho^{w}\nabla\left(\psi^{\overline{w}} + \mu^{\overline{w}}\right)$$

.

$$\frac{\mathrm{D}^{\overline{\iota}}\left(\epsilon^{s}\rho^{s}\mathbf{v}^{\overline{s}}\right)}{\mathrm{D}t} = -\epsilon^{s}\rho^{s}\mathbf{v}^{\overline{s}}\nabla\cdot\mathbf{v}^{\overline{s}} + \nabla\cdot\left(\epsilon^{s}\mathbf{t}^{\overline{s}}\right) + \epsilon^{s}\rho^{s}\mathbf{g}^{\overline{s}} + \epsilon^{w}\rho^{w}\mathbf{g}^{\overline{w}} + \epsilon^{w}\rho^{w}\mathbf{g}^{\overline{w}} + \epsilon^{w}\rho^{w}\nabla\left(\psi^{\overline{w}} + \mu^{\overline{w}}\right) - \nabla\left(\epsilon^{w}p^{w}\right) + \hat{\mathbf{R}}^{w}\cdot\mathbf{v}^{\overline{w},\overline{s}}$$

Characteristics of Behavior

- NAPLs leave a state of residual saturation in media through which they pass
- NAPLs follow a complex pattern of flow, which is importantly influenced by media heterogeneity
- LNAPLs accumulate on the top of the water table
- > DNAPLs can sink below the water
- NAPLs often reach stable configurations of locally high saturations known as pools
- NAPLs are usually sparingly soluble and DNAPL contaminants usually degrade slowly---thus are long lived in the environment

Two-Dimensional Unsaturated Downward Vertical Displacement of TCE



- •21-cm x 21-cm two-dimensional cell
- •Pooled TCE established
- •TCE dyed with Oil Red O for visualization
- •Established bottom brine layer
- •Drained to unsaturated conditions
- •0.3 pore-volume downward flush with mixture of sulfosuccinate surfantants
- •Measured 80.0% TCE removal, no visible pools
- •Reference: Hill et. al. [ES&T, 35(14), 2001]